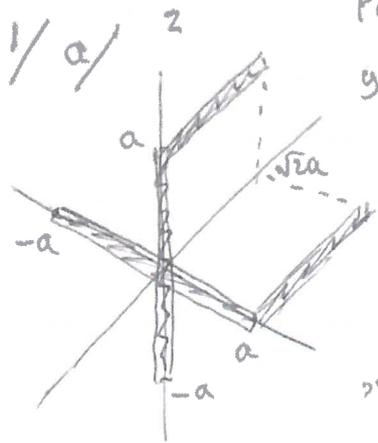


Part IIA, 3CS 2025



x-view

$$I_{xx} = \frac{1}{12} m (2a)^2 \quad (\text{CD})$$

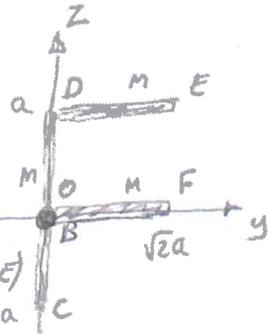
$$+ \frac{1}{12} m (a\sqrt{2})^2 + ma^2 \left(1 + \frac{\sqrt{2}}{2}\right)^2 \quad (\text{DE})$$

$$+ \frac{1}{12} m (a\sqrt{2})^2 + ma^2 \left(\frac{\sqrt{2}}{2}\right)^2 \quad (\text{BF})$$

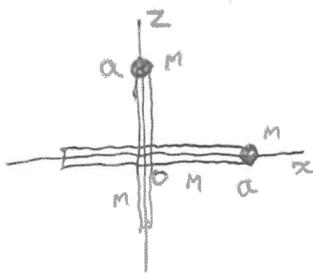
$$= ma^2 \left( \frac{1}{3} + \frac{1}{6} + 1 + \frac{1}{2} + \frac{1}{6} + \frac{1}{2} \right)$$

$$= \frac{8}{3} ma^2$$

$$I_{yz} = ma \frac{a\sqrt{2}}{2} = ma^2 \frac{\sqrt{2}}{2}$$



y-view

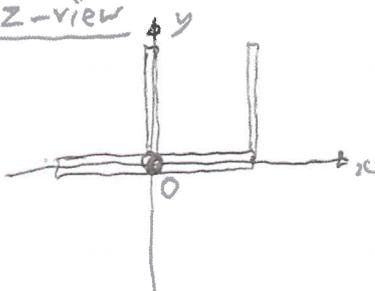


$$I_{yy} = \frac{1}{12} m (2a)^2 + \frac{1}{12} m (2a)^2 + ma^2 + ma^2$$

$$= \frac{8}{3} ma^2$$

$$I_{xz} = 0$$

z-view



Same as x-view

$$I_{zz} = \frac{8}{3} ma^2$$

$$I_{xy} = ma^2 \frac{\sqrt{2}}{2}$$

SO  $I_{xx} = I_{yy} = I_{zz} = \frac{8}{3} ma^2$

$$I_0 = ma^2 \begin{bmatrix} \frac{8}{3} & -\frac{\sqrt{2}}{2} & 0 \\ -\frac{\sqrt{2}}{2} & \frac{8}{3} & -\frac{\sqrt{2}}{2} \\ 0 & -\frac{\sqrt{2}}{2} & \frac{8}{3} \end{bmatrix}$$

N.B. this is NOT an AAA body because  $I_{yz}$  and  $I_{xy}$  are not zero. Need to use the whole inertia matrix for part (b)

(b) We're told that  $G$  is at  $\frac{1}{4}(a, \sqrt{2}a, a)$   
and that  $OG$  is a principal axis

