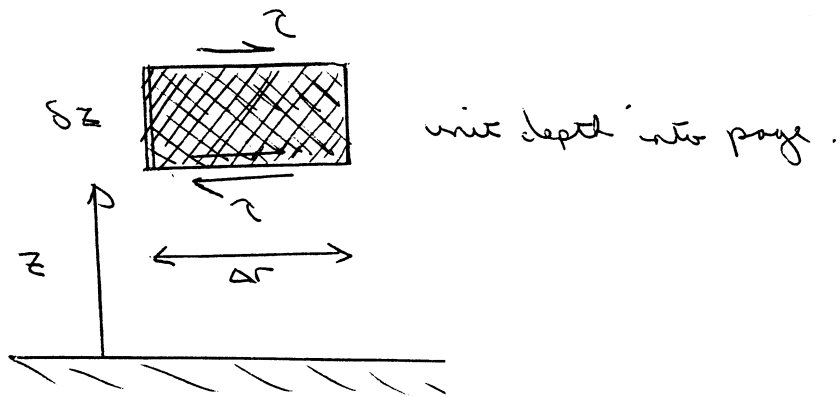


Section A

1.

(a)



net radial force:  $F_r = \tau(z+\delta z)\Delta r - \tau(z)\Delta r$

(no pressure forces as  $p = p_a$  at all points)

net radial acceleration:  $a_r = -\frac{u_0^2}{r} = -\Omega^2 r$

(ignore  $u_r \frac{\partial u_r}{\partial r}$  since  $u_r \ll u_0$ )

Newton II: 
$$\underbrace{[\tau(z+\delta z) - \tau(z)]\Delta r}_{F_r} = \underbrace{(\rho\Delta r\Delta z)}_m \underbrace{(-\Omega^2 r)}_{a_r}$$

$$\Rightarrow \underline{\underline{\tau(z+\delta z) - \tau(z) = -\rho\Delta z\Omega^2 r}}$$

(b)  $\tau = \rho v \frac{\partial u_r}{\partial z}, \quad \frac{\partial \tau}{\partial z} = -\rho\Omega^2 r$

$$\Rightarrow \cancel{\rho v} \frac{\partial^2 u_r}{\partial z^2} = -\cancel{\rho}\Omega^2 r$$

$$\Rightarrow \frac{\partial u_r}{\partial z} = -\frac{\Omega^2 r}{v}(z+A)$$

Boundary condition at  $z = h$  is  $\tau = 0 \Rightarrow \left. \frac{\partial u_r}{\partial z} \right|_h = 0$

Thus  $A = -h$  and,

$$\frac{\partial u_r}{\partial z} = \frac{\Omega^2 r}{\nu} (h - z)$$

Integrate again,

$$\underline{\underline{u_z = \frac{\Omega^2 r}{\nu} \left( hz - \frac{z^2}{2} \right)}}$$

(Constant of integration is zero since  $u_z = 0$  at  $z = 0$ )

$$(c) \quad Q = 2\pi r \int_0^h u_z dz = \frac{2\pi \Omega^2 r^2}{\nu} \left( \frac{h^3}{2} - \frac{h^3}{6} \right)$$

$$\Rightarrow \underline{\underline{Q = \frac{2\pi \Omega^2 r^2}{3\nu} h^3}}$$

Continuity of mass  $\Rightarrow Q = \text{const.}$

$$\Rightarrow \underline{\underline{h \sim r^{-2/3}}}$$

2.

$$(a) \quad \rho g h_1 + \frac{1}{2} \rho v_1^2 + \cancel{p_1} = \rho g h_2 + \frac{1}{2} \rho v_2^2 + \cancel{p_2}$$

(Bernoulli)

$$\Rightarrow \frac{1}{2} \cancel{\rho} (v_2^2 - v_1^2) = \rho g (h_1 - h_2)$$

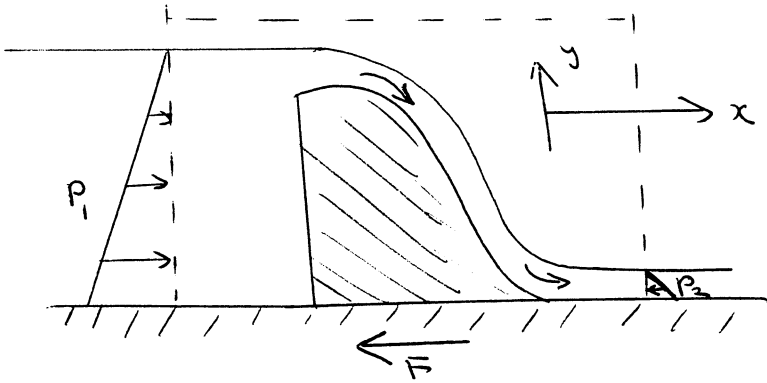
$$\Rightarrow v_2^2 \left( 1 - \left( \frac{h_2}{h_1} \right)^2 \right) = 2g (h_1 - h_2)$$

