## Question 1

- (a) The procedure to derive weak formulation consists of three steps
  - 1. Multiply the strong equation by a weight function which is equal to zero where Dirichlet boundary condition are applied but is otherwise arbitrary.
  - 2. Use integration by parts to 'shift' derivatives to the weight function.
  - 3. Insert the Neumann boundary condition
- (b) First note

$$-\left(\frac{du}{dx}\right)^2 - u\frac{d^2u}{dx^2} = -\frac{d}{dx}\left(u\frac{du}{dx}\right) \tag{1}$$

Following the three-step procedure, we write the weak from:

$$0 = \int_0^1 v(-\frac{d}{dx}(u\frac{du}{dx}) + f)dx \tag{2}$$

$$= \int_0^1 (v \frac{dv}{dx} \frac{du}{dx} + vf) dx - (v(u \frac{du}{dx}))|_0^1,$$
 (3)

where v is the test function. Plug in the boundary condition we obtain the weak formulation

$$0 = \int_0^1 \left( u \frac{du}{dx} \frac{dv}{dx} + vf \right) dx. \tag{4}$$

(c) The highest order of derivative in the weak formulation is 1, therefore, one can use a simple  $C^0$  element to model the problem. The approximate solution  $u_h$  given by the  $C^0$  element is continuous, but its first order derivative  $\frac{du_h}{dx}$  is not continuous in general.

$$\frac{\partial^4 v}{\partial x^4} + \frac{\partial^4 v}{\partial y^4} = -\frac{1}{p}$$

$$\int \frac{\partial^4 v}{\partial x^4} w dx = \int \frac{\partial}{\partial x} \left( \frac{\partial^3 v}{\partial x^3} w \right) dx - \int \frac{\partial^3 v}{\partial x^3} \frac{\partial w}{\partial x} dx$$

$$= \int \frac{\partial x}{\partial x} \left( \frac{\partial x^{3}}{\partial x^{3}} w \right) dx - \int \frac{\partial x}{\partial x} \left( \frac{\partial x^{2}}{\partial x^{2}} \frac{\partial x}{\partial x} \right) + \int \frac{\partial x^{2}}{\partial x^{2}} \frac{\partial x^{2}}{\partial x^{2}}$$

$$\int \left( \frac{\partial^2 u}{\partial x^2} \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \frac{\partial^2 w}{\partial y^2} \right) = - \int \frac{1}{2\pi} w \, dx$$

Boundary terms:

$$\int \left(\frac{\partial^3 v}{\partial x^3} w + \frac{\partial^3 v}{\partial y^3} w\right) n \, dx \qquad n-boundary normal$$

$$\left( \frac{\partial^2 v}{\partial x^2} \frac{\partial w}{\partial x} + \frac{\partial^2 v}{\partial y^2} \frac{\partial w}{\partial x} \right) n \, dx$$

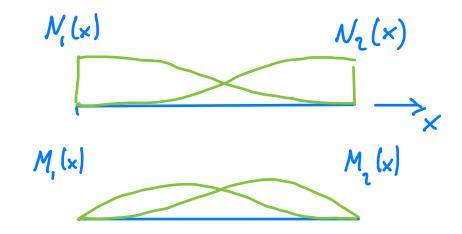
b) Each boundary has two different boundary conditions because fourth order differential equation

Displacements fixed: w = 0Sheur force prescribed:  $\frac{\partial^3 w}{\partial x^3} \neq 0$ ,  $\frac{\partial^3 w}{\partial y^3} \neq 0$ 

Slope fixed/clamped:  $\frac{\partial w}{\partial x} = 0$ 

Moments prescribed:  $\frac{\partial^2 w}{\partial x^2} \neq 0$ ,  $\frac{\partial^2 w}{\partial y^2} \neq 0$ 

c) Tensor products of Hermite shape functions



 $N_{i}(x,y) = N_{i}(x)N_{i}(y)$   $M_{i}$ 

$$N_2(x,y) = N_2(x) N_1(y)$$

$$N_3(x,y) = N_2(x) N_2(y)$$

$$N_{\downarrow} (x,y) = N_{\downarrow}(x) N_{2}(y)$$

 $M_1(x,y) = M_1(x)M_1(y)$ 

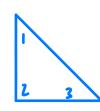
$$M_2(x,y) = M_2(x) M_i(y)$$

$$M_3(x,y) = M_2(x) M_2(y)$$

$$M_{\xi}(x_{1}y) = M_{\xi}(x) M_{\xi}(y)$$

d) We can use isoparametric mapping to obtain shape functions. However, we need to make sure that the obtained shape functions are indeed slope continuous at nodes where several elements meet.

