

EGT3  
ENGINEERING TRIPoS PART IIB

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Friday 9 May 2025 2 to 3.40

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**Module 4C9**

**CONTINUUM MECHANICS**

*Answer not more than **two** questions.*

*All questions carry the same number of marks.*

*The approximate percentage of marks allocated to each part of a question is indicated in the right margin.*

*Write your candidate number **not** your name on the cover sheet.*

**STATIONERY REQUIREMENTS**

Single-sided script paper

**SPECIAL REQUIREMENTS TO BE SUPPLIED FOR THIS EXAM**

CUED approved calculator allowed

Attachment: 4C9 Continuum Mechanics datasheet (2 pages)

Engineering Data Book

**10 minutes reading time is allowed for this paper at the start of the exam.**

**You may not start to read the questions printed on the subsequent pages of this question paper until instructed to do so.**

**You may not remove any stationery from the Examination Room.**

1 A slender cantilever beam is inclined at an angle  $\theta$  to the horizontal, as shown in Fig. 1. The beam has length  $L$ , depth  $D$  and out-of-plane width  $B$ . The position of a material point on the mid-plane of the beam is  $\mathbf{x} = x_1 \mathbf{e}_1 + x_2 \mathbf{e}_2$ . The loading is provided by self-weight and a vertical tip force of magnitude  $P$ , in the directions shown in the figure. Infinitesimal deformations can be assumed.

(a) The beam is manufactured from a graded linear elastic material with density  $\rho(x_1)$  and Young's modulus  $E(x_1)$  that vary with distance  $x_1$  along the beam.

(i) Show that the displacement field for the mid-plane can be approximated by

$$\mathbf{u}(x_1, x_2) = w \mathbf{e}_2 + \left( h - \frac{dw}{dx_1} x_2 \right) \mathbf{e}_1,$$

where  $h(x_1)$  and  $w(x_1)$  are the displacements of points on the centre-line of the beam in directions  $\mathbf{e}_1$  and  $\mathbf{e}_2$ , respectively. [10%]

(ii) Explain why the elastic strain energy density at a point depends only on the stress component  $\sigma_{11}$  and the strain component  $\varepsilon_{11}$ . [10%]

(iii) Hence, use the method of minimum potential energy to derive governing equations for the centre-line deflections  $h(x_1)$  and  $w(x_1)$ , and for the boundary conditions. A solution for the deflected shape is not required. [40%]

(b) The beam is instead manufactured from a spatially uniform linear viscoelastic material with density  $\rho$  and relaxation modulus  $E_r(t) = E_0 \exp(-E_0 t/\eta)$ .

(i) Explain briefly what is meant by the *correspondence principle*. [10%]

(ii) Using the correspondence principle, derive an expression for the time-dependent axial deflection at the tip of the cantilever,  $h(L, t)$ . [30%]

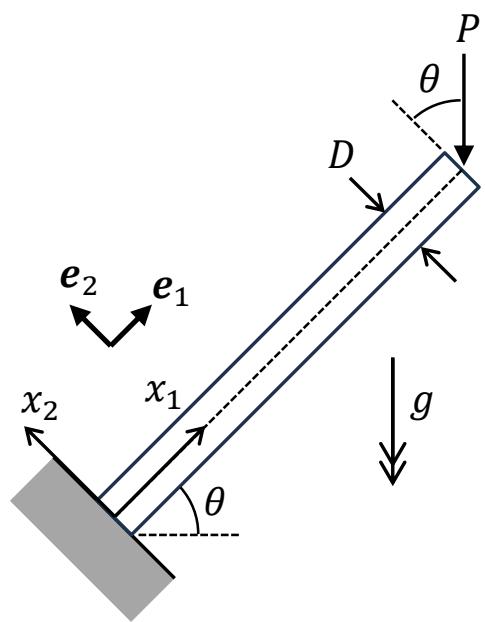


Fig. 1

2 (a) A linear elastic body with Young's modulus  $E$  and Poisson's ratio  $\nu$  undergoes infinitesimal deformation. The displacement field  $\mathbf{u}(\mathbf{x})$ , where the position  $\mathbf{x} = x_i \mathbf{e}_i$ , has components

$$\begin{aligned}u_1 &= -(ax_1 + bx_2)x_3, \\u_2 &= -(bx_1 + cx_2)x_3, \\u_3 &= \frac{1}{2} \left[ ax_1^2 + 2bx_1x_2 + cx_2^2 + \frac{\nu}{1-\nu}(a+c)x_3^2 \right],\end{aligned}$$

where  $a$ ,  $b$  and  $c$  are constants. Prove any requirements on  $a$ ,  $b$  and  $c$  to ensure that:

- (i) the deformation is compatible; and [25%]
- (ii) the body is in equilibrium in the absence of body forces. [25%]

(b) Consider a body with reference configuration  $\Omega_0$  and a current configuration  $\Omega$ . The deformation gradient is  $\mathbf{F}$  and its determinant  $J = \det \mathbf{F}$ .

