

(Q1)

(a) Reactive heat flux = conductive flux

$$\dot{w} \delta_f \Delta H_c = \lambda \frac{(T_b - T_u)}{\delta_f} \quad \Delta H_c - \text{heat of combustion}$$

$$\dot{w} \delta_f c_p (T_b - T_u) = \lambda \frac{(T_b - T_u)}{\delta_f}$$

$$\Rightarrow \boxed{\delta_f^2 = \frac{\lambda}{c_p \dot{w}}}$$

Convective mass flux = Reactive flux

$$\rho_u S_L = \dot{w} \delta_f$$

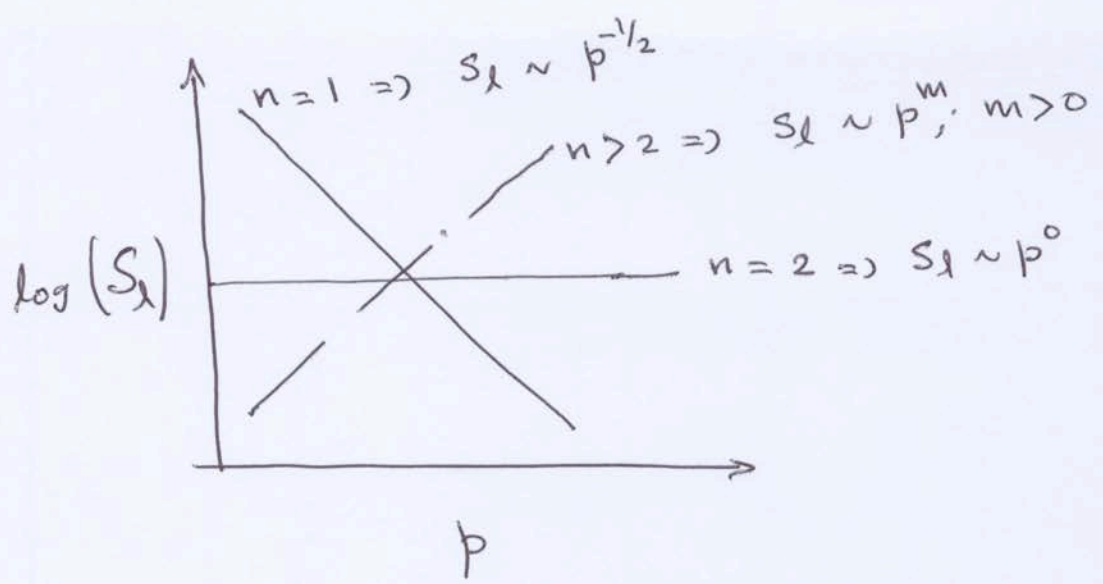
$$\Rightarrow \boxed{S_L = \frac{1}{\rho_u} \sqrt{\frac{\lambda \dot{w}}{c_p}}}$$

$$(b) \dot{w} = P^n Y_f A \exp\left(-\frac{T_a}{T_b}\right)$$

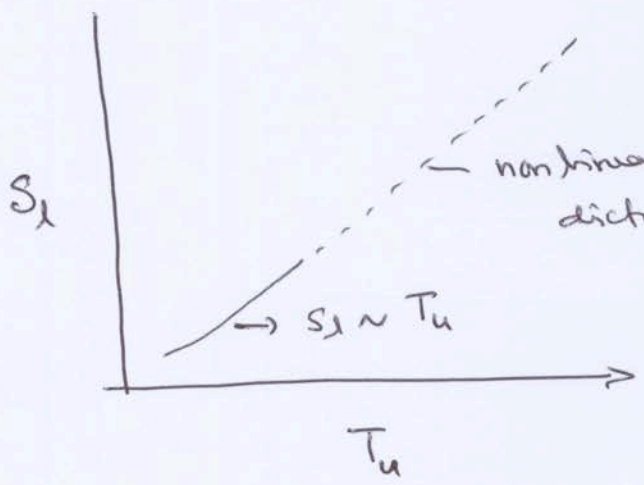
$$\Rightarrow S_L = \frac{1}{\rho_u} \sqrt{\frac{\lambda}{c_p} A Y_f} P^{n/2} \exp\left(-\frac{T_a}{2T_b}\right)$$