3a)



$$\frac{1}{2} = \int N(x)^{T} {0 \choose -1} dx$$

$$= \frac{2}{3} (0 - 1 0 - 1)^{T}$$

(Recall the volume of a tetrahedron)

$$\begin{pmatrix} 0 \\ 4 \cdot \left(-\frac{2}{3}\right) \\ 0 \\ 4 \cdot \left(-\frac{2}{3}\right) \end{pmatrix} = \begin{pmatrix} 0 \\ -\frac{8}{3} \\ 0 \\ -\frac{8}{3} \end{pmatrix}$$

$$D = \begin{pmatrix} 200 & 0 & 0 \\ 0 & 200 & 0 \\ 0 & 0 & 100 \end{pmatrix}$$

$$\frac{\partial N_1}{\partial x} = 0 \qquad \frac{\partial N_2}{\partial y} = \frac{1}{2} \qquad \frac{\partial N_3}{\partial x} = \frac{1}{2} \qquad \frac{\partial N_3}{\partial y} = 0$$

$$\frac{\partial N_2}{\partial x} = -\frac{1}{2} \qquad \frac{\partial N_2}{\partial y} = -\frac{1}{2}$$

$$B = \begin{pmatrix} 0 & 0 & -\frac{1}{2} & 0 & \frac{1}{2} & 0 \\ 0 & \frac{1}{2} & 0 & -\frac{1}{2} & 0 & 0 \\ \frac{1}{2} & 0 & -\frac{1}{2} & -\frac{1}{2} & 0 & \frac{1}{2} \end{pmatrix}$$

## Question 4

(a) Let  $[U_1, U_2]$  be the nodel displacement of the element, note from the boundary condition  $U_1 = 0$ . The semi-discretized form is given by:

$$\frac{EA}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ U_2 \end{bmatrix} + \frac{mAL}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \begin{bmatrix} 0 \\ \ddot{U}_2 \end{bmatrix} = \begin{bmatrix} 0 \\ P \end{bmatrix}$$
 (5)

which simplifies to

$$\ddot{U}_2 + \frac{3E}{mL^2}U_2 = \frac{3P_0}{mAL} \tag{6}$$

with  $U_2(0) = 0$ ,  $\dot{U}_2(0) = 0$ .

(b) Assume the solution of 6 is of the form

$$U_2(t) = A\cos(\alpha t) + B\sin(\alpha t) + C \tag{7}$$

, with  $\alpha = \sqrt{\frac{3E}{mL^2}}$ .

From the initial condition we have:

$$U_2(t=0) = A\cos(0) + B\sin(0) + C = 0$$
(8)

$$\dot{U}_2(t=0) = -A\alpha\sin(0) + B\alpha\cos(0) = 0 \tag{9}$$

Where we conclude B = 0 and A = -C.

Therefore  $U_2$  simplifies to

$$U_2(t) = -C\cos(\alpha t) + C \tag{10}$$

To obtain the value of C we substitute the above equation to the governing equation:

$$-\alpha^2 C\cos(\alpha t) + \frac{3E}{mL^2}(-C\cos(\alpha t) + C) = \frac{3P_0}{mAL}$$
(11)

Note  $\alpha^2 = \frac{3E}{mL^2}$ , then,

$$\frac{3E}{mL^2}(C) = \frac{3P_0}{mAL} \tag{12}$$

Therefore

$$C = \frac{P_0 L}{EA} \tag{13}$$

The displacement field in the bar is

$$u(x,t) = U_1(-\frac{1}{L}x+1) + U_2\frac{x}{L}.$$
(14)

With  $U_1 = 0$ , we conclude

$$u(x,t) = U_2 \frac{x}{L} = \frac{P_0}{EA} (1 - \cos(\sqrt{\frac{3E}{mL^2}}t)) \cdot x.$$
 (15)