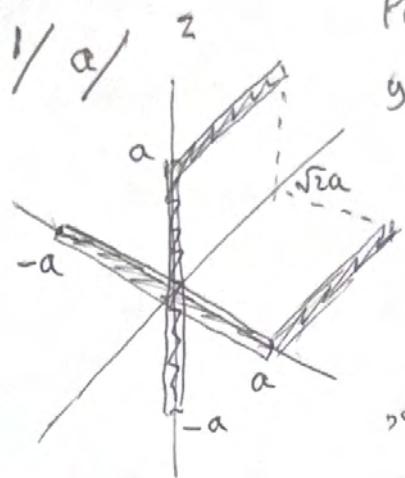




x-view

$$I_{xx} = \frac{1}{12} m (2a)^2 \quad (CD)$$

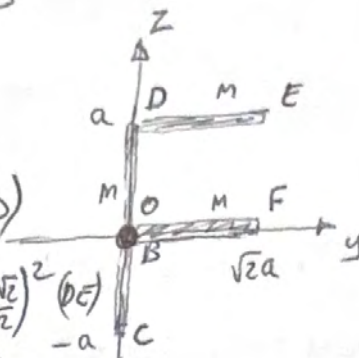
$$+ \frac{1}{12} m (a\sqrt{2})^2 + ma^2 \left(1 + \frac{\sqrt{2}}{2}\right)^2 \quad (DE)$$

$$+ \frac{1}{12} m (a\sqrt{2})^2 + ma^2 \left(\frac{\sqrt{2}}{2}\right)^2 \quad (BF)$$

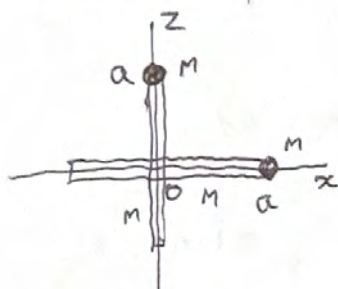
$$= ma^2 \left(\frac{1}{3} + \frac{1}{6} + 1 + \frac{1}{2} + \frac{1}{6} + \frac{1}{2} \right)$$

$$= \frac{8}{3} ma^2$$

$$I_{yz} = ma \frac{a\sqrt{2}}{2} = ma^2 \frac{\sqrt{2}}{2}$$



y-view

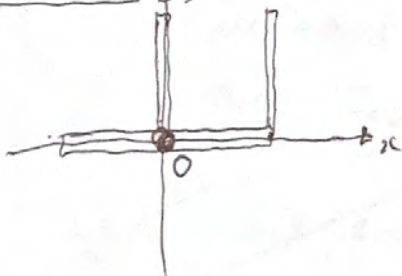


$$I_{yy} = \frac{1}{12} m (2a)^2 + \frac{1}{12} m (2a)^2 + ma^2 + ma^2$$

$$= \frac{8}{3} ma^2$$

$$I_{xz} = 0$$

z-view



Same as x-view

$$I_{zz} = \frac{8}{3} ma^2$$

$$I_{xy} = ma^2 \frac{\sqrt{2}}{2}$$

$$\text{So } I_{xx} = I_{yy} = I_{zz} = \frac{8}{3} ma^2$$

$$I_0 = ma^2 \begin{bmatrix} \frac{8}{3} & -\frac{\sqrt{2}}{2} & 0 \\ -\frac{\sqrt{2}}{2} & \frac{8}{3} & -\frac{\sqrt{2}}{2} \\ 0 & -\frac{\sqrt{2}}{2} & \frac{8}{3} \end{bmatrix}$$

N.B. this is NOT an AAA body because I_{yz} and I_{xy} are not zero. Need to use the whole inertia matrix for part (b)

(b) We're told that G is at $\frac{1}{4}(a, \sqrt{2}a, a)$
and that OG is a principal axis

