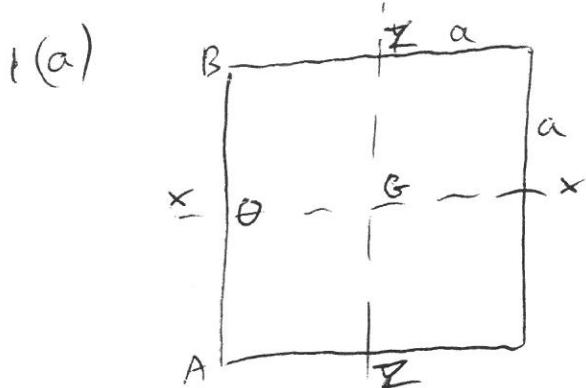


# Module 3CS 2017 Solutions

Q



From Databook

$$I_{xx} = \frac{1}{3}ma^2 = I_{zz}$$

$$I_{yy} = I_{xx} + I_{zz} = \frac{2}{3}ma^2$$

Parallel axis theorem to O

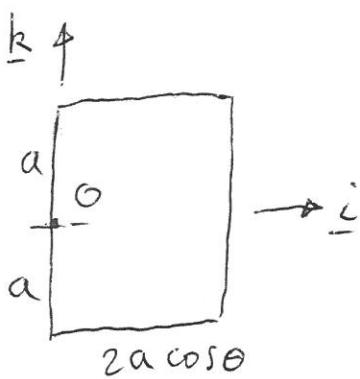
$$\therefore I_0 = \begin{bmatrix} A & 0 & 0 \\ 0 & B & 0 \\ 0 & 0 & C \end{bmatrix} \quad \text{with } A = \frac{1}{3}ma^2$$

$$B = \frac{2}{3}ma^2 + ma^2 = \frac{5}{3}ma^2$$

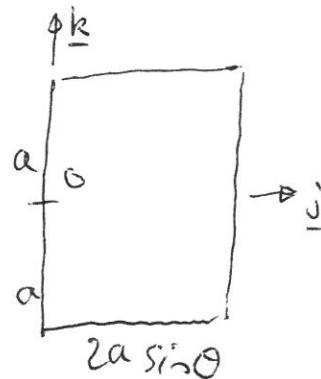
$$C = \frac{1}{3}ma^2 + ma^2 = \frac{4}{3}ma^2$$

(b)

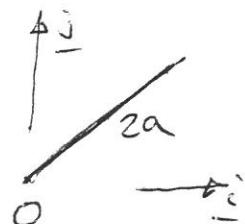
View  $i \perp k$



View  $j \perp k$



View  $i \perp j$



$$I_{yy} = \frac{4}{3}m(\cos\theta)^2 + \frac{1}{3}ma^2$$

$$I_{xx} = \text{similarly}$$

~~$$I_{zz} = \frac{4}{3}ma^2$$~~

$$= \frac{1}{3}ma^2(1 + 4\cos^2\theta)$$

$$= \frac{1}{3}ma^2(1 + 4\sin^2\theta)$$

$$I_{xz} = 0$$

$$I_{yz} = 0$$

$$I_{xy} = \int r^2 dy dm$$

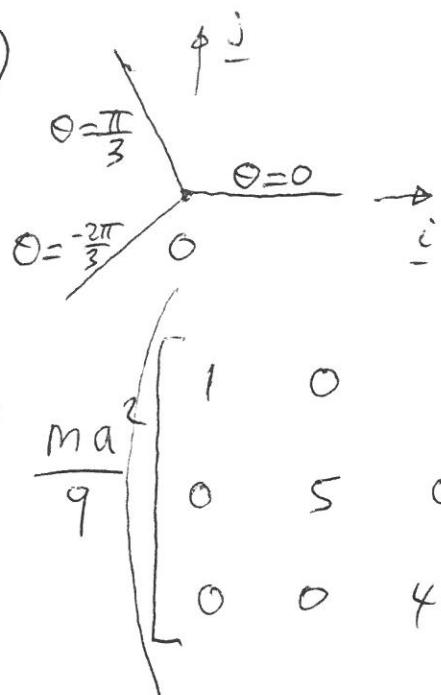
$$\therefore I_0 = \begin{bmatrix} 1 + 4\sin^2\theta & -4\cos\theta\sin\theta & 0 \\ -4\cos\theta\sin\theta & 1 + 4\cos^2\theta & 0 \\ 0 & 0 & 4 \end{bmatrix} \frac{ma^2}{3}$$

