

Reduction of pressure takes the state into the superheated vapour region. If KE changes can be ignored then process is isenthalpic. Measuring P & T at ② then gives x_1 . Otherwise, both SFEE and mass continuity are needed (dashed line). (6)

(b) (i) $\dot{m} = \rho_2 A u_2$
 $\therefore u_2 = 0.3 / (0.0662 \times \pi + 0.075^2)$
 $\therefore u_2 = 256.5 \text{ ms}^{-1}$

INTERPOLATE
FOR SPECIFIC VOLUME
 $\Rightarrow \rho_2 = 0.06618$ (3)

(ii) Essentially, due to the throttling (reduction in pressure) the density is much higher at ①, hence velocity and kinetic energy are much less. (2)

(6) (iii) SFEE: $h_1 + \frac{1}{2}u_1^2 = h_2 + \frac{1}{2}u_2^2$

Neglect $\frac{1}{2}u_1^2$ term

$\therefore h_{f1} + x_1 h_{fg1} = h_2 + \frac{1}{2}u_2^2$

$\therefore x_1 = (2602.4 \times \frac{1}{2} \times 256.5^2 \times 10^{-3} - 417.5) / 2257.9$

$x_1 = 0.982$

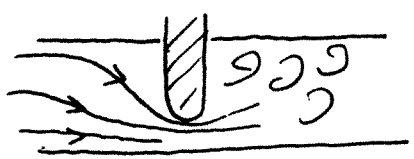
(4)

(NOTE: Including $\frac{1}{2}u_1^2$ term and solving quadratic gives $x_1 = 0.9826$)

(iv) $s_1 = (1 - x_1)^{1.302} s_{f1} + x_1^{7.36} s_{g1}$
 $= \underline{7.252 \text{ kJ/kgK}}$

$s_2 = s(0.1 \text{ bar}, 55^\circ\text{C})$
 $= \underline{8.204 \text{ kJ/kgK}}$

Increase is due to highly irreversible dissipative processes (viscous dissipation + internal heat transfer) downstream of valve:



(3)

(c) By closing valve further, mass flow rate is reduced for given pressure drop. Hence KE terms become negligible and process is isenthalpic.

(2)

Q.2. (a) Exhaust gas may be assumed to have (at least approximately) constant heat capacity. Hence exhaust gas temperature falls linearly with % of heat transferred. Water / steam temperature remains constant between P and Q due to boiling at constant pressure.

$$\begin{aligned}
 (b) \quad (i) \quad T_P &= T_S(25 \text{ bar}) + 20 \\
 &= 223.9 + 20 \\
 \underline{T_P} &= \underline{243.9 \text{ }^\circ\text{C}}.
 \end{aligned}$$

[2]

$$(ii) \quad \dot{m}_g C_{pg} (650 - 243.9) = \dot{m}_s (h(25, 450) - h_f(25))$$

$$\begin{aligned}
 \therefore \dot{m}_s &= 50 \times 1.01 \times (650 - 243.9) / (3351.25 - 962.0) \\
 \underline{\dot{m}_s} &= \underline{8.58 \text{ kg/s}}
 \end{aligned}$$

[3]

$$(iii) \quad \dot{m}_g C_{pg} (T_Q - T_P) = \dot{m}_s h_{fg}$$

$$\begin{aligned}
 \therefore T_Q &= T_P + \dot{m}_s h_{fg} / \dot{m}_g C_{pg} \\
 &= 243.9 + 8.58 \times 1839 / (50 \times 1.01) \\
 \underline{T_Q} &= \underline{556.5 \text{ }^\circ\text{C}}
 \end{aligned}$$

[2]

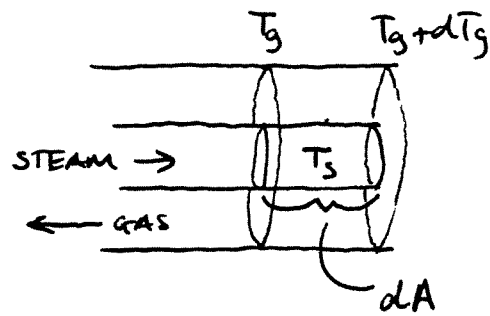
$$\begin{aligned}
 (iv) \quad \Delta \dot{S} &= \dot{m}_g C_p \ln \left(\frac{T_P}{T_Q} \right) + \dot{m}_s (s_g - s_f) \\
 &= 50 \times 1.01 \times \ln \left(\frac{243.9 + 273.15}{556.5 + 273.15} \right) \\
 &\quad + 8.58 \times (6.254 - 2.554)
 \end{aligned}$$

$$\underline{\Delta \dot{S}} = \underline{7.87 \text{ kJ/K}}.$$

[5]

Increase due to heat transfer across large temperature differences.