-

$$\underbrace{\hspace{10em}}_{\text{from } v_1 h_1 = v_2 h_2}$$

$$\Rightarrow \left\{ \begin{array}{l} \underline{\underline{v_2 = 9.28 \text{ m/s}}} \\ \underline{\underline{v_1 = 1.30 \text{ m/s}}} \end{array} \right.$$

- (b) Streamlines are straight and  $\parallel^e$  so no vertical acceleration in fluid  $\Rightarrow$  local force balance in  $z$ -direction is self-weight versus pressure gradient i.e. hydrostatic.



Apply force momentum in  $x$ -direction. Ignore shear on bed.

$$-F + \underbrace{\int_0^{h_1} p_1 dy}_0 - \underbrace{\int_0^{h_2} p_2 dy}_0 = \underbrace{(\rho h_1 v_1)}_m v_2 - \underbrace{(\rho h_1 v_1)}_m v_1$$

net external force.

$$\Rightarrow -F + \frac{1}{2} \rho g h_1^2 - \frac{1}{2} \rho g h_2^2 = \rho h_1 v_1 (v_2 - v_1)$$

$$\Rightarrow \underline{\underline{F = \frac{1}{2} \rho g (h_1^2 - h_2^2) - \rho h_1 v_1 (v_2 - v_1)}}$$

$$\Rightarrow \underline{\underline{F = 68.44 \times 10^3 \text{ N}}}$$

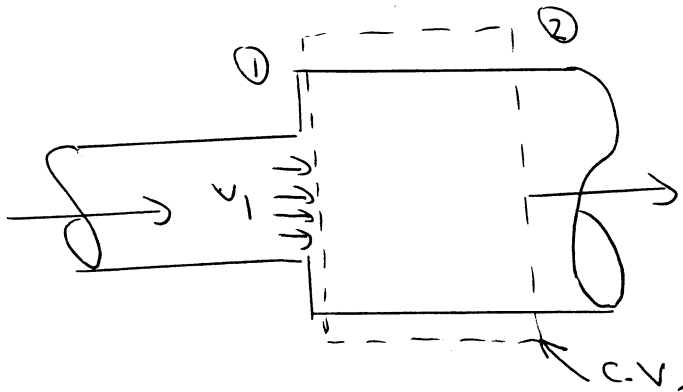
- (c) Due to friction  $v_2$  will be smaller than expected and so  $F$  will be larger.

113.

$$(a) \underbrace{\left( \frac{1}{2} \rho v_1^2 + \rho g y_1 + P_1 \right)}_{c_1} - \underbrace{\left( \frac{1}{2} \rho v_2^2 + \rho g y_2 + P_2 \right)}_{c_2}$$

$$= \underbrace{\sum \Delta C}_{\substack{\text{energy losses} \\ \text{(per unit volume)}}} - \underbrace{\sum \Delta C}_{\substack{\text{energy in from pumps} \\ \text{(per unit volume)}}$$

(b)



Apply F-M-E in streamwise direction. Ignore shear on boundary.

$$\underbrace{P_1 A_2 - P_2 A_2}_{\text{net external force}} = \dot{m} v_2 - \dot{m} v_1 = \rho v_2 A_2 (v_2 - v_1)$$

$$\Rightarrow \underline{\underline{P_1 - P_2 = \rho v_2 (v_2 - v_1)}}$$

(c) Extended Bernoulli :

$$\left( \frac{1}{2} \rho v_1^2 + P_1 \right) - \left( \frac{1}{2} \rho v_2^2 + P_2 \right) = \Delta C_{\text{loss}}$$

$$\begin{aligned} \Rightarrow |\Delta C| &= c_1 - c_2 = \frac{1}{2} \rho v_1^2 - \frac{1}{2} \rho v_2^2 + \rho v_2^2 - \rho v_2 v_1 \\ &= \frac{1}{2} \rho v_1^2 + \frac{1}{2} \rho v_2^2 + \rho v_1 v_2 \\ &= \frac{1}{2} \rho (v_1 - v_2)^2 \end{aligned}$$

$$\begin{aligned}
 \text{(d)} \quad \text{Power} &= \dot{Q} |\Delta C| \\
 &= \dot{m} \frac{1}{2} (v_1 - v_2)^2 \\
 &= \frac{1}{2} \dot{m} v_1^2 \left(1 - \frac{A_2}{A_1}\right)^2 \\
 &= \frac{1}{2} \times 10 \text{ kg/s} \times (10 \text{ m/s})^2 \times \left(1 - \frac{1}{2}\right)^2 \\
 &= 125 \text{ Watts.}
 \end{aligned}$$