Verify using

$[I_O] \underline{u} = \lambda \underline{u}$ where $\underline{u}$ is the principal axis
and λ is the principal moment
of inertia

$$ma^2 \begin{bmatrix} \frac{8}{3} & -\frac{\sqrt{2}}{2} & 0 \\ -\frac{\sqrt{2}}{2} & \frac{8}{3} & -\frac{\sqrt{2}}{2} \\ 0 & -\frac{\sqrt{2}}{2} & \frac{8}{3} \end{bmatrix} \begin{bmatrix} 1 \\ \sqrt{2} \\ 1 \end{bmatrix} = ma^2 \begin{bmatrix} \frac{8}{3} - 1 \\ -\frac{\sqrt{2}}{2} + \frac{8\sqrt{2}}{3} - \frac{\sqrt{2}}{2} \\ -1 + \frac{8}{3} \end{bmatrix} = ma^2 \begin{bmatrix} \frac{5}{3} \\ \sqrt{2}(\frac{8}{3} - 1) \\ \frac{5}{3} \end{bmatrix} = \frac{5}{3} \begin{bmatrix} 1 \\ \sqrt{2} \\ 1 \end{bmatrix} ma^2$$

So yes, OG is principal and I_{OG} is $\frac{5}{3}ma^2$

note, don't assume it's an "AAA" body
just because $I_{xx} = I_{yy} = I_{zz} = \frac{8}{3}ma^2$ are
all equal. If ~~that~~ it was an AAA body
then I_{OG} would be $\frac{8}{3}ma^2$

~~(c) we're told that the framework is "AAC" at G and
but we can deduce that the value " $\frac{5}{3}ma^2$ " that
we found in (b) ~~cannot be~~ ^{is not} the "C" value because
if it were then by parallel axes theorem the body
would be AAC~~

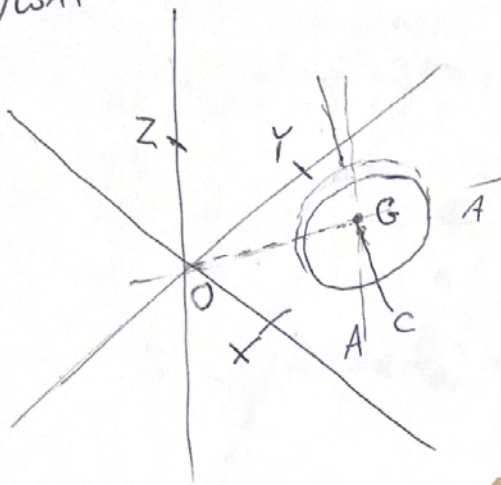
~~we~~ we have principal moments of inertia at G

as $\frac{5}{3}ma^2, \frac{5}{3}ma^2, \frac{8}{3}ma^2$

~~so the~~ which is "AAC"

so the "C" value is $\frac{8}{3}ma^2$

(C) cont



So to find the principle moments of inertia at O
use parallel axes theorem
on A and on C

$$\begin{aligned}\text{with } OG &= \sqrt{x^2 + y^2 + z^2} \\ &= \sqrt{\left(\frac{a}{4}\right)^2 + \left(\frac{a\sqrt{2}}{4}\right)^2 + \left(\frac{a}{4}\right)^2} \\ &= \frac{a}{4} (1 + 2 + 1)^{\frac{1}{2}} \\ &= \frac{a}{2}\end{aligned}$$

so at O principal moments of inertia

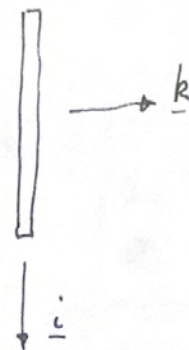
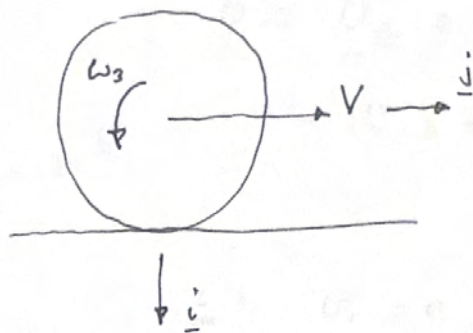
are ~~$\frac{5}{3}ma^2$~~ $\frac{5}{3}ma^2$, $\frac{5}{3}ma^2 + 4m\left(\frac{a}{2}\right)^2$, $\frac{8}{3}ma^2 + 4m\left(\frac{a}{2}\right)^2$