Verify using

$[I_0] \underline{u} = \lambda \underline{u}$  where  $\underline{u}$  is the principal axis  
and  $\lambda$  is the principal moment  
of inertia

$$ma^2 \begin{bmatrix} \frac{8}{3} & -\frac{\sqrt{2}}{2} & 0 \\ -\frac{\sqrt{2}}{2} & \frac{8}{3} & -\frac{\sqrt{2}}{2} \\ 0 & -\frac{\sqrt{2}}{2} & \frac{8}{3} \end{bmatrix} \begin{bmatrix} 1 \\ \sqrt{2} \\ 1 \end{bmatrix} = ma^2 \begin{bmatrix} \frac{8}{3} - 1 \\ -\frac{\sqrt{2}}{2} + \frac{8\sqrt{2}}{3} - \frac{\sqrt{2}}{2} \\ -1 + \frac{8}{3} \end{bmatrix} = ma^2 \begin{bmatrix} \frac{5}{3} \\ \sqrt{2}(\frac{8}{3} - 1) \\ \frac{5}{3} \end{bmatrix} = \frac{5}{3} \begin{bmatrix} 1 \\ \sqrt{2} \\ 1 \end{bmatrix} ma^2$$

So yes,  $OG$  is principal and  $I_{OG}$  is  $\frac{5}{3}ma^2$

note, don't assume it's an "AAA" body  
just because  $I_{xx} = I_{yy} = I_{zz} = \frac{8}{3}ma^2$  are  
all equal. If ~~that~~ it was an AAA body  
then  $I_{OG}$  would be  $\frac{8}{3}ma^2$

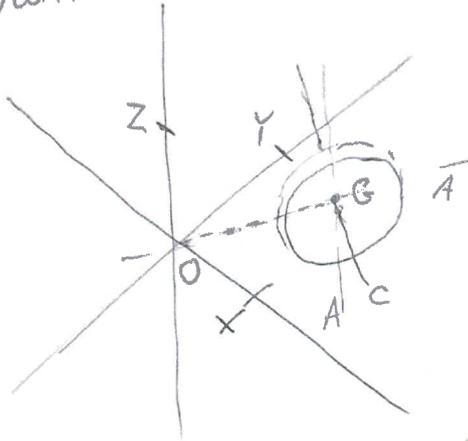
~~(c) We're told that the framework is "AAC" at  $G$  and  
but we can deduce that the value " $\frac{5}{3}ma^2$ " that  
we found in (b) <sup>is not</sup> cannot be the "C" value because  
if it were then by parallel axes theorem the body  
would be AAC~~

~~where~~ we have principal moments of inertia at  $G$

as  $\frac{5}{3}ma^2, \frac{5}{3}ma^2, \frac{8}{3}ma^2$  ~~so the~~ which is "AAC"

so the "C" value is  $\frac{8}{3}ma^2$

(C) cont



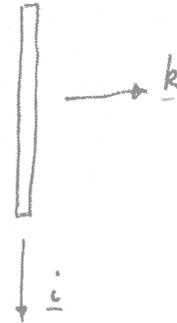
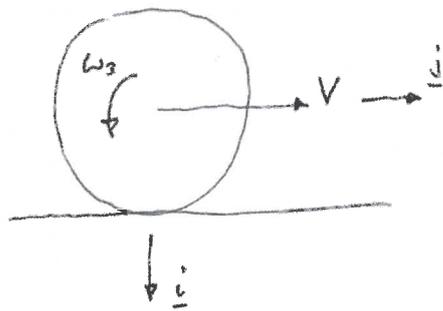
So to find the principle moments of inertia at O use parallel axes theorem on A and on C

$$\begin{aligned} \text{with } \overline{OG} &= \sqrt{x^2 + y^2 + z^2} \\ &= \sqrt{\left(\frac{a}{4}\right)^2 + \left(\frac{a\sqrt{2}}{4}\right)^2 + \left(\frac{a}{4}\right)^2} \\ &= \frac{a}{4} (1 + 2 + 1)^{\frac{1}{2}} \\ &= \frac{a}{2} \end{aligned}$$

so at O principal moments of inertia

are  ~~$\frac{5}{3}Ma^2$~~   $\frac{5}{3}Ma^2$ ,  $\frac{5}{3}Ma^2 + 4M\left(\frac{a}{2}\right)^2$ ,  $\frac{8}{3}Ma^2 + 4M\left(\frac{a}{2}\right)^2$   
 $\uparrow$   $= \frac{8}{3}Ma^2$ ,  $= \frac{11}{3}Ma^2$   
already found