$$= \int_0^{2a} r \cos\theta r \sin\theta \frac{m}{2a} d$$

$$= \frac{m \cos\theta \sin\theta}{2a} \left[ \frac{r^3}{3} \right]_0^{2a}$$

$$= \frac{4ma^2}{3} \cos\theta \sin\theta$$

1. cont (c)



Expect AAC because ②  
of three axis symmetry  
Three identical eigenvalues

$$I_0 = \frac{ma^2}{9} \left( \begin{bmatrix} 1 & 0 & 0 \\ 0 & 5 & 0 \\ 0 & 0 & 4 \end{bmatrix} + \begin{bmatrix} 1+4\left(\frac{\sqrt{3}}{2}\right)^2 - 4\left(\frac{-1}{2}\right)\frac{\sqrt{3}}{2} & 0 \\ -4\left(\frac{-1}{2}\right)\frac{\sqrt{3}}{2} & 1+4\left(\frac{-1}{2}\right)^2 & 0 \\ 0 & 0 & 4 \end{bmatrix} \right.$$

$$\theta=0$$

$$\theta=\frac{2\pi}{3}$$

$$\left. + \begin{bmatrix} 1+4\left(\frac{-\sqrt{3}}{2}\right)^2 - 4\left(\frac{-1}{2}\right)\left(\frac{-\sqrt{3}}{2}\right) & 0 \\ -4\left(\frac{-1}{2}\right)\left(\frac{-\sqrt{3}}{2}\right) & 1+4\left(\frac{-1}{2}\right)^2 & 0 \\ 0 & 0 & 4 \end{bmatrix} \right)$$

$$= \frac{ma^2}{9} \begin{bmatrix} 9 & 0 & 0 \\ 0 & 9 & 0 \\ 0 & 0 & 12 \end{bmatrix} \quad \theta = \frac{2\pi}{3}$$

$$= ma^2 \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \frac{4}{3} \end{bmatrix}$$

(d) For  $N$  plates of mass  $m$  in total

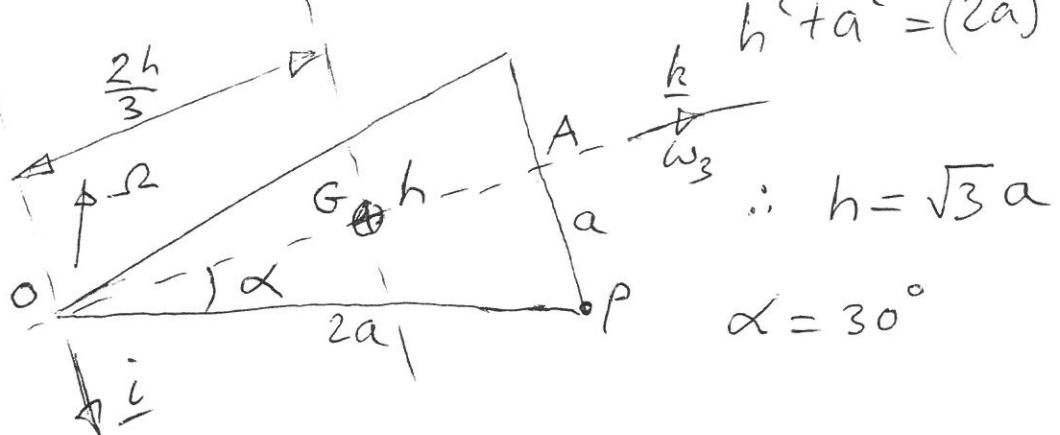
$$I_0 = ma^2 \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \frac{4}{3} \end{bmatrix} \quad \text{because if the}$$

three plates is AAC and indistinguishable from a cylinder of mass  $m$  then this must be so for  $N$  plates equispaced.