$\uparrow$
already found

$= \frac{8}{3}ma^2$

$= \frac{11}{3}ma^2$

2 (a)



Steady state

$$\underline{U} = V \underline{i}$$

$$U_1 = 0 \quad U_2 = V \quad U_3 = 0$$

$$\underline{\omega} = \omega_3 \underline{k} = -\frac{V}{a} \underline{k}$$

$$\omega_1 = 0 \quad \omega_2 = 0 \quad \omega_3 = -\frac{V}{a}$$

$$\underline{F} = -mg \underline{i}$$

$$F_1 = -mg \quad F_2 = 0 \quad F_3 = 0$$

$$\underline{R} = 0$$

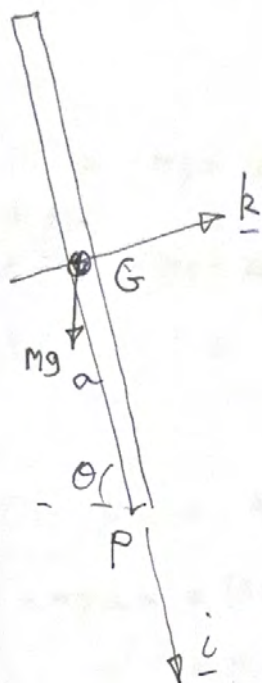
$$R_3 = 0$$

This means that in perturbed motion ($\theta = \frac{\pi}{2} - \alpha$)

these quantities are small :

$$\alpha, U_1, U_3, \omega_1, \omega_2, F_2, F_3, R_3$$

2 (b)(i) No slip at P



$$0 = \underline{v}_G + \underline{\omega} \times a \underline{i}$$

$$= (v_1 \underline{i} + v_2 \underline{j} + v_3 \underline{k}) + \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ \omega_1 & \omega_2 & \omega_3 \\ a & 0 & 0 \end{vmatrix}$$

$$\underline{i} : v_1 = 0$$

$$\underline{j} : v_2 + a\omega_3 = 0$$

$$\underline{k} : v_3 - a\omega_2 = 0$$

$$\begin{cases} v_1 = 0 \\ v_2 = -a\omega_3 \\ v_3 = a\omega_2 \end{cases}$$

(ii) $\underline{Q}_P = \underline{h}_P + \underline{r}_P \times \underline{f}$

$$\underline{f} = ma(-\omega_3 \underline{j} + \omega_2 \underline{k})$$

$$\underline{r}_P = \underline{v}_G + \underline{\Omega} \times a \underline{i}$$

$$= (v_2 + a\Omega_3) \underline{j} + (v_3 - a\omega_2) \underline{k}$$

$$\underline{r}_P \times \underline{f} = [ma\omega_2(v_2 + a\Omega_3) + ma\omega_3(v_3 - a\omega_2)] \underline{i}$$

$$= ma^2 \underline{i} (-\omega_2 \omega_3 + \omega_2 \Omega_3 + \omega_3 \omega_2 - \omega_3 \omega_2)$$

$$= ma^2 \underline{i} (\Omega_3 - \omega_3) \omega_2$$

Fast spin $\omega_3 \gg \Omega_3 \therefore \underline{r}_P \times \underline{f} = -ma^2 \omega_3 \omega_2 \underline{i}$

$$\underline{h}_P = \underline{h}_G + (\underline{r}_G - \underline{r}_P) \times \underline{f}$$

$$= A\omega_1 \underline{i} + A\omega_2 \underline{j} + C\omega_3 \underline{k} - (a \underline{i}) \times ma(-\omega_3 \underline{j} + \omega_2 \underline{k})$$

$$= A\omega_1 \underline{i} + (A\omega_2 + ma^2 \omega_2) \underline{j} + (C\omega_3 + ma^2 \omega_3) \underline{k}$$

$$= A\omega_1 \underline{i} + (A + ma^2) \omega_2 \underline{j} + (C + ma^2) \omega_3 \underline{k}$$

$$\underline{h}_P = \underline{h}_P|_{rot} + \underline{\Omega} \times \underline{h}_P$$