- (i) Show that  $\nabla_0 Q = \mathbf{F}^T \nabla q$ , where  $Q$  is a scalar field on a reference configuration and  $q$  is its push-forward to the spatial configuration. [10%]
- (ii) Consider a vector  $\mathbf{V}$  on the reference configuration. If the push-forward to the spatial configuration is given by  $\mathbf{v} = J^{-1} \mathbf{F} \mathbf{V}$ , divergences of the vector fields are preserved, i.e.  $\int_{\Omega_0} \nabla_0 \cdot \mathbf{V} d\Omega_0 = \int_{\Omega} \nabla \cdot \mathbf{v} d\Omega$ . Use this result to find the relationship between the outward unit normal vectors to  $\Omega_0$  and to  $\Omega$ . [20%]
- (iii) Prove that  $\int_{\Omega_0} \nabla_0 \cdot \mathbf{V} d\Omega_0 = \int_{\Omega} \nabla \cdot \mathbf{v} d\Omega$  when  $\mathbf{v} = J^{-1} \mathbf{F} \mathbf{V}$ .

Hint: start with  $\int_{\Omega} (\nabla \cdot \mathbf{v}) q d\Omega$ , where  $q$  is a differentiable function that goes to zero on the boundary of  $\Omega$ . [20%]

3 (a) Consider a triangle in the reference configuration with vertices  $\hat{v}_1 = (0, 0)$ ,  $\hat{v}_2 = (1, 0)$ ,  $\hat{v}_3 = (0, 1)$ . Under a transformation the vertices map to  $v_1 = (0, 0)$ ,  $v_2 = (2, 0)$ ,  $v_3 = (2, 2)$ . The deformation gradient,  $\mathbf{F}$ , is constant.

- (i) Sketch the two configurations and give the deformation map  $\phi(X)$ . [20%]
- (ii) By geometric arguments alone, give  $\det \mathbf{F}$ . [10%]
- (iii) Compute the deformation gradient. [10%]
- (iv) Consider two vector fields,  $\mathbf{W}_t$  and  $\mathbf{W}_n$ , on the reference configuration. On the edge  $\hat{v}_2-\hat{v}_3$ ,  $\mathbf{W}_t$  is tangential to the edge and  $\mathbf{W}_n$  is normal to the edge. Apply the transformations  $\mathbf{w}^{(1)} = (\det \mathbf{F})^{-1} \mathbf{F} \mathbf{W}_t$  and  $\mathbf{w}^{(2)} = \mathbf{F}^{-T} \mathbf{W}_n$ , which are both push-forward transformations, to  $\mathbf{W}_t$  and  $\mathbf{W}_n$  at a point on the  $\hat{v}_2-\hat{v}_3$  edge. [20%]
- (v) Sketch the vectors  $\mathbf{W}_t$  and  $\mathbf{W}_n$  at a point on the  $\hat{v}_2-\hat{v}_3$  edge, and on the spatial configuration for the  $\mathbf{w}^{(1)}$  and  $\mathbf{w}^{(2)}$  push-forward transformations. Comment on any significant features. [10%]

(b) What properties make a constitutive model hyperelastic? Comment on any requirements on the kinematic quantities that a constitutive model may depend on. [10%]

(c) Derive the expression for the material time derivative of a scalar quantity on the current configuration, and use this to prove that conservation of mass requires that  $\partial \rho / \partial t + \nabla \cdot (\rho \mathbf{v}) = 0$ , where  $\rho$  is the density and  $\mathbf{v}$  is the velocity.

