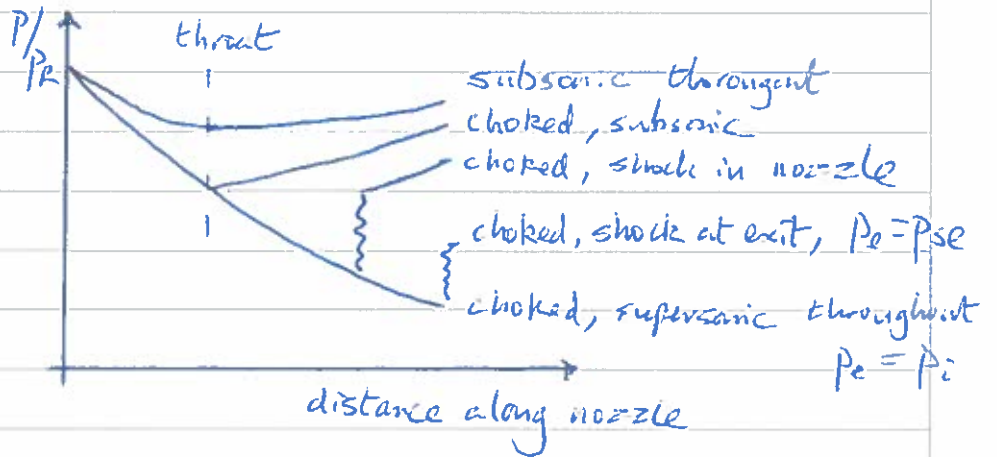


a) (i)

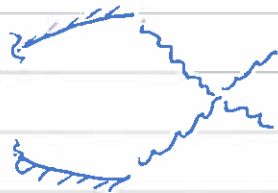
$P_R$ : reservoir pressure

$P_e$ : nozzle exit pressure

$P_v$ : vessel pressure

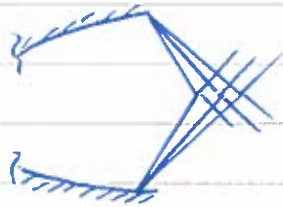


(ii) Over-expanded:  $P_i < P_v < P_{se}$



shocks in exit jet

Under-expanded:  $P_v < P_i$



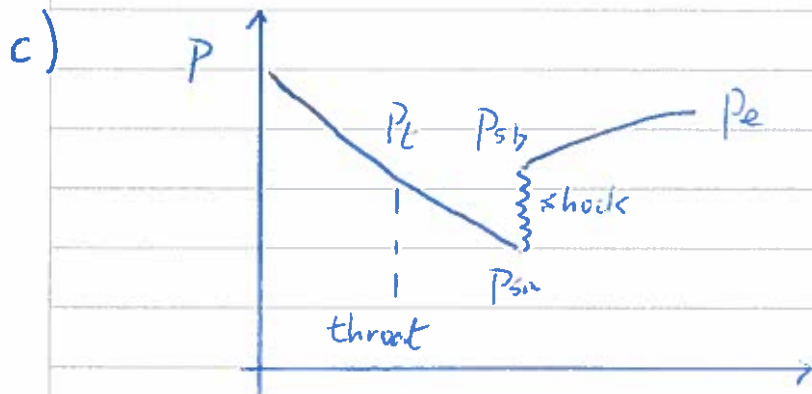
expansions in exit jet

b) (i) Choked:  $\frac{m \sqrt{c_p T_0}}{A P_R} = 1.281 \Rightarrow$  exit  $\frac{m \sqrt{c_p T_0}}{A P_R} = 0.427$

Tables (subsonic):  $M = 0.195$   $\frac{P_e}{P_R} = 0.974$

(ii) Tables (supersonic):  $M = 2.6375$   $\frac{P_e}{P_R} = 0.0473$

(iii) Shock tables:  $P_{sp} = 7.9645$   $\frac{P_0}{P_R} = 0.377$



Loss in  $p_0 \Rightarrow$  shock present with  $M = 1.89$   
(and nozzle choked)