$$\vec{\Omega} \times \vec{h}_P = \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ \omega_1 & \omega_2 & \omega_3 \\ A\omega_1 & (A+ma^2)\omega_2 & (c+ma^2)\omega_3 \end{vmatrix}$$

$$\begin{aligned} \textcircled{1} \quad \underline{Q}_P &= -mga \cos \theta \quad \underline{j} \\ \textcircled{2} \quad \therefore \begin{cases} \underline{i}: & 0 \\ \underline{j}: & -mga \cos \theta \\ \underline{k}: & 0 \end{cases} &= \begin{cases} A\dot{\omega}_1 + (c+ma^2)\omega_2\omega_3 - (A+ma^2)\omega_2\omega_3 - ma^2\omega_2\omega_3 \\ (A+ma^2)\dot{\omega}_2 + A\omega_1\omega_3 - (c+ma^2)\omega_1\omega_3 \\ (c+ma^2)\dot{\omega}_3 + (A+ma^2)\omega_1\omega_2 - A\omega_1\omega_2 \end{cases} \end{aligned}$$

$$\begin{aligned} \textcircled{1} \quad \therefore \begin{cases} 0 &= A\dot{\omega}_1 + c\omega_2\omega_3 \\ -mga \alpha &= (A+ma^2)\dot{\omega}_2 - (c+ma^2)\omega_1\omega_3 \\ 0 &= (c+ma^2)\dot{\omega}_3 \end{cases} & \begin{pmatrix} \omega_2 \omega_3 \ll \omega_3 \\ \text{is small} \\ \cos \theta = \sin \alpha = 1 \\ \omega_3 \ll \omega_3 \end{pmatrix} \end{aligned}$$

$$\textcircled{3} \quad \therefore \omega_3 \text{ is constant}$$

$$\textcircled{1} \quad \therefore \dot{\omega}_1 = -\frac{c\omega_3}{A} \omega_2 \quad \rightarrow$$

$$\textcircled{2} \quad \therefore (A+ma^2)\ddot{\omega}_2 - (c+ma^2)\omega_3 \dot{\omega}_1 + mga \dot{\alpha} = 0$$

$$\text{and } \omega_2 = \dot{\theta} = -\dot{\alpha}$$

$$\therefore (A+ma^2)\ddot{\omega}_2 + (c+ma^2)\omega_3 \frac{c\omega_3}{A} \omega_2 - mga \omega_2 = 0$$

$$\therefore (A+ma^2)\ddot{\omega}_2 + \left[(c+ma^2) \frac{c}{A} \omega_3^2 - mga \right] \omega_2 = 0$$

$$\text{SHM if } \omega_3^2 > \frac{mga}{(c+ma^2) \frac{c}{A}}$$

$$\begin{aligned} C &= ma^2 \\ A &= \frac{1}{2}ma^2 \end{aligned}$$

$$\therefore \omega_3^2 > \frac{g}{4a}, \quad V^2 > \frac{ga}{4}$$

 bike wheel

$$\begin{aligned} C &= \frac{1}{2}ma^2 \\ A &= \frac{1}{4}ma^2 \end{aligned}$$

$$\therefore \omega_3^2 > \frac{g}{3a}, \quad V^2 > \frac{ga}{3}$$

 disc

wobbling frequency for a wheel, $C = ma^2$
 $A = \frac{1}{2}ma^2$

$$\therefore \left(\frac{1}{2} + 1\right) \ddot{\omega}_2 + \left[(1+1)2\omega_3^2 - \frac{g}{a}\right] \omega_2 = 0$$

$$\therefore \boxed{\frac{3}{2} \ddot{\omega}_2 + \left[4\omega_3^2 - \frac{g}{a}\right] \omega_2 = 0}$$

$$U_3 = a\omega_2$$

$$U_2 = -a\omega_3$$

$$\therefore \boxed{\ddot{U}_3 + \frac{2}{3} \left[4V^2 - ga\right] U_3 = 0}$$

$$(\text{wobble frequency})^2 = \frac{2}{3} \left[4V^2 - ga\right]$$

$$\therefore \text{need } V^2 > \frac{ga}{4}$$

3 (a) (i) The question gives the necessary angular velocities but need to add $\dot{\Theta}$ to the $\underline{j}$ component