(e) Consider boundary layer - dissipation stays finite as  $v \rightarrow 0$  because  $\delta \rightarrow 0$  at a rate  $v^{-1/2}$ . Expect internal shear layers to form which become thinner as  $v \rightarrow 0$  so that dissipation stays finite.

4 (a)

$$\Delta p = f(\rho, v, \mu, L, d)$$

$$\text{Parameters} = 6, \quad \text{Dimension} = 3 \quad (M, L, T)$$

$$\text{Groups} = 3$$

$$\text{Spec: } \pi_1 = \frac{\Delta p}{\frac{1}{2}\rho v^2}, \quad \pi_2 = \frac{L}{d}, \quad \pi_3 = \frac{\rho v d}{\mu}$$

$$\Rightarrow \pi_1 = f(\pi_2, \pi_3)$$

$$\Rightarrow \Delta p = \frac{1}{2}\rho v^2 f\left(\frac{L}{d}, \frac{\rho v d}{\mu}\right)$$

But  $\Delta p \propto L$  so  $f$  must be such that,

$$\Delta p = \frac{1}{2}\rho v^2 \frac{L}{d} f\left(\frac{\rho v d}{\mu}\right)$$

          
Darcy friction factor

Laminar flow,  $\Delta p \neq f(\rho)$  so only possibility is

$$(Re) = \frac{k}{\rho v d} \text{ so } \rho \text{ cancels. Thus,}$$

$$\Delta p = \frac{1}{2} \cancel{\rho} v^2 \frac{L}{d} \frac{k\cancel{\mu}}{\cancel{\rho} v d}$$

$$= \frac{k}{2} \frac{\mu v L}{d^2}$$

$$(b) \quad \Delta p = f(\rho, v, \mu, L, d, \epsilon)$$

$$\text{Parameters} = 7, \quad \text{Dimensions} = 3$$

$$\Rightarrow \text{Groups} = 4$$

$$\pi_1 = \frac{\Delta p}{\frac{1}{2}\rho v^2}, \quad \pi_2 = \frac{L}{d}, \quad \pi_3 = \frac{\rho v d}{\mu}$$

(as before)

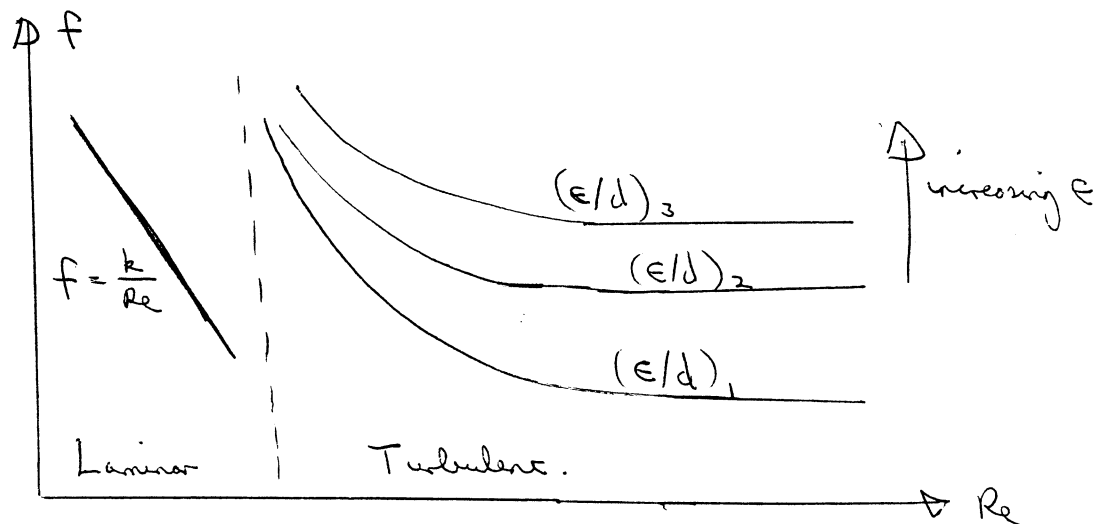
$$\text{New group: } \pi_4 = \frac{\epsilon}{d}$$

$$\pi_1 = f(\pi_2, \pi_3, \pi_4)$$

$$\Rightarrow \Delta p = \frac{1}{2}\rho v^2 f\left(\frac{L}{d}, \frac{\rho v d}{\mu}, \frac{\epsilon}{d}\right)$$