(3)

## Q2 Conical shell

a/



$$OG = \frac{2}{3}h = \frac{2}{3}\sqrt{3}a$$

AAC body

$$A = m \left( \frac{a^2}{4} + \frac{h^2}{18} \right) + m \overline{OG}^2$$

$$\boxed{C = \frac{ma^2}{2}}$$

$$= ma^2 \left( \frac{1}{4} + \frac{3}{18} + \frac{4}{9} 3 \right)$$

$$\boxed{A = \frac{7}{4}ma^2}$$

b/ No slip at P

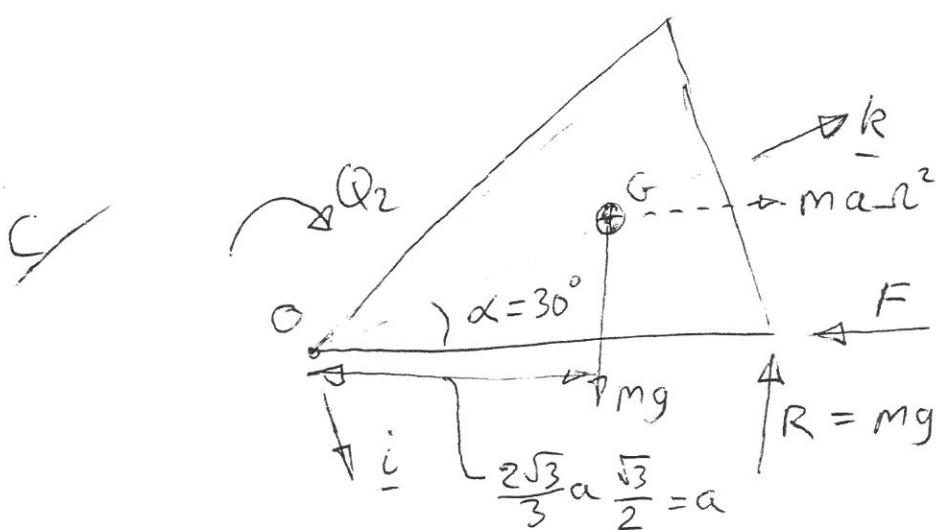
$$\underline{\omega}_P = \underline{\omega}_A + \omega_3 \underline{k} \times \underline{a}_c = 0$$

(note  $\overline{OA} \cos \alpha$   
is needed to  
get  $\underline{v}_A$ )

$$\Rightarrow -2h \cos \alpha + \omega_3 a = 0$$

$$\therefore \omega_3 = -\Omega \sqrt{3} \frac{\sqrt{3}}{2}$$

$$= -\frac{3\Omega}{2}$$



$$\begin{aligned} Q_2 &= -mg 2a \\ &\quad + m g a \\ &= -mg a \end{aligned}$$

(4)