$$\frac{P_t}{P_R} = \underline{\underline{0.528}}$$

$$\frac{P_{sa}}{P_R} = \underline{\underline{0.152}} \quad (\text{tables})$$

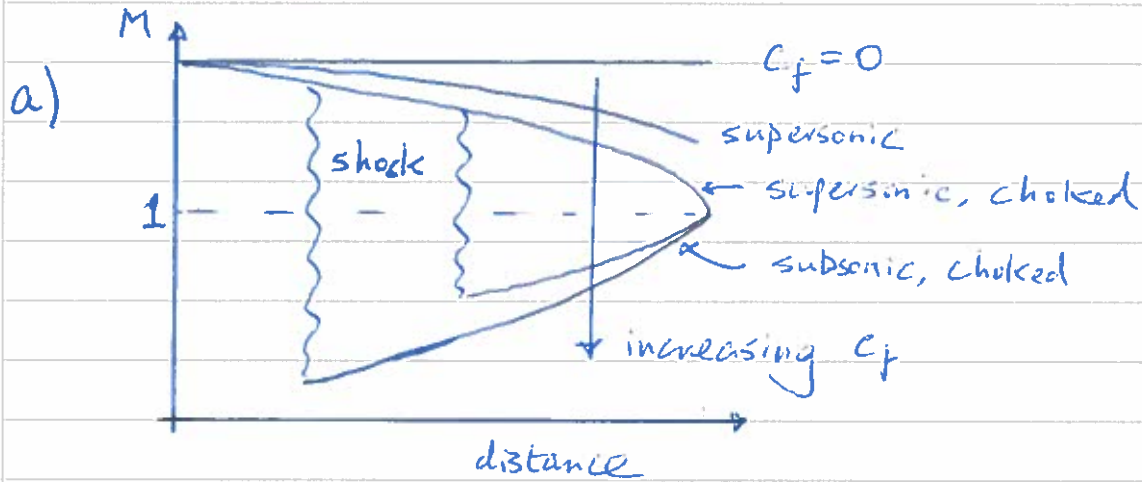
Shock:  $\frac{P_{sb}}{P_{sa}} = 4.0008$

$$\frac{P_{sb}}{P_R} = \underline{\underline{0.608}}$$

At exit  $\frac{m \sqrt{c_p \bar{T}_0}}{A p_0} = \frac{m \sqrt{c_p \bar{T}_0}}{A p_R} \frac{p_R}{p_0} = 0.5531 \Rightarrow M = 0.260$   
(tables)

$$\frac{P_e}{P_0} = 0.954$$

$$\frac{P_e}{P_R} = \frac{P_e}{P_0} \frac{P_0}{P_R} = \underline{\underline{0.737}}$$



b) (i) Exit choked :  $\frac{\dot{m} \sqrt{C_p T_0}}{A_{pe}} = 2.4249$  (tables)

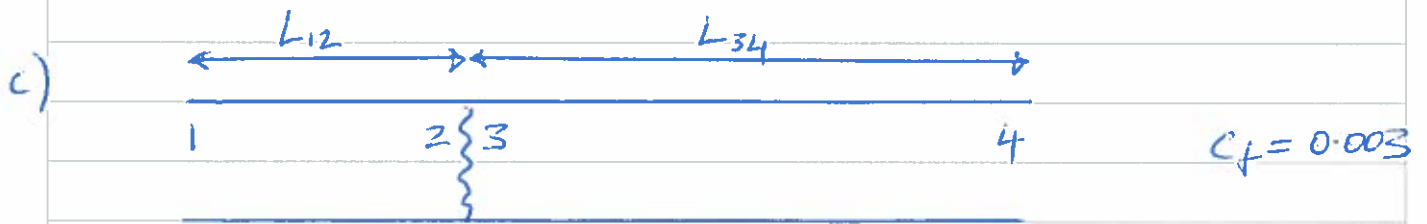
Inlet  $\frac{\dot{m} \sqrt{C_p T_0}}{A_{p0}} = \frac{p_e}{p_0} \frac{\dot{m} \sqrt{C_p T_0}}{A_{pe}} = 0.8981$

Inlet  $M$  : Subsonic  $0.459$       Supersonic  $1.788$  (tables)

(ii)  $\frac{4c_f L_{max}}{D}$  :  $1.4594$        $0.23796$

$L_{max} = 5.9 \text{ m}$  ,       $D = 0.2 \text{ m}$

$\Rightarrow c_f = \underline{\underline{0.0124}}$       or       $\underline{\underline{0.0020}}$



with  $\frac{4C_f L_{max}}{D} = 0.23796$   $L_{max} = 3.966 \text{ m now}$

At 2, for  $L_{12} = 0.8 \text{ m}$   $L_{max} = 3.166 \text{ m}$   $\frac{4C_f L_{max}}{D} = 0.18996$   
 $L_{34} = 5.1 \text{ m}$

Tables:  $M = 1.65$   $M_s = 0.654$   $\frac{4C_f (L_{max})_3}{D} = 0.314$

$(L_{max})_3 = 5.2 \text{ m} > L_{34}$ . Not choked, so  $L_{12}$  too small.

