

Question 1

(a) Multiply by a test function: Let $v(x)$ be a test function satisfying $v(0) = v(1) = 0$. Multiply the equation by $v(x)$ and integrate:

$$\int_0^1 \left(\frac{d^4 w}{dx^4} + w \right) v \, dx = \int_0^1 v \, dx.$$

Apply integration by parts twice to the fourth-order term. Boundary terms vanish due to $v(0) = v(1) = 0$ and $w''(0) = w''(1) = 0$:

$$\int_0^1 \frac{d^4 w}{dx^4} v \, dx = \int_0^1 \frac{d^2 w}{dx^2} \frac{d^2 v}{dx^2} \, dx.$$

Substitute back to obtain:

$$\int_0^1 \frac{d^2 w}{dx^2} \frac{d^2 v}{dx^2} \, dx + \int_0^1 w v \, dx = \int_0^1 v \, dx \quad \forall v \in H_0^2(0, 1).$$

(b) First we check the boundary conditions with the trial solution:

1. $w(0) = 0 \implies C_3 = 0$.
2. $w(1) = 0 \implies C_3 = 0$.
3. $\frac{d^2 w}{dx^2}(0) = 0$ and $\frac{d^2 w}{dx^2}(1) = 0$ are automatically satisfied by the trial solution.

We now take a test function $v(x) = \sin(\pi x)$:

1. Compute derivatives:

$$w''(x) = -\pi^2 C_1 \sin(\pi x) - 9\pi^2 C_2 \sin(3\pi x), \quad v''(x) = -\pi^2 \sin(\pi x).$$

2. First term in weak formulation:

$$\int_0^1 w'' v'' \, dx = \pi^4 C_1 \int_0^1 \sin^2(\pi x) \, dx = \frac{\pi^4 C_1}{2}.$$

3. Second term in weak formulation:

$$\int_0^1 w v \, dx = C_1 \int_0^1 \sin^2(\pi x) \, dx = \frac{C_1}{2}.$$

4. Right-hand side:

$$\int_0^1 v \, dx = \frac{2}{\pi}.$$

5. Combine terms:

$$\frac{\pi^4 C_1}{2} + \frac{C_1}{2} = \frac{2}{\pi} \implies \boxed{C_1 = \frac{4}{\pi(\pi^4 + 1)}}.$$

We now take another test function $v(x) = \sin(3\pi x)$

1. Compute derivatives:

$$w''(x) = -\pi^2 C_1 \sin(\pi x) - 9\pi^2 C_2 \sin(3\pi x), \quad v''(x) = -9\pi^2 \sin(3\pi x).$$

2. First term in weak formulation:

$$\int_0^1 w'' v'' dx = 81\pi^4 C_2 \int_0^1 \sin^2(3\pi x) dx = \frac{81\pi^4 C_2}{2}.$$

3. Second term in weak formulation:

$$\int_0^1 w v dx = C_2 \int_0^1 \sin^2(3\pi x) dx = \frac{C_2}{2}.$$

4. Right-hand side:

$$\int_0^1 v dx = \frac{2}{3\pi}.$$

5. Combine terms:

$$\frac{81\pi^4 C_2}{2} + \frac{C_2}{2} = \frac{2}{3\pi} \implies \boxed{C_2 = \frac{4}{3\pi(81\pi^4 + 1)}}.$$

Therefore,

$$\boxed{C_1 = \frac{4}{\pi(\pi^4 + 1)}, \quad C_2 = \frac{4}{3\pi(81\pi^4 + 1)}, \quad C_3 = 0.}$$

- a) - Insufficient boundary conditions (possible rigid body deformations)
- Too few quadrature points
 - Inverted elements
 - Something wrong with material (e.g. zero Young's mod)
- b) In FE equilibrium satisfied only weakly, in the weak form we use only a limited set of functions as test and trial functions.
- c) Gaussian quadrature is very efficient for polynomials and almost polynomials. Needs far fewer evaluation points than other schemes.