Gyroscope equations, Q<sub>2</sub> equation is  
useful

$$Q_2 = A\dot{\alpha}_2 + (A\alpha_3 - c\omega_3)\alpha_4$$

$$\alpha_1 = -\dot{\phi} \sin \theta$$

$$\alpha_2 = \dot{\theta}$$

$$\alpha_3 = \dot{\phi} \cos \theta$$

$$\text{with } \dot{\phi} = \alpha$$

$$\theta = \frac{\pi}{2} - \alpha \quad \dot{\theta} = 0$$

$$\therefore \alpha_1 = -\alpha \cos \alpha = -\alpha \frac{\sqrt{3}}{2}$$

$$\alpha_2 = 0$$

$$\alpha_3 = \alpha \sin \alpha = \alpha / 2$$

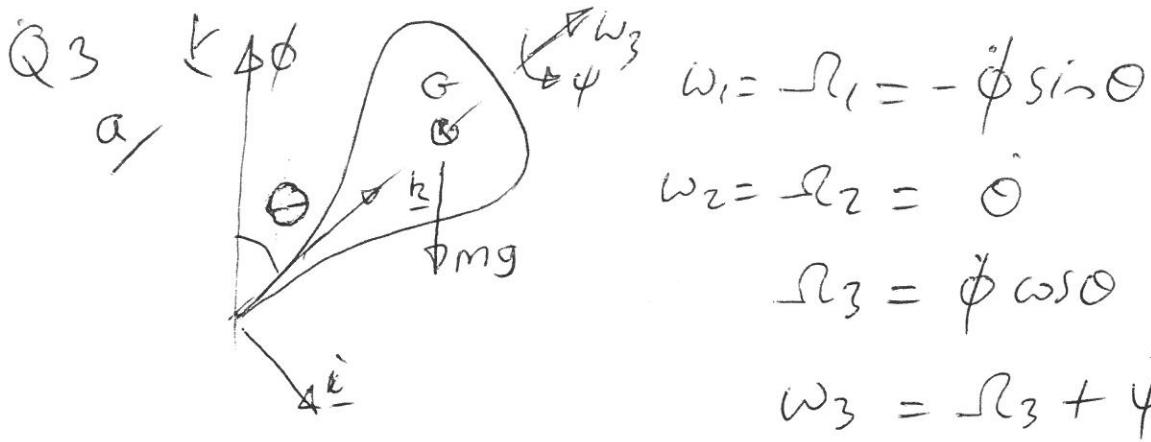
$$\therefore -mg\alpha = 0 + \left( \frac{7}{4}ma^2 \frac{\alpha}{2} - \frac{ma^2}{2} \left( -\frac{3\alpha}{2} \right) \right) \left( -\alpha \frac{\sqrt{3}}{2} \right)$$

$$= \left( \frac{7}{2} + 3 \right) ma^2 \alpha^2 \frac{\sqrt{3}}{8}$$

$$\therefore \alpha^2 = \frac{16}{13\sqrt{3}} \frac{g}{a} = \frac{16\sqrt{3}}{39} \frac{g}{a}$$

$$\begin{aligned} d/ \mu &= \frac{F}{R} = \frac{ma\alpha^2}{mg} \\ &= 16\sqrt{3}/39 \end{aligned}$$

(5)



b/  $T = \frac{1}{2} A (\dot{\phi} \sin \theta)^2 + \frac{1}{2} A \dot{\theta}^2 + \frac{1}{2} C (\dot{\phi} \cos \theta + \dot{\psi})^2$

$$V = mg a (\cos \theta - 1)$$

Lagrange:  $\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{q}_i} \right) - \frac{\partial T}{\partial q_i} + \frac{\partial V}{\partial z_i} = Q_i$

$\psi: \frac{\partial T}{\partial \dot{\psi}} = C(\dot{\phi} \cos \theta + \dot{\psi})$

$$\frac{\partial T}{\partial \dot{\psi}} = 0$$

$$\frac{\partial V}{\partial \psi} = 0$$

conservation of moment  
of momentum about k

$$\therefore \boxed{C(\dot{\phi} \cos \theta + \dot{\psi}) = h, \text{ a constant}}$$

$\phi: \frac{\partial T}{\partial \dot{\phi}} = A \sin^2 \theta \dot{\phi} + C(\dot{\phi} \cos \theta + \dot{\psi}) \cos \theta$

$$\frac{\partial T}{\partial \dot{\phi}} = 0$$

$$\frac{\partial V}{\partial \phi} = 0$$

conservation of moment  
of momentum about k

$$\therefore \boxed{A \sin^2 \theta \dot{\phi} + h \cos \theta = H, \text{ a constant}}$$