Try  $L_{12} = 1.0 \text{ m}$   $(L_{max})_2 = 2.966 \text{ m}$ ,  $\left(\frac{4C_f L_{max}}{D}\right)_2 = 0.1780$   
 $L_{34} = 4.9 \text{ m}$

Tables:  $M = 1.616$   $M_s = 0.6638$   $\frac{4C_f (L_{max})_3}{D} = 0.2884$

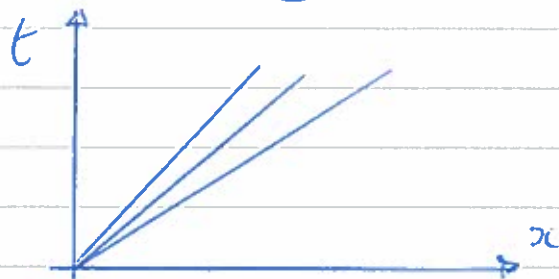
$(L_{max})_3 = 4.8 \text{ m} < L_{34}$  Predicts choked before end, so  $L_{12}$  too large.

$\therefore$  True  $L_{12}$  between  $0.8 \text{ m}$  and  $1.0 \text{ m}$

≡

a) See lecture notes.

b) (i) Expansion wave system propagates rightwards from  $x=0$ :



(ii)  $p_e = 1 \text{ bar}$  (e)  $\rightarrow$  (P)  $p_p = 2.5 \text{ bar}$   $V_p = 0 \text{ m/s}$   
 $T_p = 283.15 \text{ K}$

$$a = \sqrt{\gamma R T} \Rightarrow a_p = 337.3 \text{ m/s}$$

$$T_e = \left( \frac{p_e}{p_p} \right)^{\frac{\gamma-1}{\gamma}} T_p \quad (\text{isentropic}) \Rightarrow T_e = 217.9 \text{ K} \quad a_e = 295.9 \text{ m/s}$$

$$\text{Across R-R waves: } V_e - \frac{2}{\gamma-1} a_e = V_p - \frac{2}{\gamma-1} a_p$$

$$V_e = -206.9 \text{ m/s}$$

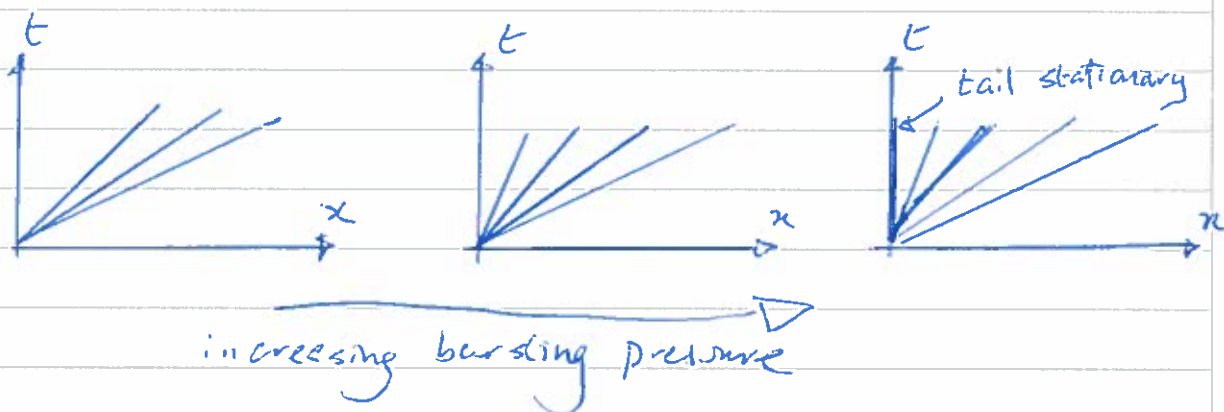
$$\dot{m} = \rho_e A |V_e| \quad \text{with} \quad \rho_e = \frac{p}{R T_e} = 1.599 \text{ kg/m}^3$$

$$A = 0.03142$$

$$\dot{m} = \underline{\underline{10.4 \text{ kg/s}}}$$

c) (i) Higher pressure ratio  $\Rightarrow$  lower  $T_e \Rightarrow$  lower  $a_e$

Speed of tail wave ( $a_e + V_e$ ) reduces.