$$\text{but } \rho_u = \frac{P}{RT_u}$$

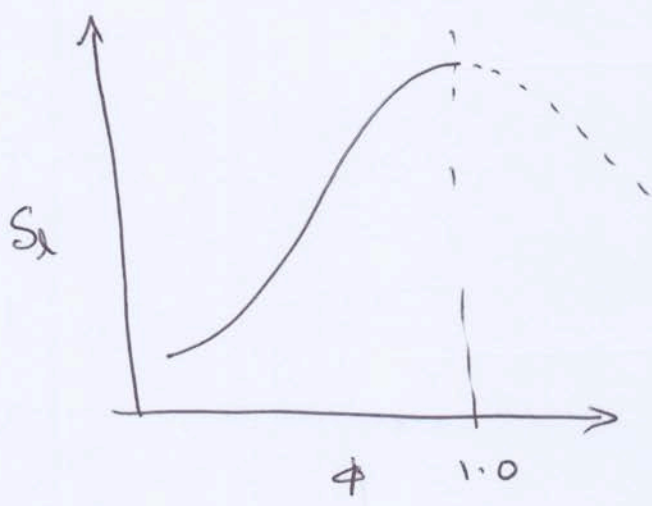
$$\Rightarrow S_L = \sqrt{\frac{\lambda}{c_p} A Y_f R^2} T_u P^{\frac{n-2}{2}} \exp\left(-\frac{T_a}{2T_b}\right)$$



$S_1$  vs  $P$  variation depends on the overall order of the reaction. The available experimental data suggests that  $S_1 \sim P^{-1/2}$ , a drop in  $S_1$  with  $P$  over a typical range of interest.



non-linear behaviour dictated by the exponential term as  $T_u \uparrow$   $T_b \uparrow$  so the exponential term dominates.



As  $\phi \uparrow$   $Y_f \uparrow$  but  $T_b$  will also increase so, the final behaviour is dictated by the compounded effect of  $Y_f$  &  $T_b$ . Also  $A$  can change with  $\phi$ .

for the rich side ( $\phi > 1$ ), the  $w$  expression doesn't fully represent the chemistry as it is good for lean mixtures.

(c) mass balance across the flame

$$\rho_u S_L = \rho_b \bar{P}_B u_b$$

if  $P = \text{const}$  & molecular wt of the mixture is also constant then

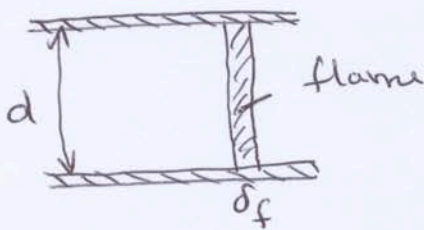
$$\Rightarrow u_b = \frac{T_b}{T_u} S_L$$

$$\frac{\rho_u}{\rho_b} = \frac{T_b}{T_u} \text{ through}$$

$$u_b - S_L = S_L \left( \frac{T_b - T_u}{T_u} \right) = \tau S_L$$

$$\therefore \boxed{u_b = (1 + \tau) S_L} \text{ as required.}$$

(d)



heat generated in the flame }  $\approx$  Conducted away through the wall

$$\dot{w} c_p (T_b - T_u) \frac{\pi d^2}{4} \delta_f$$

$$\approx \lambda \pi d \delta_f \frac{(T_b - T_w)^2}{d}$$

using  $\dot{w} = \frac{\rho_u S_L}{\delta_f}$  and rearranging

$$\Rightarrow d \approx \left\{ \frac{8 \lambda (T_b - T_w) \delta_f}{\rho_u c_p S_L (T_b - T_u)} \right\}^{1/2}$$