2 (a)



Steady state

$$\underline{U} = V \underline{j}$$

$$\underline{\omega} = \omega_3 \underline{k} = -\frac{V}{a} \underline{k}$$

$$\underline{F} = -mg \underline{i}$$

$$\underline{Q} = 0$$

$$U_1=0 \quad U_2=V \quad U_3=0$$

$$\omega_1=0 \quad \omega_2=0 \quad \omega_3=-\frac{V}{a}$$

$$F_1=-mg \quad F_2=0 \quad F_3=0$$

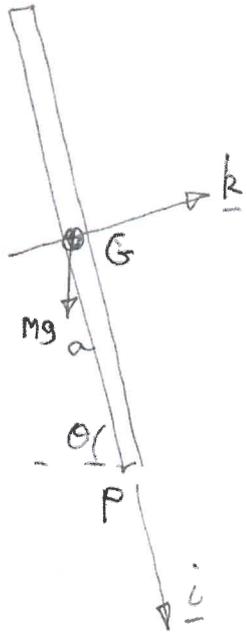
$$\Omega_3=0$$

This means that in perturbed motion ( $\theta = \frac{\pi}{2} - \alpha$ )

these quantities are small :

$$\alpha, U_1, U_3, \omega_1, \omega_2, F_2, F_3, \Omega_3$$

2 (b)(i) No slip at P



$$0 = \underline{v}_G + \underline{\omega} \times a \underline{i}$$

$$= (v_1 \underline{i} + v_2 \underline{j} + v_3 \underline{k}) + \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ \omega_1 & \omega_2 & \omega_3 \\ a & 0 & 0 \end{vmatrix}$$

$$\underline{i} : v_1 = 0$$

$$\underline{j} : v_2 + a\omega_3 = 0$$

$$\underline{k} : v_3 - a\omega_2 = 0$$

$$\therefore v_2 = -a\omega_3$$

$$\therefore v_3 = a\omega_2$$

$$\boxed{v_1 = 0}$$

$$(ii) \underline{Q}_P = \underline{h}_P + \underline{\dot{\Gamma}}_P \times \underline{P}$$

$$\underline{P} = ma(-\omega_3 \underline{j} + \omega_2 \underline{k})$$

$$\underline{\dot{\Gamma}}_P = \underline{v}_G + \underline{\Omega} \times a \underline{i}$$

$$= (v_2 + a\Omega_3) \underline{j} + (v_3 - a\Omega_2) \underline{k}$$

$$\underline{\dot{\Gamma}}_P \times \underline{P} = (ma\omega_2(v_2 + a\Omega_3) + ma\omega_3(v_3 - a\Omega_2)) \underline{i}$$

$$= ma^2 \underline{i} (-\omega_2\omega_3 + \omega_2\Omega_3 + \omega_3\omega_2 - \omega_3\Omega_2)$$

$$= ma^2 \underline{i} (\Omega_3 - \omega_3)\omega_2$$

Fast spin  $\omega_3 \gg \Omega_3 \therefore \underline{\dot{\Gamma}}_P \times \underline{P} = -ma^2 \omega_3 \omega_2 \underline{i}$

$$\underline{h}_P = \underline{h}_G + (\underline{r}_G - \underline{r}_P) \times \underline{P}$$

$$= A\omega_1 \underline{i} + A\omega_2 \underline{j} + C\omega_3 \underline{k} - (a \underline{i}) \times ma(-\omega_3 \underline{j} + \omega_2 \underline{k})$$

$$= A\omega_1 \underline{i} + (A\omega_2 + ma^2\omega_2) \underline{j} + (C\omega_3 + ma^2\omega_3) \underline{k}$$

$$= A\omega_1 \underline{i} + (A + ma^2)\omega_2 \underline{j} + (C + ma^2)\omega_3 \underline{k}$$

$$\underline{h}_P = \underline{h}_P|_{nt} + \underline{\Omega} \times \underline{h}_P$$