(ii) Max. possible mass flow when tail stationary (choked)

$$V_e = -a_e \Rightarrow -a_e - \frac{2}{\gamma-1} a_e = -\frac{2}{\gamma-1} a_p ; a_e = 281.1 \text{ m/s}$$

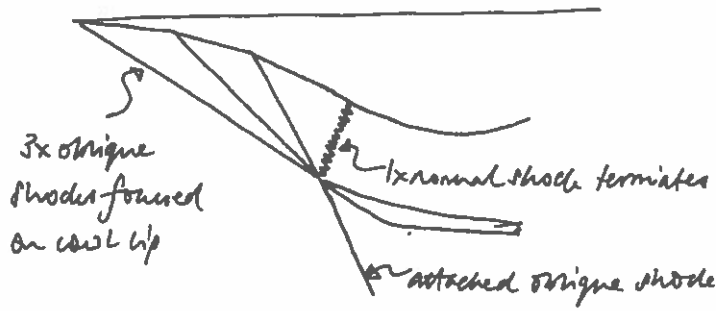
$$T_e = a_e^2 / \gamma R = 196.6 \text{ K} \quad P_p = \left( \frac{T_p}{T_e} \right)^{\frac{\gamma}{\gamma-1}} P_e = 3.583 \text{ bar}$$

$$\rho_e = \frac{P_e}{RT_e} = 1.772 \text{ kg/m}^3 \quad \dot{m} = 15.65 \text{ kg/s} \quad \text{at} \quad P_p = 3.58 \text{ bar}$$

(iii) Further expansion in pipe not possible; exit pressure rises above ambient and expansion continues outside:



4a)(i)



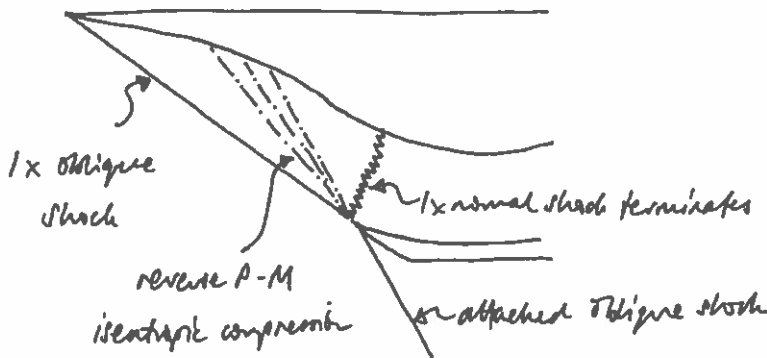
$M = 2.40 \Rightarrow \theta_{max} = 28.7^\circ$

$3 \times 8^\circ \Rightarrow 24^\circ$  flow turning by oblique shocks.

$4^\circ$  cow lip wedge angle

$\hookrightarrow$  cow lip @  $28^\circ$  to free stream

$\therefore$  attached shock possible



Isentropic compression (i.e. reversed Prandtl-Meyer expansion) replaces oblique shocks 2 & 3.

ii) x iii)

$M_{in}$	$\delta$	$M_{out}$	$p_1/p$	$p_{01}/p_0$
2.40	$8^\circ$	2.08	1.63	0.988
2.08	$8^\circ$	1.79	1.56	0.991
1.79	$8^\circ$	1.51	1.50	0.993
1.51	normal		2.49	0.927
			<u>9.50</u>	<u>90.1%</u>

For case (b) replace shocks 2 & 3 with reversed P-M, shock 1 as before.

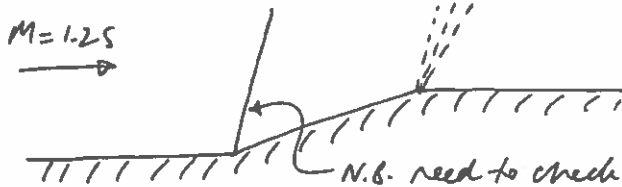
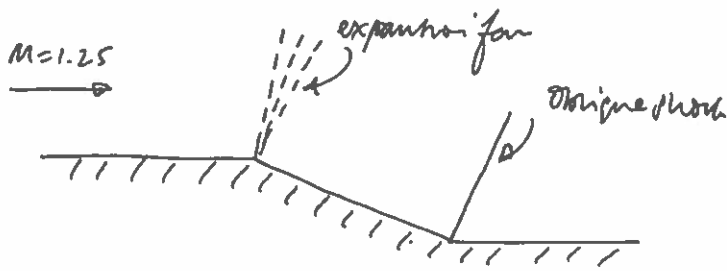
$M$	$\nu$	$p/p_0$	$p_{01}/p_0$	$p_1/p$
2.08	28.6	0.113		
1.52	12.6	0.263		
1.52	normal		0.922	2.54
			<u>91.4%</u>	<u>9.65</u>

$\left\{ \frac{P^+}{\Delta p/p} = 2.33 \right. \quad [0.263/0.113]$

Case (b) is superior design, so use this as starting point.