for the flame to be quenched (or when it is quenched)  
 $T_w = T_u$

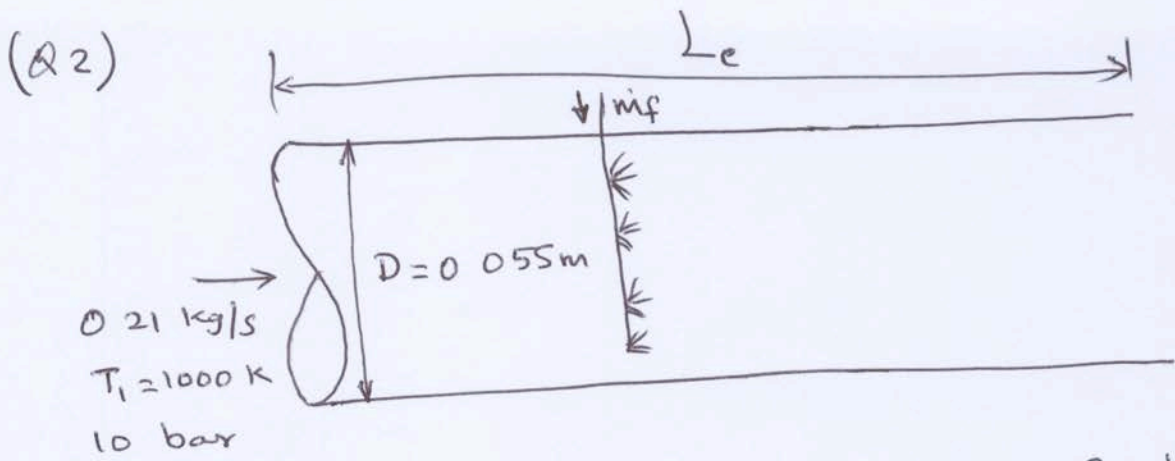
$$\Rightarrow d \leq 2\sqrt{2} \left\{ \frac{\lambda}{\rho_u c_p} \frac{\delta_f}{S_L} \right\}^{1/2}$$

but  $\frac{\lambda}{\rho_u c_p} = \alpha$  thermal diffusivity

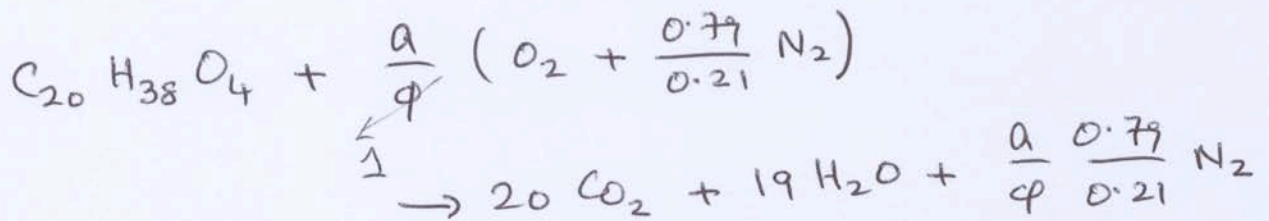
$$\& \alpha \sim \delta_f S_L$$

$$\Rightarrow \boxed{d \leq 2\sqrt{2} \delta_f}$$

in the limit  $d \approx 2\sqrt{2} \delta_f$  for the flame to be quenched



$$(a) \phi = \frac{(m_f/m_a)_{\text{actual}}}{(m_f/m_a)_{\text{st.}}} = \frac{(0.01/0.21)}{(\quad)_{\text{st.}}}$$



O atom balance:

$$4 + 2a = 40 + 19 \Rightarrow$$

$$\boxed{a = \frac{55}{2}}$$

$$\therefore \left( \frac{m_f}{m_a} \right)_{\text{st.}} = \frac{342}{\frac{55}{2} \left( 32 + \frac{0.79}{0.21} 28 \right)} = 0.0906$$

$$\therefore \phi = \frac{(0.01/0.21)}{0.0906} = 0.526 //$$

$$\boxed{\phi = 0.526}$$

(b)



fuel droplet

$$d = 30 \times 10^{-6} \text{ m}$$

(5)

$$m = \rho_f \frac{4}{3} \pi \left(\frac{d_d}{2}\right)^3$$

$$\frac{dm}{dt} = -\dot{m}'' A = -\rho_f \frac{B}{4d_d} 4\pi \frac{d_d^2}{4}$$

$$\rho_f \frac{4}{3} \pi \left(\frac{d_d}{2}\right)^2 \frac{3}{2} \frac{dd_d}{dt} = -\rho_f \frac{B}{4d_d} 4\pi \frac{d_d^2}{4}$$