$$\underline{\tau} \times \underline{h}_P = \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ \omega_1 & \omega_2 & \omega_3 \\ A\omega_1 & (A+ma^2)\omega_2 & (A+ma^2)\omega_3 \end{vmatrix}$$

$$\begin{aligned} \textcircled{1} \quad \underline{Q}_P &= -mga \sin\theta \underline{j} \\ \textcircled{2} \quad \therefore \begin{cases} \underline{i}: & 0 \\ \underline{j}: & -mga \sin\theta \\ \underline{k}: & 0 \end{cases} &= \begin{cases} A\dot{\omega}_1 + (c+ma^2)\omega_2\omega_3 - (A+ma^2)\omega_2\omega_3 \\ (A+ma^2)\dot{\omega}_2 + A\omega_1\omega_3 - (c+ma^2)\omega_1\omega_3 \\ (c+ma^2)\dot{\omega}_3 + (A+ma^2)\omega_1\omega_2 - A\omega_1\omega_2 \end{cases} \end{aligned}$$

$$\begin{aligned} \textcircled{1} \quad \therefore \begin{cases} 0 &= A\dot{\omega}_1 + c\omega_2\omega_3 \\ -mga\alpha &= (A+ma^2)\dot{\omega}_2 - (c+ma^2)\omega_1\omega_3 \\ 0 &= (c+ma^2)\dot{\omega}_3 \end{cases} \quad \left( \begin{array}{l} \omega_2, \omega_3 \ll \omega_1 \\ \text{is small} \end{array} \right) \\ \textcircled{2} \quad \left( \begin{array}{l} \sin\theta = \sin\alpha \approx \alpha \\ \omega_3 \ll \omega_1 \end{array} \right) \end{aligned}$$

③  $\therefore \omega_3$  is constant

$$\textcircled{1} \quad \therefore \dot{\omega}_1 = -\frac{c\omega_3}{A}\omega_2 \quad \downarrow$$

$$\textcircled{2} \quad \therefore (A+ma^2)\ddot{\omega}_2 - (c+ma^2)\omega_3\dot{\omega}_1 + mga\alpha = 0$$

$$\text{and } \omega_2 = \dot{\theta} = -\dot{\alpha}$$

$$\therefore (A+ma^2)\ddot{\omega}_2 + (c+ma^2)\omega_3 \frac{c\omega_3}{A}\omega_2 + mga\omega_2 = 0$$

$$\therefore (A+ma^2)\ddot{\omega}_2 + \left[ (c+ma^2)\frac{c}{A}\omega_3^2 - mga \right] \omega_2 = 0$$

$$\text{SHM if } \omega_3^2 > \frac{mga}{(c+ma^2)\frac{c}{A}}$$

$$C = ma^2 \\ A = \frac{1}{2}ma^2$$

$$\therefore \omega_3^2 > \frac{g}{4a}, \quad v^2 > \frac{ga}{4} \quad \text{bike wheel}$$

$$C = \frac{1}{2}ma^2 \\ A = \frac{1}{4}ma^2$$

$$\therefore \omega_3^2 > \frac{g}{3a}, \quad v^2 > \frac{ga}{3} \quad \text{disc}$$

wobbling frequency for a wheel,  $C = ma^2$   
 $A = \frac{1}{2}ma^2$

$$\therefore \left(\frac{1}{2} + 1\right) \ddot{\omega}_2 + \left[(1+1)2\omega_3^2 - \frac{g}{a}\right] \omega_2 = 0$$

$$\therefore \boxed{\frac{3}{2} \ddot{\omega}_2 + \left[4\omega_3^2 - \frac{g}{a}\right] \omega_2 = 0}$$

$$U_3 = a\omega_2$$

$$U_2 = -a\omega_3$$

$$\therefore \boxed{\ddot{U}_3 + \frac{2}{3} \left[4V^2 - ga\right] U_3 = 0}$$

$$(\text{wobble frequency})^2 = \frac{2}{3} \left[4V^2 - ga\right]$$

$$\therefore \text{need } V^2 > \frac{ga}{4}$$

3 (a) (i) The question gives the necessary angular velocities but need to add  $\dot{\theta}$  to the  $\underline{j}$  component