$$\Omega_1 = -\Omega \cos \lambda \sin \Theta$$

$$\Omega_2 = \Omega \sin \lambda + \dot{\Theta}$$

$$\Omega_3 = \Omega \cos \lambda \cos \Theta$$

Gyro equations :

$$Q_1 = A\dot{\Omega}_1 - (A\Omega_3 - C\omega_3)\Omega_2$$

$$Q_2 = A\dot{\Omega}_2 + (A\Omega_3 - C\omega_3)\Omega_1$$

$$Q_3 = C\dot{\omega}_3$$

Fast spin and frictionless $\therefore \omega_3 = \omega$
 $\therefore A\Omega_3 \ll C\omega_3 = \text{const}$

$$Q_2 = 0 \quad (\text{free to turn about } \underline{j})$$

$$\therefore A\dot{\Omega}_2 - C\omega_3\Omega_1 = 0$$

$$\therefore \underline{A\ddot{\Theta} + C\omega_3\Omega \cos \lambda \sin \Theta = 0}$$

(ii) Steady state, $\ddot{\Theta} = 0 \quad \therefore \sin \Theta = 0$
 $\therefore \Theta = 0 \text{ or } \pi$

$$\Theta = \text{small (around zero)} \quad \therefore A\ddot{\Theta} + C\omega_3\Omega \cos \lambda \Theta = 0$$

This is SHM and therefore stable
 when $\Theta = 0$, aligned with North

3(a)(ii) cont

Now try $\Theta = \pi + x$ (around $\Theta = \pi$, x small)

$$\sin \Theta = \sin(\pi + x) = -\sin x$$

$$\therefore A\ddot{x} - C\omega_3 \Omega \cos \lambda x = 0$$

This is unstable so $\Theta = \pi$ (south pointing)
is not stable

(iii) The non-zero couple component is Q_1

and in steady state $\Theta = 0 \quad \therefore \dot{\Omega}_1 = 0$

$$\therefore \dot{Q}_1 = 0$$

$$\begin{aligned}\therefore Q_1 &= C\omega_3 \Omega_2 \\ &= \underline{\underline{C\omega \Omega \sin \lambda}}\end{aligned}$$

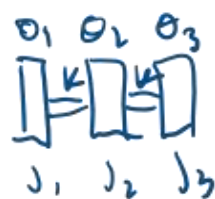
[Aside: at the equator, $\lambda = 0 \quad \therefore Q_1 = 0$
because the rotor stays aligned with
the Earth's spin

at (near) the North pole, $\lambda = 90^\circ \quad \therefore Q_1 = C\omega \Omega$

which is the classic gyro couple as in

Part IB mechanics, ie the couple
needed to keep a gyro precessing at
rate Ω and spin ω]

AC685 1 Q56



a) $q_1 = \theta_1 \quad q_2 = \theta_2 \quad q_3 = \theta_3$

$$T = \frac{1}{2} j_1 \dot{\theta}_1^2 + \frac{1}{2} j_2 \dot{\theta}_2^2 + \frac{1}{2} j_3 \dot{\theta}_3^2$$

$$T = \frac{1}{2} \underline{\dot{q}}^T \underline{M} \underline{\dot{q}}$$

$$\rightarrow \underline{M} = \begin{bmatrix} j_1 & 0 & 0 \\ 0 & j_2 & 0 \\ 0 & 0 & j_3 \end{bmatrix}$$

b) $H(\theta, p_\theta) = T + V$ given (a)

$$V = \frac{1}{2} k (\theta_1 - \theta_2)^2 + \frac{1}{2} k (\theta_2 - \theta_3)^2$$

$$p_{\theta_i} = \frac{\partial T}{\partial \dot{\theta}_i} = j_i \dot{\theta}_i \rightarrow \dot{\theta}_i = \frac{p_{\theta_i}}{j_i}$$