(6)

$$\theta: \frac{\partial T}{\partial \dot{\theta}} = A \ddot{\theta}$$

$$\begin{aligned} \frac{\partial T}{\partial \theta} &= A \dot{\phi}^2 \sin \theta \cos \theta - c \dot{\phi} \sin \theta (\dot{\phi} \omega_0 + \dot{\psi}) \\ &= A \dot{\phi}^2 \sin \theta \cos \theta - h \dot{\phi} \sin \theta \end{aligned}$$

$$\frac{\partial V}{\partial \theta} = -mg a \sin \theta$$

$$\therefore \boxed{A \ddot{\theta} - A \dot{\phi}^2 \sin \theta \cos \theta + h \dot{\phi} \sin \theta - mg a \sin \theta = 0} \quad [3]$$

c/ Steady precession,  $\dot{\theta} = \text{constant}$

$\dot{\phi} = \text{constant}$

$\dot{\psi} = \text{constant}$

$$\text{Use } [3] \quad \therefore A \dot{\phi}^2 \cos \theta - h \dot{\phi} + mg a = 0$$

$$\therefore A \dot{\phi}^2 \cos \theta - c(\dot{\phi} \cos \theta + \dot{\psi}) \dot{\phi} + mg a = 0 \quad [4]$$

$$\therefore (A - c) \cos \theta \dot{\phi}^2 - c \dot{\psi} \dot{\phi} + mg a = 0$$

$$\therefore \dot{\phi} = \frac{c \dot{\psi} \pm \sqrt{(c \dot{\psi})^2 - 4(A - c) mg a \cos \theta}}{2(A - c) \cos \theta} \quad [5]$$

For fast spin (ie  $\dot{\psi}$  large) [1] gives  $h \approx c \dot{\psi}$

Then [4] gives  $-c \dot{\psi} \dot{\phi} + mg a \approx 0$

(assuming  $\dot{\phi} \ll \dot{\psi}$  - ie slow precessio.)

$$\therefore \boxed{\dot{\phi} \approx \frac{mg a}{c \dot{\psi}}} \quad [6]$$

- The second solution use ④ but this time with  $\phi$  comparable with  $\dot{\psi}$   
 $\therefore$  neglect  $mga$

$$\therefore A \dot{\phi} \cos \omega \approx C(\dot{\phi} \cos \omega + \dot{\psi})$$

$$\therefore \psi(A-C) \cos \varphi = C \psi$$

$$\therefore \dot{\phi} = \frac{c\dot{\psi}}{(A-c)\cos\Omega}$$

Alternatively Note  $\sqrt{1+\epsilon} \simeq 1 + \frac{1}{2}\epsilon$   
 use [5]

$$\text{So } \dot{\phi} = c\dot{\psi} \pm c\ddot{\psi} \sqrt{1 - \frac{4(A-c)}{(c\dot{\psi})^2} \cos \omega m \theta}$$

$$\text{The } "-" \text{ solution gives } \dot{\phi} = \frac{\frac{1}{2} \frac{4(1-\epsilon) \cos \theta m \omega}{(\epsilon - \dot{\phi})^2}}{2(1-\epsilon) \cos \theta}$$

$$\approx \frac{mg a}{c \dot{\varphi}}$$

$$\text{The "+" solution gives } \dot{\phi} = \frac{2c\dot{\varphi}}{2(A-c)\cos\theta} \quad \left. \begin{array}{l} \text{BEFORE} \\ \text{*} \end{array} \right\}$$

$$= \frac{c\dot{\varphi}}{(A-c)\cos\theta} \quad \text{*}$$

NOTE : For  $\dot{\psi} \gg \dot{\phi}$  then  $\dot{\psi} \approx \omega_3$

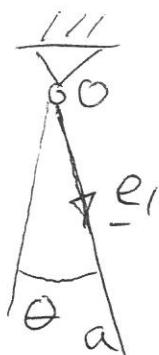
$$\text{So } \dot{\phi} \approx \frac{m g a}{C w_3} \quad \text{as well}$$

but for  $\phi$  of same order as  $\psi$  they use

$$\omega_3 = \dot{\phi} \cos\theta + \dot{\psi} \quad \text{to give} \quad \dot{\phi} = \frac{c\omega_3}{A \cos\theta}$$

4/ The had hit is to get the velocity of  $B$ . Hence the "show that" for  $T$ .