But  $\Delta p \propto L$  so as before this becomes,

$$\Delta p = \frac{1}{2}\rho v^2 \frac{L}{d} f\left(Re, \frac{\epsilon}{d}\right)$$







Heat flux by conduction along the metal per unit length along  $z$ :

$$q_c = -\lambda_m e \frac{dT}{dx}$$

Convection heat flux on both sides per unit length along  $z$

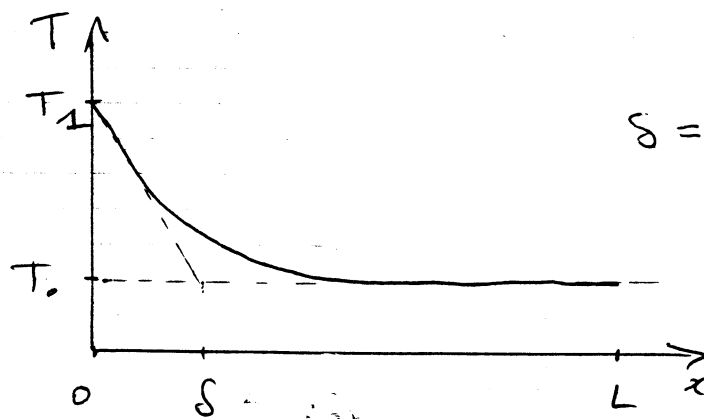
$$q_{out} = 2h dx (T - T_0)$$

Heat conservation :  $\frac{dq_c}{dx} = -q_{out}$

$$\therefore \frac{d^2T}{dx^2} = \frac{2h}{\lambda_m e} (T - T_0)$$

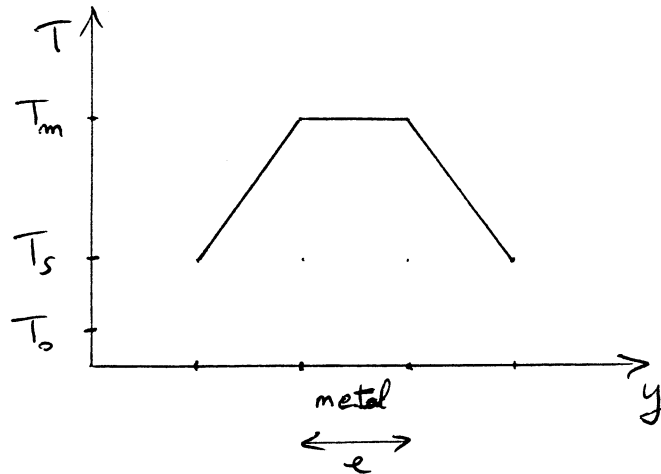
The solution is a linear combination of the two fundamental exponential solutions, but the growing exponential is set to zero.

$$T - T_0 = (T_i - T_0) e^{-\sqrt{\frac{2h}{\lambda_m e}} x} = (T_i - T_0) e^{-x/\delta}, \quad \delta = \sqrt{\frac{\lambda e}{2h}}$$



$$\delta = 0.14 \text{ m}$$

ii) With coating.



Heat conduction in the insulating material leads to a linear variation of temperature.

The heat transferred by conduction in the insulating coating must be transferred by convection in air

$$d_i \frac{T_m - T_s}{e} = h (T_s - T_0)$$

$$\therefore T_s = \frac{\frac{d_i T_m}{e} + h T_0}{\frac{d_i}{e} + h}$$

The convection heat on both sides becomes now:

$$q_{out} = 2h (T_s - T_0) = 2h \frac{T_m - T_0}{1 + \frac{he}{d_i}}$$

$$\frac{he}{d_i} = 10$$

iii) The rest of the analysis is identical with  $\frac{d_m e}{2h}$  being replaced by  $\frac{11 d_m e}{2h}$ . New  $\delta = \sqrt{\frac{11 \lambda e}{2h}} = 0.47 \text{ m}$ .