$$\Rightarrow \frac{dd_d}{dt} = -\frac{B}{2d_d}$$

$$\int_0^t \Rightarrow d_d^2(t) = d_d^2(t=0) - Bt$$

$$d_d(t) = 0 \text{ @ } t_{\text{evap}} \Rightarrow t_{\text{evap}} = \frac{d_d^2(t=0)}{B}$$

$$t_{\text{evap}} = \frac{(30 \times 10^{-6})^2}{1.5 \times 10^{-7}} = 6 \times 10^{-3} \text{ s}$$

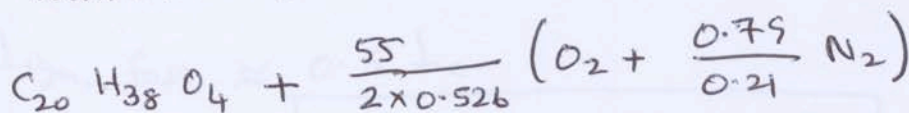
$$t_{\text{evap}} = 6 \text{ ms}$$

$$L_v = U t_{\text{evap}}$$

$$\dot{m} = \rho U A$$

$$A = \frac{\pi}{4} \times (0.055)^2 = 2.3758 \times 10^{-3} \text{ m}^2$$

The mixture is



$$\Rightarrow x_f = \frac{1}{(1 + 52.281 + 52.281 \times \frac{0.75}{0.21})} = 0.004$$

$$x_{\text{O}_2} = 0.2092; \quad x_{\text{N}_2} = 0.7868$$

$$\Rightarrow MW_{\text{mix}} = 30.092$$

$$\rho = \frac{PMW_{mix}}{RT_1} = \frac{10^6 \times 30.092}{8314 \times 1000} = 3.619 \text{ Kg/m}^3 \quad (6)$$

$$U = \frac{(0.21 + 0.01)}{3.619 \times 2.3758 \times 10^{-3}} = 25.58 \text{ m/s.}$$

$$\therefore L_v = 25.58 \times 6 \times 10^{-3} = 15.35 \text{ cm.}$$

$$L_v = 15.35 \text{ cm.}$$

$$\text{also } \underline{L_v < L_c.}$$

$$(c) \quad L_{ign} = U \tau_{ign} \quad \tau_{ign} \sim e^{-T_a/T_1}$$

$$\frac{L_{ign, bio}}{L_{ign, fossil}} = \frac{\tau_{ign, bio}}{\tau_{ign, fossil}} = \frac{e^{-1.05T_a/T_1}}{e^{-T_a/T_1}}$$

$$= e^{(0.05T_a/T_1)} = 2.117$$

$$\Rightarrow L_{ign, bio} = 2.117 L_{ign, fossil}$$

$$L_{ign, fossil} \approx 0.5 L_c$$

$$\Rightarrow \boxed{L_{ign, bio} \approx 1.059 L_c.}$$

$\Rightarrow$  Auto ignition occurs outside the combustor, which is unacceptable. So combustor length for the biofuel use can be increased, which is impractical.  $T_1$  can be increased.

$$(d) \frac{L_{ign, bio}}{L_{ign, foss}} = \frac{e^{-1.05 T_a / T_{i, bio}}}{e^{-T_a / T_{i, foss}}} = 1$$

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$$\Rightarrow T_{i, bio} = 1.05 T_{i, foss}$$

This implies that  $T_i$  has to be increased to 1050 K to allow safe operation using the biofuel

This increase of 50 K in the mixture temperature can be achieved either by

- (i) pre heating the fuel or
- (ii) increasing the compression ratio.

Q3 (a)  $\text{NH}_3$  / air flames have much lower flame speed than hydrocarbon-fuelled flames. Hence both engines & gas turbines must be adapted, or the fuel must be strengthened with a "promoter", such as a little  $\text{CH}_4$  or  $\text{H}_2$ . The  $\text{H}_2$  could come from the  $\text{NH}_3$  itself, so some transformation of  $\text{NH}_3$  into  $\text{N}_2$  &  $\text{H}_2$  (i.e. chemical cracking) is needed. This adds complexity & cost.