$$H = \frac{1}{2} \frac{p_{\theta_1}^2}{j_1} + \frac{1}{2} \frac{p_{\theta_2}^2}{j_2} + \frac{1}{2} \frac{p_{\theta_3}^2}{j_3} + \frac{1}{2} k (\theta_1 - \theta_2)^2 + \frac{1}{2} k (\theta_2 - \theta_3)^2$$

$$c) \quad p_{\text{tot}} = p_1 + p_2 + p_3 = J_1 \dot{\theta}_1 + J_2 \dot{\theta}_2 + J_3 \dot{\theta}_3$$

$$\left\{ p_{\text{tot}}, H \right\} = \underbrace{\frac{\partial p_{\text{tot}}}{\partial \theta_1}}_{=0} \frac{\partial H}{\partial p_{\theta_1}} - \underbrace{\frac{\partial p_{\text{tot}}}{\partial \theta_1}}_{=1} \underbrace{\frac{\partial H}{\partial \theta_1}}_{=k(\theta_1 - \theta_2)} +$$

$$+ \underbrace{\frac{\partial p_{\text{tot}}}{\partial \theta_2}}_{=0} \frac{\partial H}{\partial p_{\theta_2}} - \underbrace{\frac{\partial p_{\text{tot}}}{\partial \theta_2}}_{=1} \underbrace{\frac{\partial H}{\partial \theta_2}}_{=k(\theta_2 - \theta_1 - \theta_3)} +$$

$$+ \underbrace{\frac{\partial p_{\text{tot}}}{\partial \theta_3}}_{=0} \frac{\partial H}{\partial p_{\theta_3}} - \underbrace{\frac{\partial p_{\text{tot}}}{\partial \theta_3}}_{=1} \underbrace{\frac{\partial H}{\partial \theta_3}}_{=k(\theta_3 - \theta_2)}$$

$$= k(\theta_1 - \theta_2) + k(2\theta_2 - \theta_1 - \theta_3) + k(\theta_3 - \theta_2) \\ = 0$$

total momentum is conserved

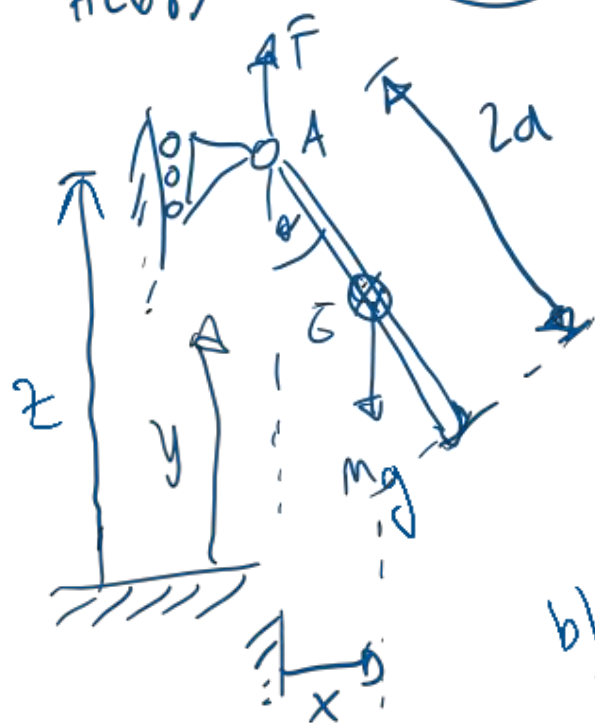
Justification: there are no external generalised torques $\rightarrow$ under an offset in angle (one in each q_i), the system is invariant.

AC685

(a4)

$$F = mg(1 + b \cos \omega t)$$

$$q_1 = z \quad q_2 = \theta$$



$$a) \quad \delta W = F \delta z$$

$$\delta W = mg(1 + b \cos \omega t) \delta z$$

$$Q_z = mg(1 + b \cos \omega t)$$

$$Q_\theta = 0$$

$$b) \quad \frac{\partial V_{\text{eff}}}{\partial z} = -Q_z \Rightarrow V_{\text{eff}} = -\int Q_z dz = -z mg(1 + b \cos \omega t)$$