Note:  $\dot{J} = J \nabla \cdot \mathbf{v}$ . [20%]

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## ENGINEERING TRIPPOS PART IIB

### Module 4C9 Continuum Mechanics

#### Data sheet

##### Indicial notation

A repeated index implies summation

$$\mathbf{a} = a_i \mathbf{e}_i \quad \mathbf{a} \cdot \mathbf{b} = a_i b_i$$

$\mathbf{c} = \mathbf{a} \times \mathbf{b}$  can be written as  $c_i = e_{ijk} a_j b_k$

$\mathbf{A} = \mathbf{a} \otimes \mathbf{b}$  can be written as  $A_{ij} = a_i b_j$

Kronecker delta:  $\delta_{ij} = 1$  for  $i = j$ , and  $\delta_{ij} = 0$  for  $i \neq j$

Note that  $\delta_{ij} = \mathbf{e}_i \cdot \mathbf{e}_j$

Permutation symbol:  $e_{ijk} = 1$  when  $i, j, k$  are in cyclic order

$e_{ijk} = -1$  when  $i, j, k$  are in anti-cyclic order

$e_{ijk} = 0$  when any indices repeat

$e - \delta$  identity:  $e_{ijk} e_{ipq} = \delta_{jp} \delta_{kq} - \delta_{jq} \delta_{kp}$

$\text{grad } \phi = \nabla \phi = \phi_i \mathbf{e}_i$

$\text{div } \mathbf{v} = \nabla \cdot \mathbf{v} = v_{i,i}$

$\text{curl } \mathbf{v} = \nabla \times \mathbf{v} = e_{ijk} v_{k,j} \mathbf{e}_i$

Gauss's theorem (the divergence theorem):

$$\int_V \frac{\partial A_{ij}}{\partial x_j} dV = \oint_S A_{ij} n_j dS$$

Stokes's theorem:

$$\int_S e_{ijk} \frac{\partial A_{pk}}{\partial x_j} n_i dS = \oint_C A_{pk} dx_k$$

## Isotropic linear elasticity

Equilibrium:  $\sigma_{ij,j} + b_i = 0$  ,  $\sigma_{ij} = \sigma_{ji}$

Compatibility:  $\varepsilon_{ij,kp} + \varepsilon_{kp,ij} - \varepsilon_{pj,ki} - \varepsilon_{ki,pj} = 0$

Constitutive relationships:  $\sigma_{ij} = \frac{E}{(1+\nu)} \varepsilon_{ij} + \frac{\nu E}{(1+\nu)(1-2\nu)} \varepsilon_{kk} \delta_{ij}$

Lame's constants:  $\mu = G = \frac{E}{2(1+\nu)}$  ,  $\lambda = \frac{\nu E}{(1+\nu)(1-2\nu)}$

The strain energy density  $U$  is given by:  $\sigma_{ij} = \frac{\partial U}{\partial \varepsilon_{ij}}$

At equilibrium, the potential energy  $\Pi$  is minimised. Hence, for any small kinematically admissible perturbation  $\delta u_i$ :

$$\delta \Pi = \int_V \delta U dV - \int_S t_i^e \delta u_i dS - \int_V b_i \delta u_i dV = 0$$

Definitions:  $\sigma_{ij}$  is the stress tensor,  $\varepsilon_{ij}$  is the infinitesimal strain tensor,  $b_i$  is the body force vector,  $t_i^e$  is the external traction vector and  $u_i$  is the displacement vector.

## Isotropic linear viscoelasticity

Relaxation modulus,  $E_r(t)$ :

if  $\varepsilon(t) = \varepsilon_0 H(t)$ , where  $H(t) = \begin{cases} 0 & t < 0 \\ 1 & t > 0 \end{cases}$  , then  $\sigma(t) = \varepsilon_0 E_r(t)$

Creep compliance,  $J_c(t)$ :

if  $\sigma(t) = \sigma_0 H(t)$ , where  $H(t) = \begin{cases} 0 & t < 0 \\ 1 & t > 0 \end{cases}$  , then  $\varepsilon(t) = \sigma_0 J_c(t)$

The Laplace transforms of  $E_r(t)$  and  $J_c(t)$  are related by:  $\bar{E}_r(s) \bar{J}_c(s) = \frac{1}{s^2}$

Boltzmann superposition principle in 1D:

$$\sigma(t) = \int_0^t \frac{\partial \varepsilon(\tau)}{\partial \tau} E_r(t-\tau) d\tau$$

$$\varepsilon(t) = \int_0^t \frac{\partial \sigma(\tau)}{\partial \tau} J_c(t-\tau) d\tau$$

Correspondence principle: in the Laplace domain, the viscoelastic solution corresponds to the elastic solution, with the substitution  $E \rightarrow s\bar{E}_r(s)$  ,  $\nu \rightarrow s\bar{\nu}_r(s)$  (for any time-dependent moduli).