$\text{NH}_3$  flames are weak, hence the possibility of "NH<sub>3</sub>-slip" is higher.  $\text{NH}_3$  is toxic.

Also, there may be strong emission of nitrogen oxides ( $\text{NO}$ ,  $\text{NO}_2$ ,  $\text{N}_2\text{O}$ ), especially  $\text{N}_2\text{O}$  may be large.  $\text{N}_2\text{O}$  is a potent greenhouse gas, which means some of the key advantage of  $\text{NH}_3$  (i.e. the absence of C, hence no  $\text{CO}_2$ ) is negated.

$\text{NH}_3$  gas turbines are under development at present. They have the potential to reduce  $\text{CO}_2$  in our energy system.



(b) The stability of a combustor depends, approximately, on the Damköhler number. This can be defined as

$$Da = \frac{\text{physical time}}{\text{chemical time}} = \frac{L/u}{\nu/S_L^2}$$

$L$ : lengthscale of combustor  
 $u$ : inlet velocity  
 $\nu$ : kinematic viscosity  
 $S_L$ : laminar flame speed

If  $Da$  is to stay the same as  $P_0$  increases to  $2P_0$ , and since  $L$  is fixed, and since residence time ( $=L/u$ ) stays the same, this means  $\nu/S_L^2$  must remain constant, from condition 0 ( $P_0, T_0$ ) to condition 1 ( $2P_0, T_1$ ).

$$\nu = \frac{\mu}{\rho} \quad \mu \sim \mu_0 \left(\frac{T_0}{T}\right)^\alpha \quad (\alpha \approx 0.7)$$

$$P = \rho RT \Rightarrow P_0 = \rho_0 RT \quad 2$$

$$P_1 = \rho_1 RT_1 \Rightarrow 2P_0 = \rho_1 RT_1$$

$$\Rightarrow \rho_1 = \frac{2P_0}{RT_1} \Rightarrow \nu_1 = \nu_0 \cdot 2 \left(\frac{T_0}{T_1}\right)^\alpha \left(\frac{T_0}{T_1}\right)^{-1} = 2\rho_0 \frac{T_0}{T_1}$$

$$\Rightarrow S_L = S_{L0} \left(\frac{2P_0}{P_0}\right)^m \left(\frac{T_1}{T_0}\right)^n \Rightarrow S_L^2 = S_{L0}^2 (2)^{2m} \left(\frac{T_1}{T_0}\right)^{2n}$$

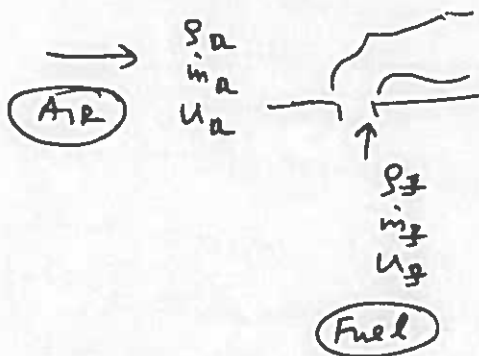
$\Rightarrow$  If  $\frac{\nu}{S_L^2}$  stays constant, we must have:

$$2 \cdot (2)^{-2m} \cdot \left(\frac{T_0}{T_1}\right)^{-1-\alpha+2n} = 1 \quad (\text{or}) \quad 2^{(1-2m)} \left(\frac{T_1}{T_0}\right)^{1+\alpha-2n} = 1$$

Depending on  $\alpha, m, n$ , we get  $T_1/T_0$ .

This problem demonstrates the complexity of scaling-up combustion problems.