- b) 1) Reduce leading ramp angle
- 2) Increase total turning ahead of normal shock (but note implications for cow drag)

5a) (i)



N.B. need to check that  $4^\circ$  turning is possible with attached shock.

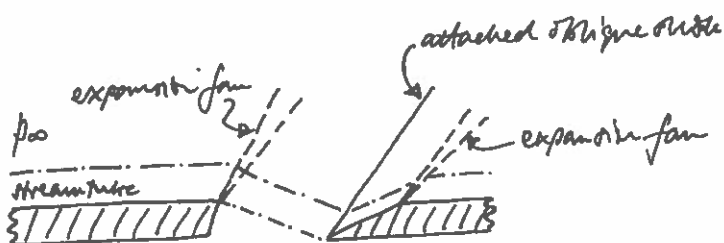
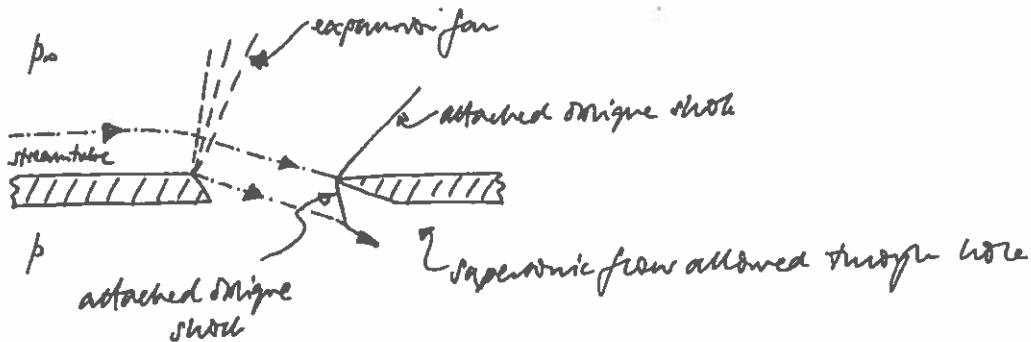
(ii) Case (a)  $M = 1.25$   $\nu = 4.83 \Rightarrow 4^\circ$  turning fan  $\Rightarrow \nu = 8.83$ .  $M = 1.394$ ,  $\frac{p_0}{p_{0a}} = 0.99920$

Case (b)  $M = 1.25$   $\frac{p_0}{p_{0a}} = 0.99882$ .

(iii) Case (b) more affected as shock detaches from corner @  $M = 1.15$ , but remains attached in case (a) [for latter,  $\nu_2 = 6.38 \Rightarrow M_2 > 1.3 \Rightarrow \theta_{max} > 6.662^\circ$ ]

b)

(i)



no flow through hole at this supersonic flow condition:

(ii) For a set combination of  $M$  and  $p/p_0$  the hole in case (d) can be "educated" to aerodynamically close, while allowing flow through in a subsonic condition. It can be used as a bleed valve with no moving parts or outside control system.



$$a) (i) \quad u_{i-1} = u_i - \frac{\partial u}{\partial x} \Delta x + \frac{1}{2} \frac{\partial^2 u}{\partial x^2} \Delta x^2 - \frac{1}{6} \frac{\partial^3 u}{\partial x^3} \Delta x^3 + O(\Delta x^4)$$

$$u_{i-2} = u_i - 2 \frac{\partial u}{\partial x} \Delta x + 2 \frac{\partial^2 u}{\partial x^2} \Delta x^2 - \frac{4}{3} \frac{\partial^3 u}{\partial x^3} \Delta x^3 + O(\Delta x^4)$$

$$\text{so } 3u_i^n - 4u_{i-1}^n + u_{i-2}^n = 2 \frac{\partial u}{\partial x} \Delta x - \frac{2}{3} \frac{\partial^3 u}{\partial x^3} \Delta x^3 + \dots$$

and difference eqn is

$$\frac{\partial u}{\partial t} \Delta t + \frac{\partial^2 u}{\partial x^2} \frac{\Delta t^2}{2} = -a \Delta t \left[ \frac{\partial u}{\partial x} - \frac{1}{3} \frac{\partial^3 u}{\partial x^3} \Delta x^2 \right]$$

$$\frac{\partial u}{\partial t} + a \frac{\partial u}{\partial x} = - \frac{\partial^2 u}{\partial x^2} \frac{\Delta t}{2} + \frac{\partial^3 u}{\partial x^3} \frac{a \Delta x^2}{3}$$