$$\underline{v}_B = a \underline{e}_1 + b \underline{e}_2$$



view along  $\underline{k}$   
 $\underline{k}$  out of page

$$\underline{v}_B = a(\dot{\theta} \underline{k}) \times \underline{e}_1 + b(\dot{\theta} \underline{k} - \dot{\phi} \underline{e}_1) \times \underline{e}_2$$

$$\text{Define } \underline{k} \times \underline{e}_1 = \underline{e}_1^*$$

which is in the plane  
of paper

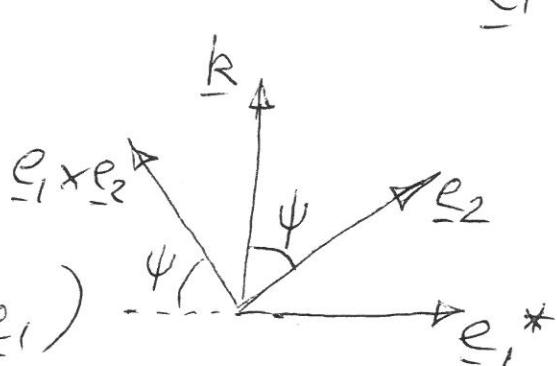


(note AB  
isn't in the  
plane of paper)

view along  $\underline{e}_1$

Resolve  $\underline{v}_B$   
into  $\underline{e}_1$   $\underline{e}_1^*$   $\underline{k}$   
directions

$$(\text{note } \underline{k} \times \underline{e}_2 = -\sin \psi \underline{e}_1)$$



$\underline{e}_1$  out of  
page

$$\underline{v}_B = a \dot{\theta} \underline{e}_1^* + b \dot{\theta} (-\sin \psi \underline{e}_1) - b \dot{\psi} (\underline{e}_1 \times \underline{e}_2)$$

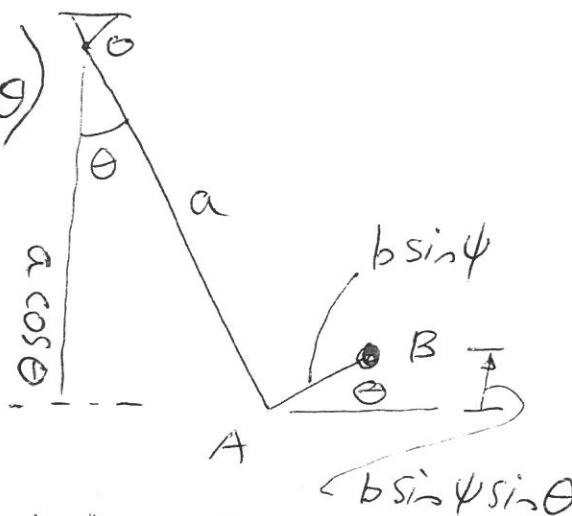
$$\text{note } \underline{e}_1 \times \underline{e}_2 = -\cos \psi \underline{e}_1^* + \sin \psi \underline{k}$$

$$\therefore \underline{v}_B = (a \dot{\theta} + b \dot{\psi} \cos \psi) \underline{e}_1^* - b \dot{\theta} \sin \psi \underline{e}_1 - b \dot{\psi} \sin \psi \underline{k}$$

$$\begin{aligned}
 \text{So } T &= \frac{1}{2}m \left( (a\dot{\theta} + b\dot{\psi}\cos\psi)^2 + (b\dot{\theta}\sin\psi)^2 + (b\dot{\psi}\sin\psi)^2 \right) \quad (9) \\
 &= \frac{1}{2}m \left( a^2\dot{\theta}^2 + 2ab\dot{\theta}\dot{\psi}\cos\psi + b^2\dot{\psi}^2\cos^2\psi \right. \\
 &\quad \left. + b^2\dot{\theta}^2\sin^2\psi + b^2\dot{\psi}^2\sin^2\psi \right) \\
 &= \frac{1}{2}m \left( a^2\dot{\theta}^2 + b^2\dot{\psi}^2 + b^2\dot{\theta}^2\sin^2\psi + 2ab\dot{\theta}\dot{\psi}\cos\psi \right)
 \end{aligned}$$

