

EGT3  
ENGINEERING TRIPOS PART IIB

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Monday 3 May 2021 1.30 to 3.10

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**Module 4C9**

**CONTINUUM MECHANICS**

*Answer not more than **two** questions.*

*All questions carry the same number of marks.*

*The **approximate** percentage of marks allocated to each part of a question is indicated in the right margin.*

*Write your candidate number **not** your name on the cover sheet and at the top of each answer sheet.*

**STATIONERY REQUIREMENTS**

Write on single-sided paper.

**SPECIAL REQUIREMENTS TO BE SUPPLIED FOR THIS EXAM**

CUED approved calculator allowed.

Attachment: 4C9 datasheet (2 pages).

You are allowed access to the electronic version of the Engineering Data Books.

**10 minutes reading time is allowed for this paper at the start of the exam.**

**The time taken for scanning/uploading answers is 15 minutes.**

**Your script is to be uploaded as a single consolidated pdf containing all answers.**

1 A thin-walled circular tube of length  $L$ , radius  $R$  and wall thickness  $t$  is loaded by a torque  $T$  as shown in Fig. 1. The tube is constrained such that the displacement  $\mathbf{u} = 0$  at  $x_1 = 0$ . The initial position  $\mathbf{x}$  and displacement  $\mathbf{u}$  of a point on the tube wall (assuming plane stress, and infinitesimal strains and rotations) are given by

$$\mathbf{x} = x_1 \mathbf{e}_1 + R \mathbf{e}_3, \quad \mathbf{u} = u_2(x_1) \mathbf{e}_2.$$

Due to the symmetry, this can describe any material point following a rotation of the orthonormal basis vectors  $(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$  about the  $\mathbf{e}_1$  direction.

The tube is 3D-printed from a linear elastic material. Due to printing defects, the Young's modulus and the wall thickness both vary with axial position, i.e.  $E(x_1)$  and  $t(x_1)$ . At a given  $x_1$ , these properties are constant around the circumference of the tube. The Poisson ratio  $\nu$  is constant everywhere.

(a) Show that the strain energy density  $U(x_1)$  is given by:

$$U = \frac{E(x_1)}{4(1+\nu)} (u_{2,1})^2.$$

[20%]

(b) Show that the potential energy  $\Pi$  is given by:

$$\Pi = \int_0^L \frac{\pi R}{2(1+\nu)} E(x_1) t(x_1) (u_{2,1})^2 dx_1 - \frac{T}{R} u_2(L).$$

[20%]

(c) Using the method of minimum potential energy, derive the governing equation for  $u_2(x_1)$  along the length of the tube, and the associated boundary conditions. [50%]

(d) Measurements provide data for  $E(x_1)$  and  $t(x_1)$ . Explain, without further calculation, how an expression for the rotation of the tube at  $x_1 = L$  would be obtained. [10%]

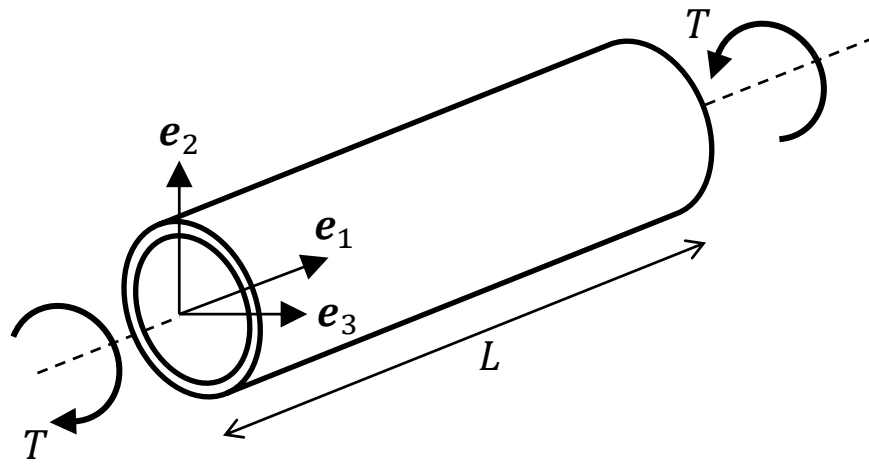


Fig. 1

2 (a) The position vector  $\mathbf{x}$  is defined using orthonormal basis vectors  $(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$ . Use index notation to answer the following.