1st order in time

2nd order in space

(ii) Points that can be made (not all required):

- For  $a > 0$ , this is an upwinding scheme; hence expect stability for sufficiently small  $a \Delta t / \Delta x$
- $O(\Delta t)$  term can be written  $-a^2 \frac{\partial^2 u}{\partial x^2} \frac{\Delta t}{2}$ , which represents negative diffusion  $\Rightarrow$  destabilising

$O(\Delta x^2)$  term represents dispersion, which can cause oscillatory behaviour

- a) (ii) cont'd • A sawtooth stability analysis, with  $u_i^n = \varepsilon^n (-1)^i$ , yields

$$\left| \frac{u_i^{n+1}}{\varepsilon^n} \right| = \left| 1 - 4a \frac{\Delta t}{\Delta x} \right|$$

suggesting stability for  $a > 0$  and  $a \Delta t / \Delta x < 1/2$

b) 
$$u_{i+1} - u_{i-1} = 2 \frac{\partial u}{\partial x} \Delta x + \frac{2}{6} \frac{\partial^3 u}{\partial x^3} \Delta x^3 + O(\Delta x^5)$$

$$u_{i+1} - 2u_i + u_{i-1} = \frac{2}{2} \frac{\partial^2 u}{\partial x^2} \Delta x^2 + O(\Delta x^4)$$

Difference eqn becomes

$$\frac{\partial u}{\partial t} \Delta t + \frac{\partial^2 u}{\partial t^2} \frac{\Delta t^2}{2} = -a \Delta t \left[ \frac{\partial u}{\partial x} + \frac{\partial^3 u}{\partial x^3} \frac{\Delta x^2}{6} \right] + \beta \frac{\partial^2 u}{\partial x^2} \Delta x^2$$

with 1st-order error  $\beta \frac{\Delta x^2}{\Delta t} \frac{\partial^2 u}{\partial x^2} - \frac{1}{2} \frac{\partial^2 u}{\partial t^2} \Delta t$

From underlying PDE,  $\frac{\partial^2 u}{\partial t^2} = a^2 \frac{\partial^2 u}{\partial x^2}$  to leading order

$$\text{1st-order error} = \left( \beta \frac{\Delta x^2}{\Delta t} - \frac{a^2 \Delta t}{2} \right) \frac{\partial^2 u}{\partial x^2}$$

which is eliminated for  $\beta = \frac{1}{2} \left( \frac{a \Delta t}{\Delta x} \right)^2$

(i) Elliptic

(ii) Parabolic

(iii) Hyperbolic

See lecture notes for commentary on boundary/initial conditions and solution methods.

# 7b) Half Question

①/①

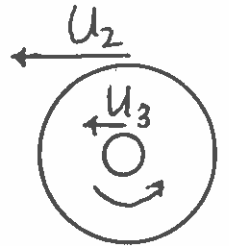
2 - rotor inlet  
3 - rotor exit

$$h_{03} - h_{02} = \frac{1}{2}(U_3^2 - U_2^2) - \frac{1}{2}(V_{3,rel}^2 - V_{2,rel}^2) + \frac{1}{2}(V_3^2 - V_2^2)$$

large power output  $\Rightarrow h_{03} - h_{02}$  large negative.

(i)  $\frac{1}{2}(U_3^2 - U_2^2)$  LARGE NEGATIVE  $\Rightarrow U_3 < U_2$

$\Rightarrow$  RADIAL INFLOW TURBINE

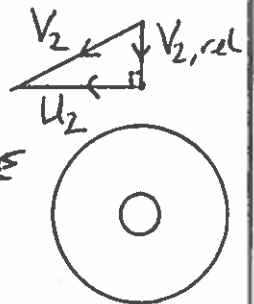


(ii)  $-\frac{1}{2}(V_{3,rel}^2 - V_{2,rel}^2)$  LARGE NEGATIVE

$\Rightarrow V_{2,rel}$  SMALL AS POSSIBLE

$\Rightarrow$  RADIAL FLOW IN RELATIVE FRAME

$\Rightarrow \alpha_{2,rel} = 0, V_{r2} = U_2$

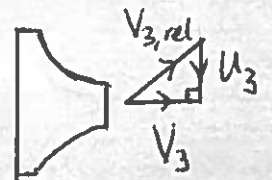


(iii)  $\frac{1}{2}(V_3^2 - V_2^2)$  SMALL AS POSSIBLE,  $\dot{\omega}$  negative