$$\Omega_1 = -\Omega \cos \lambda \sin \theta$$

$$\Omega_2 = \Omega \sin \lambda + \dot{\theta}$$

$$\Omega_3 = \Omega \cos \lambda \cos \theta$$

Gyro equations :

$$Q_1 = A \dot{\Omega}_1 - (A \Omega_3 - C \omega_3) \Omega_2$$

$$Q_2 = A \dot{\Omega}_2 + (A \Omega_3 - C \omega_3) \Omega_1$$

$$Q_3 = C \dot{\omega}_3$$

Fast spins and frictionless  $\therefore \omega_3 = \omega$   
 $= \text{const}$   
 $\therefore A \Omega_3 \ll C \omega_3$

$$Q_2 = 0 \quad (\text{free to turn about } \underline{j})$$

$$\therefore A \dot{\Omega}_2 - C \omega_3 \Omega_1 = 0$$

$$\therefore \underline{A \ddot{\theta} + C \omega_3 \Omega \cos \lambda \sin \theta = 0}$$

(ii) steady state,  $\ddot{\theta} = 0 \quad \therefore \sin \theta = 0$

$$\therefore \theta = 0 \text{ or } \pi$$

$$\theta = \text{small (around zero)} \quad \therefore A \ddot{\theta} + C \omega_3 \Omega \cos \lambda \theta = 0$$

This is SHM and therefore stable  
 when  $\theta = 0$ , aligned with North

3(a)(ii) cont

Now try  $\theta = \pi + \alpha$  (around  $\theta = \pi$ ,  $\alpha$  small)

$$\sin \theta = \sin(\pi + \alpha) = -\sin \alpha$$

$$\therefore A \ddot{x} - C \omega_3 \Omega \cos \lambda x = 0$$

This is unstable so  $\theta = \pi$  (south pointing)  
is not stable

(iii) The non-zero couple component is  $Q_1$

and in steady state  $\theta = 0 \quad \therefore \dot{\theta}_1 = 0$

$$\therefore \dot{\theta}_1 = 0$$

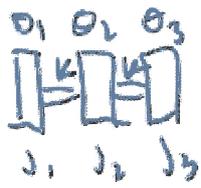
$$\begin{aligned} \therefore Q_1 &= C \omega_3 \Omega_2 \\ &= \underline{C \omega \Omega \sin \lambda} \end{aligned}$$

[Aside: at the equator,  $\lambda = 0 \quad \therefore Q_1 = 0$   
because the rotor stays aligned with  
the Earth's spin

at (near) the north pole,  $\lambda = 90^\circ \quad \therefore Q_1 = C \omega \Omega$

which is the classic gyro couple as in  
Part IB mechanics, ie the couple  
needed to keep a gyro precessing at  
rate  $\Omega$  and spin  $\omega$  ]

AC685, Q56



a)  $q_1 = \theta_1, q_2 = \theta_2, q_3 = \theta_3$

$$T = \frac{1}{2} J_1 \dot{\theta}_1^2 + \frac{1}{2} J_2 \dot{\theta}_2^2 + \frac{1}{2} J_3 \dot{\theta}_3^2$$

$$T = \frac{1}{2} \underline{\dot{q}}^T \underline{M} \underline{\dot{q}}$$

$$\rightarrow \underline{M} = \begin{bmatrix} J_1 & 0 & 0 \\ 0 & J_2 & 0 \\ 0 & 0 & J_3 \end{bmatrix}$$

b)  $H(\theta, p_\theta) = T + V$  given (a)

$$V = \frac{1}{2} k (\theta_1 - \theta_2)^2 + \frac{1}{2} k (\theta_2 - \theta_3)^2$$

$$p_{\theta_i} = \frac{\partial T}{\partial \dot{\theta}_i} = J_i \dot{\theta}_i \rightarrow \dot{\theta}_i = \frac{p_{\theta_i}}{J_i}$$