(i) If  $\phi(\mathbf{x})$  is a scalar field, show that  $\nabla \times (\nabla \phi) = 0$ . [10%]

(ii) For rotated orthonormal basis vectors  $\hat{\mathbf{e}}_i = R_{ij}\mathbf{e}_j$ , show that

$$R_{ip}R_{jp} = R_{pi}R_{pj} = \delta_{ij}.$$

[20%]

(iii) For the tensor  $\mathbf{A} = A_{ij}\mathbf{e}_i \otimes \mathbf{e}_j = \hat{A}_{ij}\hat{\mathbf{e}}_i \otimes \hat{\mathbf{e}}_j$ , derive an expression for the relationship between components  $A_{ij}$ ,  $\hat{A}_{ij}$ , and  $R_{ij}$ . [20%]

(b) The Green–Lagrange strain is given by  $\mathbf{E} := (1/2)(\mathbf{F}^T\mathbf{F} - \mathbf{I})$ , where  $\mathbf{F}$  is the deformation gradient and  $\mathbf{I}$  is the identity tensor. Express the Green–Lagrange strain using index notation, and use this to show that the Green–Lagrange strain reduces to the linearised strain measure for linearised kinematics. [10%]

(c) Consider a body in 3D that undergoes an arbitrary translation, plus a rotation about the  $X_3$  axis.

(i) Compute the deformation gradient  $\mathbf{F}$ . [10%]

(ii) Without calculation, give  $\det \mathbf{F}$ . Explain your reasoning. [5%]

(iii) By considering an axis-aligned unit cube, demonstrate that  $ds = J\mathbf{F}^{-T} d\mathbf{S}$  (Nanson’s formula) holds. [20%]

(iv) Compute the Green–Lagrange strain. [5%]

- 3 (a) The deformation gradient  $\mathbf{F}$  can be decomposed as:

$$\mathbf{F} = \mathbf{R}\mathbf{U} = \mathbf{V}\mathbf{R},$$

where  $\mathbf{R}$  is a rotation, and  $\mathbf{U}$  and  $\mathbf{V}$  are symmetric. This is known as the polar decomposition.

- (i) Using the singular value decomposition of a matrix  $\mathbf{A} = \mathbf{W}\mathbf{\Sigma}\mathbf{Z}^T$ , where  $\mathbf{W}$  and  $\mathbf{Z}$  are orthogonal matrices and  $\mathbf{\Sigma}$  holds the singular values of  $\mathbf{A}$ , show that the polar decomposition of the deformation gradient exists. [30%]
- (ii) Show that the stretch tensor  $\mathbf{U}$  is symmetric and that all eigenvalues of  $\mathbf{U}$  are positive. Comment on the physical significance of the eigenvalues being positive. [20%]
- (iii) Comment on the significance of the polar decomposition for the definition of strain tensors. [10%]

- (b) Starting from the balance of momentum in the material configuration

$$\rho_0 \ddot{\boldsymbol{\phi}} = \nabla_X \cdot \mathbf{P} + \rho_0 \mathbf{b}_m,$$

where  $\rho_0$  is the reference density,  $\boldsymbol{\phi}$  is the deformation map,  $\mathbf{P}$  is the first Piola–Kirchhoff stress and  $\mathbf{b}_m$  is the material body force, show that  $\mathbf{P}$  is work-conjugate to  $\dot{\mathbf{F}}$ . [10%]

- (c) The Hu–Washizu variational principle involves a total potential energy functional of the form

$$\Pi(\boldsymbol{\phi}, \mathbf{F}, \mathbf{P}) := \int_{\Omega_0} \psi(\mathbf{F}) \, dV + \int_{\Omega_0} \mathbf{P} : (\nabla_X \boldsymbol{\phi} - \mathbf{F}) \, dV - \Pi_{\text{external}}(\boldsymbol{\phi}),$$

where  $\boldsymbol{\phi}$ ,  $\mathbf{F}$  and  $\mathbf{P}$  are treated as independent variables and  $\Pi_{\text{external}}$  is associated with loadings. If equilibrium corresponds to a stationary point of  $\Pi$ , give any relationships that must exist between  $\psi$ ,  $\boldsymbol{\phi}$ ,  $\mathbf{F}$  and  $\mathbf{P}$  at equilibrium. [30%]

**END OF PAPER**

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# ENGINEERING TRIPOS PART IIB