$$d) \ i) \ x = \sum_{i=1}^4 N_i(\xi, \eta) x_i = N_2(\xi, \eta) + \frac{3}{4} N_3(\xi, \eta)$$

$$= \frac{1}{16} (7 - \eta + 7\xi - \eta\xi)$$

$$y = \frac{3}{4} N_3(\xi, \eta) + N_4(\xi, \eta)$$

$$= \frac{1}{16} (7 + 7\eta - \xi - \eta\xi)$$

$$ii) \quad J = \begin{pmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} \\ \frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta} \end{pmatrix} = \frac{1}{16} \begin{pmatrix} 7-\eta & -1-\eta \\ -1-\xi & 7-\xi \end{pmatrix} \quad 4$$

$$iii) \quad x(0,0) = \frac{7}{16} \quad y(0,0) = \frac{7}{16}$$

$$s(0,0) = e^{\frac{49}{256}}$$

$$J(0,0) = \frac{1}{16} \begin{pmatrix} 7 & -1 \\ -1 & 7 \end{pmatrix} \quad \det J(0,0) = \frac{3}{16}$$

$$f = 4 \cdot \frac{3}{16} e^{\frac{49}{256}} \frac{1}{4} \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \end{pmatrix} = \frac{3}{16} e^{\frac{49}{256}} \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \end{pmatrix}$$

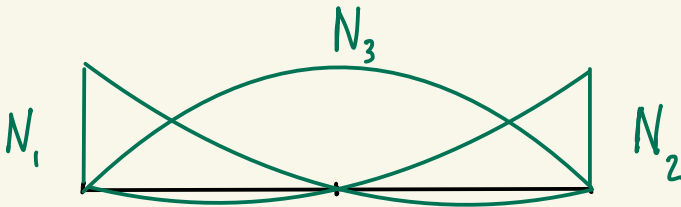
Obviously, using a single quadrature point gives only a very rough approx. Need more Gauss points, possibly 2×2 .

Q3) a)

$$N_1(\xi) = 1/2 \xi(\xi-1)$$

$$N_2(\xi) = 1/2 \xi(\xi+1)$$

$$N_3(\xi) = -(\xi-1)(\xi+1)$$



$$b) \quad N_3(\xi, \eta) = \frac{1}{4} \xi \eta (\xi+1)(\eta+1)$$

$$N_9(\xi, \eta) = (\xi-1)(\xi+1)(\eta-1)(\eta+1)$$

$$N_8(\xi, \eta) = -\frac{1}{2}(\eta-1)(\eta+1) \xi(\xi-1)$$

c) Curved boundaries and higher order convergence for smooth problems.

d) i) Weak form

$$\int (\underline{\alpha} \nabla v - \nabla \cdot (\beta \nabla v) + f) w \, d\Omega = 0$$

$$\beta \int \nabla \cdot \nabla v w d\Omega = \beta \int \nabla \cdot (\nabla v w) d\Omega - \beta \int \nabla v \cdot \nabla w d\Omega \quad 6$$

$$\int \nabla \cdot (\nabla v w) d\Omega = \int_{\Gamma_N} w \nabla v \cdot \underline{n} d\Gamma + \int_{\Gamma_D} w \nabla v \cdot \underline{n} d\Gamma$$

Choose $w=0$ on Γ_D

$$\Rightarrow \int (\underline{\alpha} \cdot \nabla v w - \beta \nabla v \cdot \nabla w) d\Omega + \int f w d\Omega + \int_{\Gamma_N} w \nabla v \cdot \underline{n} d\Gamma = 0$$

$$ii) \quad v = \sum_{i=1}^3 N_i(x, y) v_i \quad w = \sum_{i=1}^3 N_i(x, y) w_i$$

$$N_1(x, y) = 1 - x - y$$

$$N_2(x, y) = x$$

$$N_3(x, y) = y$$

$$\nabla v = \begin{pmatrix} -v_1 + v_2 \\ -v_1 + v_3 \end{pmatrix} = \begin{pmatrix} -1 & 1 & 0 \\ -1 & 0 & 1 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \\ v_3 \end{pmatrix} = \underline{B} \underline{v}$$

Stiffness corresponding to $\beta \nabla v \cdot \nabla w$

$$\frac{1}{2} \beta \underline{B}^T \underline{B} = \frac{\beta}{2} \begin{pmatrix} 2 & -1 & -1 \\ -1 & 1 & 0 \\ -1 & 0 & 1 \end{pmatrix}$$