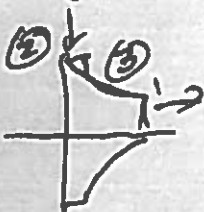
$\Rightarrow V_3^2$  MINIMUM

$\Rightarrow \alpha_3 = 0$

NO EXIT SWIRL  $V_{\theta 3}$



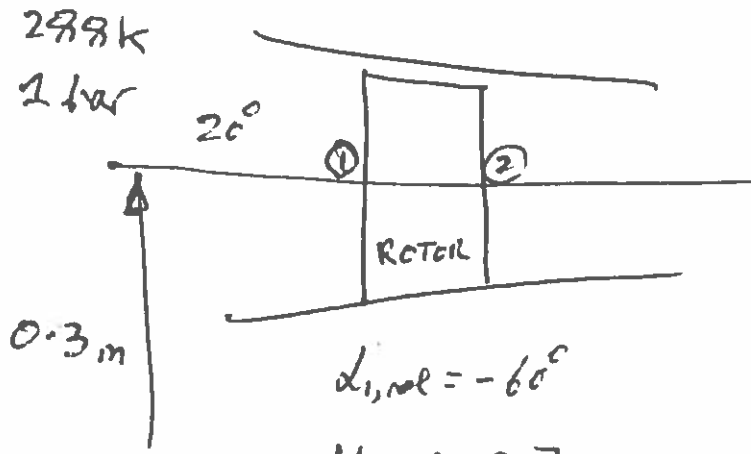
(iv)  $\eta = \frac{U_2 V_{\theta 2} - U_3 V_{\theta 3}}{U_2^2} = \frac{U_2 \cdot U_2 - 0}{U_2^2} = 1$



N.B. Also derivable from expression given for  $\Delta h_0$ , using velocity triangles as sketched.

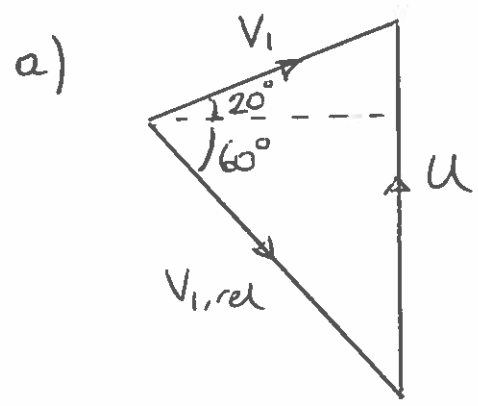
8)

①/④



$v_x = \text{const.}$

$\alpha_{1,rel} = -60^\circ$   
 $M_{1,rel} = 0.7$   
 $\gamma_p = 0.04$   
 $\alpha_{2,rel} = -30^\circ$

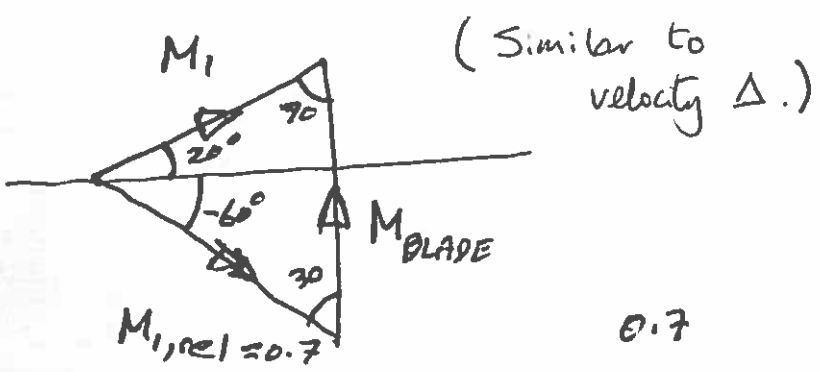


b)  $\tan d_{rel} = \tan d - \frac{1}{\phi}$

$\phi = (\tan d - \tan d_{rel})^{-1}$

$\phi = (\tan 20^\circ - \tan(-60^\circ))^{-1} = \underline{\underline{0.4771}}$

c) (k)