## Module 4C9 Continuum Mechanics

### Data sheet

#### Indicial notation

A repeated index implies summation

$$\mathbf{a} = a_i \mathbf{e}_i \quad \mathbf{a} \cdot \mathbf{b} = a_i b_i$$

$$\mathbf{c} = \mathbf{a} \times \mathbf{b} \text{ can be written as } c_i = e_{ijk} a_j b_k$$

$$\mathbf{A} = \mathbf{a} \otimes \mathbf{b} \text{ can be written as } A_{ij} = a_i b_j$$

$$\text{Kronecker delta: } \delta_{ij} = 1 \text{ for } i = j, \text{ and } \delta_{ij} = 0 \text{ for } i \neq j$$

$$\text{Note that } \delta_{ij} = \mathbf{e}_i \cdot \mathbf{e}_j$$

$$\text{Permutation symbol: } e_{ijk} = 1 \text{ when } i, j, k \text{ are in cyclic order}$$

$$e_{ijk} = -1 \text{ when } i, j, k \text{ are in anti-cyclic order}$$

$$e_{ijk} = 0 \text{ when any indices repeat}$$

$$e - \delta \text{ identity: } e_{ijk} e_{ipq} = \delta_{jp} \delta_{kq} - \delta_{jq} \delta_{kp}$$

$$\text{grad } \phi = \nabla \phi = \phi_{,i} \mathbf{e}_i$$

$$\text{div } \mathbf{v} = \nabla \cdot \mathbf{v} = v_{i,i}$$

$$\text{curl } \mathbf{v} = \nabla \times \mathbf{v} = e_{ijk} v_{k,j} \mathbf{e}_i$$

Gauss's theorem (the divergence theorem):

$$\int_V \frac{\partial A_{ij}}{\partial x_j} dV = \oint_S A_{ij} n_j dS$$

Stokes's theorem:

$$\int_S e_{ijk} \frac{\partial A_{pk}}{\partial x_j} n_i dS = \oint_C A_{pk} dx_k$$

## Isotropic linear elasticity

Equilibrium:  $\sigma_{ij,j} + b_i = 0$  ,  $\sigma_{ij} = \sigma_{ji}$

Compatibility:  $\varepsilon_{ij,kp} + \varepsilon_{kp,ij} - \varepsilon_{pj,ki} - \varepsilon_{ki,pj} = 0$

Constitutive relationships:  $\sigma_{ij} = \frac{E}{(1+\nu)} \varepsilon_{ij} + \frac{\nu E}{(1+\nu)(1-2\nu)} \varepsilon_{kk} \delta_{ij}$

Lame's constants:  $\mu = G = \frac{E}{2(1+\nu)}$  ,  $\lambda = \frac{\nu E}{(1+\nu)(1-2\nu)}$

The strain energy density  $U$  is given by:  $\sigma_{ij} = \frac{\partial U}{\partial \varepsilon_{ij}}$

At equilibrium, the potential energy  $\Pi$  is minimised. Hence, for any small kinematically admissible perturbation  $\delta u_i$  :

$$\delta \Pi = \int_V \delta U dV - \int_S t_i^e \delta u_i dS - \int_V b_i \delta u_i dV = 0$$

Definitions:  $\sigma_{ij}$  is the stress tensor,  $\varepsilon_{ij}$  is the infinitesimal strain tensor,  $b_i$  is the body force vector,  $t_i^e$  is the external traction vector and  $u_i$  is the displacement vector.

## Isotropic linear viscoelasticity

Relaxation modulus,  $E_r(t)$ :

if  $\varepsilon(t) = \varepsilon_0 H(t)$  , where  $H(t) = \begin{cases} 0 & t < 0 \\ 1 & t > 0 \end{cases}$  , then  $\sigma(t) = \varepsilon_0 E_r(t)$

Creep compliance,  $J_c(t)$ :

if  $\sigma(t) = \sigma_0 H(t)$  , where  $H(t) = \begin{cases} 0 & t < 0 \\ 1 & t > 0 \end{cases}$  , then  $\varepsilon(t) = \sigma_0 J_c(t)$

The Laplace transforms of  $E_r(t)$  and  $J_c(t)$  are related by:  $\bar{E}_r(s) \bar{J}_c(s) = \frac{1}{s^2}$

Boltzmann superposition principle in 1D:

$$\sigma(t) = \int_0^t \frac{\partial \varepsilon(\tau)}{\partial \tau} E_r(t - \tau) d\tau$$

$$\varepsilon(t) = \int_0^t \frac{\partial \sigma(\tau)}{\partial \tau} J_c(t - \tau) d\tau$$

Correspondence principle: in the Laplace domain, the viscoelastic solution corresponds to the elastic solution, with the substitution  $E \rightarrow s \bar{E}_r(s)$  ,  $\nu \rightarrow s \bar{\nu}_r(s)$  (for any time-dependent moduli).