$$c) \quad L(z, \theta, \dot{z}, \dot{\theta}, t) = T - V + \underbrace{z mg(1 + b \cos \omega t)}_{V_{\text{eff}}}$$

$$V = mg(z - a \cos \theta)$$

$$T = \frac{1}{2} m \dot{x}_G^2 + \frac{1}{2} m \dot{y}_G^2 + \frac{1}{2} I_G \dot{\theta}^2$$

$$\begin{bmatrix} x_G \\ y_G \end{bmatrix} = \begin{bmatrix} a \sin \theta \\ z - a \cos \theta \end{bmatrix}$$

$$\rightarrow \begin{bmatrix} \dot{x}_G \\ \dot{y}_G \end{bmatrix} = \begin{bmatrix} a \cos \theta \dot{\theta} \\ \dot{z} + a \sin \theta \dot{\theta} \end{bmatrix}$$

$$T = \frac{1}{2} m (a \cos \theta \dot{\theta})^2 + \frac{1}{2} m (\dot{z} + a \sin \theta \dot{\theta})^2 + \frac{1}{2} I_G \dot{\theta}^2$$

$$T = \frac{1}{2} \left(\frac{4M a^2}{3} \right) \dot{\theta}^2 + \frac{1}{2} m \dot{z}^2 + \frac{1}{2} 2M \dot{z} a \sin \theta \dot{\theta}$$

$$\text{with } J_{\theta} = \frac{1}{12} m (2a)^2 = \frac{ma^2}{3}$$

i) for small θ

$$\sin \theta \approx \theta$$

$$T = \frac{1}{2} \left(\frac{4M a^2}{3} \right) \dot{\theta}^2 + \frac{1}{2} m \dot{z}^2 + m \dot{z} a \theta$$

$$p_{\theta} = \frac{\partial T}{\partial \dot{\theta}} = \frac{4M a^2}{3} \dot{\theta} + m \dot{z} a \theta$$

$$p_z = \frac{\partial T}{\partial \dot{z}} = m \dot{z} + m a \dot{\theta} \theta \rightarrow = 0$$

ii) for small θ , find Hamiltonian

$$H(p_z, p_{\theta}, z, \theta, t) = p_z \dot{z} + p_{\theta} \dot{\theta} - T + V - Q_{\theta} z$$

potential energy
coupled
by
 Q_{θ}

$$\dot{z} = \frac{p_z}{m}$$

$$(p_{\theta} - m a \theta \dot{z}) = \frac{4M a^2}{3} \dot{\theta}$$

$$\Rightarrow \dot{\theta} = \frac{p_{\theta} - m a \theta \frac{p_z}{m}}{\frac{4M a^2}{3}}$$

$$\frac{4M a^2}{3}$$

$$H = \frac{p_z^2}{m} + \frac{p_\theta (p_\theta - a\dot{\theta} p_z)}{\frac{4m a^2}{3}} - \frac{1}{2} \frac{p_z^2}{m} - \frac{1}{2} \left(\frac{4m a^2}{3} \right) \frac{(p_\theta - a\dot{\theta} p_z)^2}{\left(\frac{4m a^2}{3} \right)^2}$$

$$+ m \left(\frac{p_z}{m} \right) \dot{\theta} a + mg(z - a \cos \theta)$$

$$- mg(1 + b \cos \omega t) z$$

Since

$$\frac{\partial V_z}{\partial z} = -Q_z$$

additional time varying term
caused by
additional force

$$e) \text{ (i) } H = T + V$$

Since the T can be written in terms of the mass matrix, one would be inclined to say $T + V$ is a statement of energy conservation.

No, however, the energy of ^{the system} V is not conserved because of the time varying force leading to V_{osc} .

$$\text{(ii) } \frac{d}{dt} f(\underline{p}, \underline{q}, t) = \frac{\partial f}{\partial t} + \{f, H\}$$

$$\rightarrow \frac{d}{dt} H(\underline{p}_2, \underline{p}_0, \underline{r}, \theta, t) = \frac{\partial H}{\partial t} + \{H, H\}$$

$$\{H, H\} \neq 0 \quad H \text{ is not constant}$$

$\Rightarrow$ Energy is not conserved.