(c) The jet cross flow is determined by the mass & momentum ratios between the two streams



$$\dot{m}_f = N \cdot \left(\frac{\pi d^2}{4}\right) \rho_f u_f$$

$$J_f = N \left(\frac{\pi d^2}{4}\right) \rho_f u_f^2$$

$$\Rightarrow J_f = \dot{m}_f u_f$$

We are told that when we switch from fuel 1 to fuel 2, power must stay the same, i.e.  $\dot{m}_{f,1} Q_1 = \dot{m}_{f,2} Q_2$  ( $Q$ : calorific value)

$$Q_{CH_4} = 50 \text{ MJ/kg} \quad (Q_1)$$

$$Q_{H_2} = 120 \text{ MJ/kg} \quad (Q_2)$$

$$\dot{m}_{a,1} = \dot{m}_{a,2} \quad \& \quad J_{a,1} = J_{a,2} \quad (u_{a,1} = u_{a,2})$$

Since momentum flow rate ratio must also stay the same, then  $(\dot{m}_{f,1} u_1) = (\dot{m}_{f,2} u_2)$

$$\Rightarrow u_{f,2} = u_{f,1} \frac{Q_2}{Q_1}$$

Since  $J_f$  stays constant,

$$N_1 \rho_1 \frac{\pi d_1^2}{4} u_1^2 = N_2 \rho_2 \frac{\pi d_2^2}{4} u_2^2$$

$$d_1^2 N_1 = \left(\frac{Q_2}{Q_1}\right)^2 d_2^2 N_2 \Rightarrow \boxed{\frac{d_2}{d_1} = \sqrt{\frac{N_1 \rho_1 Q_2}{N_2 \rho_2 Q_1}}}$$

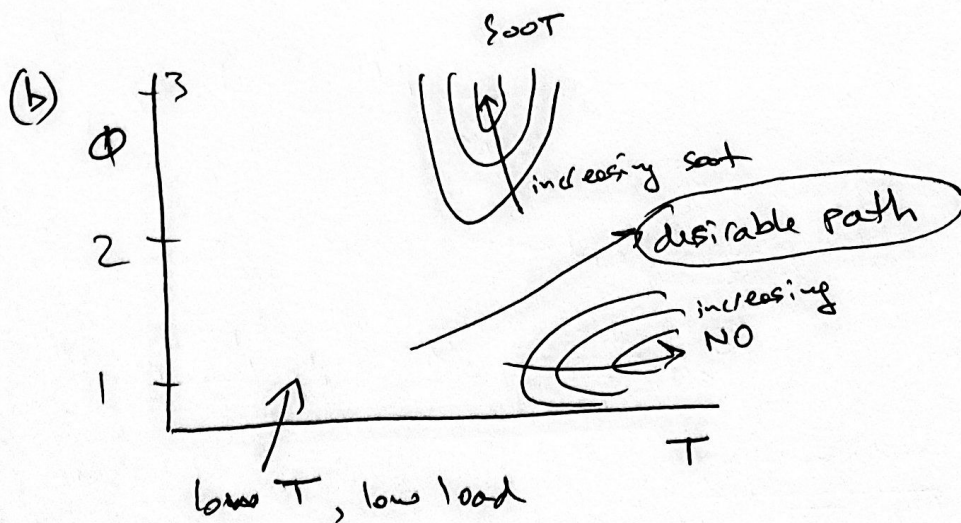
$\Rightarrow N$  &  $d$  are related in the new combustor

Q4 (a) 1	$\eta$	NO <sub>x</sub>	PM	weight
Diesel engine	~45-50%	low with aftertreatment	high aftertreatment needed	high
Gasturb	~40% (single cycle) ~60-64% (combined cycle)	can be very low, no aftertreatment needed	can be low	low

Both can be fuel-flexible (i.e. can be made to burn most fuels).

Maintenance is much higher in diesel engine.

In single cycle, exhaust T of gt gases is high, therefore a combined cycle can be implemented. A lot of the fuel energy is lost in cooling water in the diesel engine & this is low-grade heat that cannot be used easily.



Soot is produced at high T & rich pockets.  
NO<sub>x</sub> is produced at high T & lean regions & long residence times.

To reduce NO<sub>x</sub>, we must operate lean & at short residence times, or include EGR so that flame T reduces.

More premixing (fine droplets, high-P injection) leads to reduced soot.

"low-T combustion", multiple injection strategies lead to low NO<sub>x</sub> & low soot.

If we overdo EGR, we will start having a lot of CO & unburnt hydrocarbon emissions.