$$H = \frac{1}{2} \frac{p_{\theta_1}^2}{J_1} + \frac{1}{2} \frac{p_{\theta_2}^2}{J_2} + \frac{1}{2} \frac{p_{\theta_3}^2}{J_3} + \frac{1}{2} k (\theta_1 - \theta_2)^2 + \frac{1}{2} k (\theta_2 - \theta_3)^2$$

$$c) \quad P_{\text{TOT}} = P_1 + P_2 + P_3 = J_1 \dot{\theta}_1 + J_2 \dot{\theta}_2 + J_3 \dot{\theta}_3$$

$$\left\{ P_{\text{TOT}}, H \right\} = \underbrace{\frac{\partial P_{\text{TOT}}}{\partial \theta_1}}_{=0} \frac{\partial H}{\partial \theta_1} - \underbrace{\frac{\partial P_{\text{TOT}}}{\partial \theta_1}}_{=1} \underbrace{\frac{\partial H}{\partial \theta_1}}_{=k(\theta_1 - \theta_2)} +$$

$$+ \underbrace{\frac{\partial P_{\text{TOT}}}{\partial \theta_2}}_{=0} \frac{\partial H}{\partial \theta_2} - \underbrace{\frac{\partial P_{\text{TOT}}}{\partial \theta_2}}_{=1} \underbrace{\frac{\partial H}{\partial \theta_2}}_{=k(\theta_2 - \theta_1 - \theta_3)} +$$

$$+ \underbrace{\frac{\partial P_{\text{TOT}}}{\partial \theta_3}}_{=0} \frac{\partial H}{\partial \theta_3} - \underbrace{\frac{\partial P_{\text{TOT}}}{\partial \theta_3}}_{=1} \underbrace{\frac{\partial H}{\partial \theta_3}}_{=k(\theta_3 - \theta_2)}$$

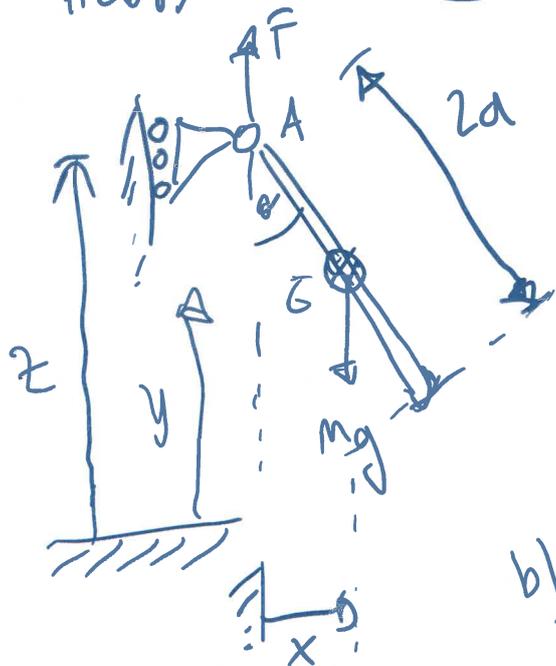
$$= k(\theta_1 - \theta_2) + k(2\theta_2 - \theta_1 - \theta_3) + k(\theta_3 - \theta_2) \\ = 0$$

total momentum is conserved

Justification: there are no external generalised torques  $\rightarrow$  under an offset in angle (one in each  $\theta_i$ ), the system is invariant.

AC685

(24)



$$F = mg(1 + b \cos \omega t)$$

$$q_1 = z \quad q_2 = \theta$$

$$a) \quad \delta W = F \delta z$$

$$\delta W = mg(1 + b \cos \omega t) \delta z$$

$$Q_z = mg(1 + b \cos \omega t)$$

$$Q_\theta = 0$$

$$b) \quad \frac{\partial V_{\text{eff}}}{\partial z} = -Q_z \Rightarrow V_{\text{eff}} = - \int Q_z dz = -z mg(1 + b \cos \omega t)$$

$$c) \quad L(z, \theta, \dot{z}, \dot{\theta}, t) = T - V + \underbrace{z mg(1 + b \cos \omega t)}_{V_{\text{eff}}}$$