Stiffness corresponding to $\alpha \cdot \nabla u \cdot \nabla w$

$$\alpha \cdot \nabla u = \alpha_x (-u_1 + u_2) + \alpha_y (-u_1 + u_3)$$

$$\begin{pmatrix} \alpha_x - \alpha_y \\ \alpha_x \\ \alpha_y \end{pmatrix} \cdot \int \begin{pmatrix} 1-x-y \\ x \\ y \end{pmatrix}^T dx dy$$

Note that

$$\int_0^1 \int_0^{1-x} (1-x-y) dy dx = \frac{1}{6} \quad \int_0^1 \int_0^{1-x} x dy dx = \frac{1}{6}$$

\Rightarrow

$$\frac{1}{6} \begin{pmatrix} \alpha_x - \alpha_y & \alpha_x - \alpha_y & \alpha_x - \alpha_y \\ \alpha_x & \alpha_x & \alpha_x \\ \alpha_y & \alpha_y & \alpha_y \end{pmatrix}$$

Question 4

(a)

The semi-discrete finite element system is:

$$\mathbf{M}\ddot{\mathbf{a}} + \mathbf{K}\mathbf{a} = \mathbf{f}.$$

Using the average acceleration scheme:

$$\begin{aligned}\mathbf{a}_{n+1} &= \mathbf{a}_n + \Delta t \dot{\mathbf{a}}_n + \frac{\Delta t^2}{4} (\ddot{\mathbf{a}}_n + \ddot{\mathbf{a}}_{n+1}), \\ \dot{\mathbf{a}}_{n+1} &= \dot{\mathbf{a}}_n + \frac{\Delta t}{2} (\ddot{\mathbf{a}}_n + \ddot{\mathbf{a}}_{n+1}).\end{aligned}$$

Substitute the equation of motion $\ddot{\mathbf{a}}_{n+1} = \mathbf{M}^{-1}(\mathbf{f}_{n+1} - \mathbf{K}\mathbf{a}_{n+1})$ into the displacement update:

$$\mathbf{a}_{n+1} = \mathbf{a}_n + \Delta t \dot{\mathbf{a}}_n + \frac{\Delta t^2}{4} (\ddot{\mathbf{a}}_n + \mathbf{M}^{-1}(\mathbf{f}_{n+1} - \mathbf{K}\mathbf{a}_{n+1})).$$

Rearrange terms and multiply through by \mathbf{M} :

$$\left(\mathbf{M} + \frac{\Delta t^2}{4}\mathbf{K}\right)\mathbf{a}_{n+1} = \mathbf{M}\mathbf{a}_n + \Delta t\mathbf{M}\dot{\mathbf{a}}_n + \frac{\Delta t^2}{4}\mathbf{M}\ddot{\mathbf{a}}_n + \frac{\Delta t^2}{4}\mathbf{f}_{n+1}.$$

Substitute $\ddot{\mathbf{a}}_n = \mathbf{M}^{-1}(\mathbf{f}_n - \mathbf{K}\mathbf{a}_n)$:

$$\boxed{\left(\mathbf{M} + \frac{\Delta t^2}{4}\mathbf{K}\right)\mathbf{a}_{n+1} = \left(\mathbf{M} - \frac{\Delta t^2}{4}\mathbf{K}\right)\mathbf{a}_n + \Delta t\mathbf{M}\dot{\mathbf{a}}_n + \frac{\Delta t^2}{4}(\mathbf{f}_n + \mathbf{f}_{n+1}).}$$

(b)

This is an implicit method, as it requires solving a system with $\mathbf{M} + \frac{\Delta t^2}{4}\mathbf{K}$. It is unconditionally stable for linear systems, with second order accuracy ($O(\Delta t^2)$).

(c)

A lumped mass matrix is used because its inversion is trivial and hence can significantly enhance the efficiency in explicit/semi-explicit methods. Lumped mass matrix can be constructed by row-sum lumping: $M_{ii}^{\text{lumped}} = \sum_j M_{ij}$, from the consistent mass matrix M .