and

$$V = mg(a(1-\cos\theta) + b\sin\psi\sin\theta)$$



∴:

$$\frac{\partial T}{\partial \dot{\theta}} = m(a^2 + b^2\sin^2\psi)\dot{\theta} + mab\dot{\psi}\cos\psi$$

$$\frac{\partial T}{\partial \theta} = 0 \quad \frac{\partial V}{\partial \theta} = mg(a\sin\theta + b\sin\psi\cos\theta)$$

$$\begin{aligned}
 \therefore \frac{d}{dt} \left[ m(a^2 + b^2\sin^2\psi)\dot{\theta} + mab\dot{\psi}\cos\psi \right] + mg(a\sin\theta + b\sin\psi\cos\theta) \\
 = 0
 \end{aligned}$$

$$\begin{aligned}
 \therefore ma^2\ddot{\theta} + mb^2\dot{\theta}\dot{\psi}\sin\psi\cos\psi + mab\cos\psi\ddot{\psi} \\
 + mb^2\sin^2\psi\ddot{\theta} - mab\dot{\psi}^2\sin\psi \\
 + mga\sin\theta + mgbs\sin\psi\cos\theta = 0
 \end{aligned}$$

(This was important to get full marks because the "show that" for  $\psi$  was too easy to fudge!)

$$\psi: \frac{\partial T}{\partial \dot{\psi}} = mb^2 \ddot{\psi} + mab\dot{\theta} \cos \psi$$

(10)

$$\frac{\partial T}{\partial \psi} = mb^2 \dot{\theta}^2 \sin \psi \cos \psi - mab \dot{\theta} \dot{\psi} \sin \psi$$

$$\frac{\partial V}{\partial \psi} = mg b \cos \psi \sin \theta$$

$$\begin{aligned} \frac{d}{dt} [mb^2 \ddot{\psi} + mab \dot{\theta} \cos \psi] - mb^2 \dot{\theta}^2 \sin \psi \cos \psi \\ + mab \dot{\theta} \dot{\psi} \sin \psi \\ + mg b \cos \psi \sin \theta = 0 \end{aligned}$$

$$\therefore b \ddot{\psi} + a \dot{\theta}^2 \cos \psi - a \dot{\theta} \dot{\psi} \sin \psi - b \dot{\theta}^2 \sin \psi \cos \psi \\ + a \dot{\theta} \dot{\psi} \sin \psi + g \cos \psi \sin \theta = 0$$

$$\therefore \boxed{b \ddot{\psi} + a \dot{\theta}^2 \cos \psi - b \dot{\theta}^2 \sin \psi \cos \psi + g \cos \psi \sin \theta = 0}$$

$\checkmark$  linearize  $T$  &  $V$  and neglect small terms

$$T = \frac{1}{2}m(a^2 \dot{\theta}^2 + b^2 \dot{\psi}^2 + 2ab \dot{\theta} \dot{\psi})$$

$$V = mg(\frac{1}{2}a\dot{\theta}^2 + b\dot{\theta}\psi)$$

$$\therefore [M] = m \begin{bmatrix} a^2 & ab \\ ab & b^2 \end{bmatrix} \text{ and } [k] = mg \begin{bmatrix} a & b \\ b & 0 \end{bmatrix}$$

(11)

4/d/ for natural frequencies

$$|(K - M\omega^2)| = 0$$

$$\therefore \begin{vmatrix} mg\alpha - ma^2\omega^2 & mgb - mab\omega^2 \\ mgb - mab\omega^2 & 0 - mb^2\omega^2 \end{vmatrix} = 0$$

$$\therefore (g\alpha - a^2\omega^2)(-b^2\omega^2) - (gb - ab\omega^2)^2 = 0$$

$$\therefore (g - a\omega^2) a\omega^2 + (g - a\omega^2)^2 = 0$$

$$\therefore (g - a\omega^2)(a\omega^2 + g - a\omega^2) = 0$$

$$\therefore g(g - a\omega^2) = 0$$

$$\therefore \omega^2 = \frac{g}{a}$$

There appears to be only  
one solution!