$$V = mg(z - a \cos \theta)$$

$$T = \frac{1}{2} m \dot{x}_G^2 + \frac{1}{2} m \dot{y}_G^2 + \frac{1}{2} J_G \dot{\theta}^2$$

$$\begin{bmatrix} x_G \\ y_G \end{bmatrix} = \begin{bmatrix} a \sin \theta \\ z - a \cos \theta \end{bmatrix}$$

$$\rightarrow \begin{bmatrix} \dot{x}_G \\ \dot{y}_G \end{bmatrix} = \begin{bmatrix} a \cos \theta \dot{\theta} \\ \dot{z} + a \sin \theta \dot{\theta} \end{bmatrix}$$

$$T = \frac{1}{2} m (a \cos \theta \dot{\theta})^2 + \frac{1}{2} m (\dot{z} + a \sin \theta \dot{\theta})^2 + \frac{1}{2} J_G \dot{\theta}^2$$

$$T = \frac{1}{2} \left( \frac{4M}{3} a^2 \right) \dot{\theta}^2 + \frac{1}{2} m \dot{z}^2 + \frac{1}{2} \cancel{2M} \dot{z} a \sin \theta \dot{\theta}$$

$$\text{with } J_{\theta} = \frac{1}{12} m (2a)^2 = \frac{ma^2}{3}$$

d) for small  $\theta$   $\sin \theta \approx \theta$

$$T = \frac{1}{2} \left( \frac{4ma^2}{3} \right) \dot{\theta}^2 + \frac{1}{2} m \dot{z}^2 + m \dot{z} \theta a$$

$$p_{\theta} = \frac{\partial T}{\partial \dot{\theta}} = \frac{4ma^2}{3} \dot{\theta} + m \dot{z} \theta a$$

$$p_z = \frac{\partial T}{\partial \dot{z}} = m \dot{z} + \cancel{ma \theta} \dot{\theta} \rightarrow = 0$$

dii) for small  $\theta$ , find Hamiltonian

$$H(p_z, p_{\theta}, z, \theta, t) = p_z \dot{z} + p_{\theta} \dot{\theta} - T + V - \underbrace{Q_z t}_{\text{potential energy coupled by } Q_t}$$

$$\dot{z} = \frac{p_z}{m}$$

$$(p_{\theta} - ma \theta \dot{z}) = \frac{4ma^2}{3} \dot{\theta}$$

$$\Rightarrow \dot{\theta} = \frac{p_{\theta} - \cancel{ma \theta} \frac{p_z}{m}}{\frac{4ma^2}{3}}$$

$$\frac{4ma^2}{3}$$

$$H = \frac{P_z^2}{M} + \frac{P_\theta (P_\theta - a\theta P_z)}{\frac{4M a^2}{3}} - \frac{1}{2} \frac{P_z^2}{M} - \frac{1}{2} \left( \frac{4M a^2}{3} \right) \frac{(P_\theta - a\theta P_z)^2}{\left| \frac{4M a^2}{3} \right|^2}$$

$$+ M \left( \frac{P_z}{M} \right) \theta a + Mg(z - a \cos \theta)$$

$$- Mg(1 + b \cos \omega t) z$$

since

$$\frac{\partial V_z}{\partial z} = -Q_z$$

additional time varying term  
caused by additional force

$$e) \text{ (i) } H = T + V$$

Since the  $T$  can be written in terms of the mass matrix, one would be inclined to say  $T+V$  is a statement of energy conservation.

No, however, the energy of <sup>the system</sup>  $V$  is not conserved because of the time varying force leading to  $V_{var}$ .

$$\text{(ii) } \frac{d}{dt} f(p, q, t) = \frac{\partial f}{\partial t} + \{f, H\}$$

$$\rightarrow \frac{d}{dt} H(p_2, p_\theta, z, \theta, t) = \frac{\partial H}{\partial t} + \underbrace{\{H, H\}}_{=0}$$

$\neq 0 \Rightarrow H$  is not constant

$\Rightarrow$  Energy is not conserved.