$\frac{M_{1,rel}}{\sin 70^\circ} = \frac{M_1}{\sin 30^\circ}$

$0.7$   
 $\downarrow$   
 $M_1 = M_{1,rel} \frac{\sin 30^\circ}{\sin 70^\circ}$

$M_1 = \underline{\underline{0.3725}}$

(d)

$$\gamma = 1.4 \quad R = 287 \text{ J/kgK}$$

$$C_p = 1005 \text{ J/kgK}$$

(2/4)

$$\frac{T_{01}}{T_1} = \left(1 + \frac{\gamma-1}{2} M_1^2\right)$$

$$T_0 = \frac{T_{01}}{1 + \frac{\gamma-1}{2} 0.3725^2} = \frac{288}{1.0278} = \underline{\underline{280.2 \text{ K}}}$$

$$P_1 = \left(\frac{T_1}{T_{01}}\right)^{\frac{\gamma}{\gamma-1}} P_{01} = 1 \text{ bar} \left(\frac{280.2}{288.0}\right)^{3.5} = \underline{\underline{0.9084 \text{ bar}}}$$

$$V_1 = M_1 \sqrt{\gamma R T_1} = 0.3725 \sqrt{1.4 \times 287 \times 280.2}$$

$$\underline{\underline{V_1 = 124.9 \text{ m/s}}}$$

$$V_{x1} = V_1 \cos \alpha_1 = 124.9 \times \cos 20^\circ = \underline{\underline{117.4 \text{ m/s}}}$$

$$V_{z1} = V_1 \sin \alpha_1 = 124.9 \times \sin 20^\circ = \underline{\underline{42.7 \text{ m/s}}}$$

$$\text{(e)} \quad \phi = \frac{V_x}{U} \Rightarrow U = \frac{V_x}{\phi} = \frac{117.4}{0.4771} = \underline{\underline{246.1 \text{ m/s}}}$$

$$U = r \Omega \Rightarrow \Omega = \frac{U}{r} = \frac{246.1}{0.3} = \underline{\underline{820.3 \text{ rad/s}}}$$

(7833.6 rpm)

(3)/4

(f)

$$\frac{T_{01,rel}}{T_1} = \left(1 + \frac{\gamma-1}{2} M_{1,rel}^2\right)$$

$$T_{01,rel} = 280 \cdot 2 \left(1 + \frac{\gamma-1}{2} 0.7^2\right) = 280 \cdot 2 \times 1.098$$

$$= \underline{\underline{307.7 \text{ K}}}$$

$$P_{02,rel} = P_1 \left(\frac{T_{02,rel}}{T_1}\right)^{\frac{\gamma}{\gamma-1}} = 0.9084 \times (1.098)^{3.5} = \underline{\underline{1.2600 \text{ bar}}}$$

Flow as fixed/constant radius

$$\Rightarrow T_{02,rel} = T_{01,rel} = \underline{\underline{307.7 \text{ K}}}$$

LOSS COEFFICIENT

$$0.04 = \gamma_p = \frac{P_{01,rel} - P_{02,rel}}{P_{01,rel} - P_1} \quad (\text{COMPRESSOR})$$

$$P_{02,rel} = P_{01,rel} - 0.04(P_{01,rel} - P_1)$$

$$= 1.2600 - 0.04(1.2600 - 0.9084)$$

$$\underline{\underline{P_{02,rel} = 1.2459 \text{ bar}}}$$

$$v_{x2} = v_{x1} = 117.4 \text{ m/s}$$

$$v_{\theta 2,rel} = v_{x2} \tan \alpha_{2,rel} = 117.4 \tan(-30^\circ)$$

$$\underline{\underline{v_{\theta 2,rel} = -67.8 \text{ m/s}}}$$

(4) (4)

(f) cont.

$$T_2 = T_{02,rel} - \frac{(v_{x2}^2 + v_{\theta 2,rel}^2)}{2c_p}$$

$$T_2 = 307.7 - \frac{(117.4^2 + (-67.8)^2)}{2 \times 1005} = \underline{298.6 \text{ K}}$$

~~$$P_2 = P_{02,rel} \left( \frac{T_2}{T_{02,rel}} \right)^{\frac{\gamma}{\gamma-1}}$$~~

$$P_2 = P_{02,rel} \left( \frac{T_2}{T_{02,rel}} \right)^{\frac{\gamma}{\gamma-1}} = 1.2459 \left( \frac{298.6}{307.7} \right)^{3.5} = \underline{1.1216 \text{ bar}}$$

$$\Rightarrow \frac{P_2}{P_1} = \frac{1.1216}{0.9084} = \underline{\underline{1.235}}$$

$$v_{\theta 2} = U + v_{\theta 2,rel} = 246.1 - 67.8 = \underline{178.3 \text{ m/s}}$$

$$\text{stage loading} = \frac{\Delta h_0}{U^2} = \frac{U(v_{\theta 2} - v_{\theta 1})}{U^2}$$

$$\psi = \frac{178.3 - 42.7}{246.1} = \underline{\underline{0.55}}$$

This is rather high, would expect  $\psi \approx 0.4$