(c)



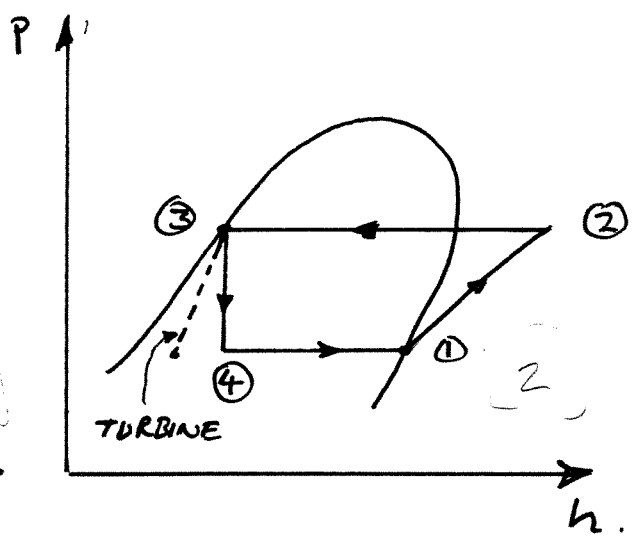
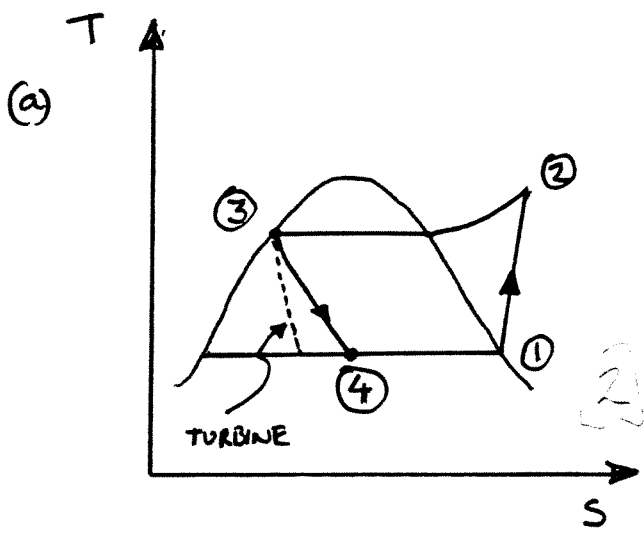
$$\dot{m}_g c_{pg} dT_g = dQ = U (T_g - T_s) dA$$

$$\therefore \frac{dT_g}{T_g - T_s} = \frac{U dA}{\dot{m}_g c_{pg}}$$

$$\begin{aligned} \therefore A &= \frac{\dot{m}_g c_{pg}}{U} \ln \left(\frac{T_g - T_s}{T_p - T_s} \right) \\ &= \frac{50 \times 1.01}{0.2} \times \ln \left(\frac{556.5 - 223.9}{20} \right) \end{aligned}$$

$$\underline{A = 710 \text{ m}^2}$$

Q3.



[2]

[2]

(b) $s_1 = 0.6991 = s_g (0.261 \text{ MN/m}^2)$
 $h_1 = 185.4 = h_{g1}$

$$h_{2s} = 199.6 + \frac{(0.6991 - 0.6854)}{(0.7321 - 0.6854)} \times (214.3 - 199.6)$$

$$= 203.91 \quad [\text{i.e., superheated by } 5.87 \text{ K}]$$

$$\therefore \Delta h_c = \frac{(203.91 - 185.4)}{0.7} = 26.44 \text{ kJ/kg.}$$

$$\therefore h_2 = 211.8 \text{ kJ/kg.}$$

$$h_3 = 64.6 \text{ kJ/kg.} \quad (h_3 = h_f)$$

$$\therefore \Delta h_{\text{cond}} = 147.24 \text{ kJ/kg}$$

[5]

Mass flow rate, \dot{m} = $\frac{\dot{Q}}{\Delta h_{\text{cond}}} = \frac{5}{147.24}$
 $= \underline{\underline{0.034 \text{ kg/s.}}}$

Power input, \dot{w}_c = $\frac{\Delta h_c}{\Delta h_{\text{cond}}} \times \dot{Q} = \frac{26.44}{147.24} \times 5000$
 $= \underline{\underline{897 \text{ W}}} \quad (= \dot{m} \Delta h_{\text{comp}})$

[2]

$$(c) \quad \text{PER} = \frac{\Delta h_{\text{COND}}}{\Delta h_c} = \underline{\underline{5.58}}$$

$$\text{PER}_{\text{IDEAL}} = \frac{T_c}{T_c - T_E} = \frac{303}{303 - 268} = \underline{\underline{8.65}}$$

Differences: (i) Irreversibility in compressor and throttle

(ii) Temperature at inlet to condenser goes above 303 K. [3]

$$s_3 = 0.2399$$

$$(d) \quad h_{4s} = 31.4 + \frac{(0.2399 - 0.1251) \cdot (185.4 - 31.4)}{(0.6991 - 0.1251)}$$

$$= 62.2 \quad [20\% \text{ liquid}] \text{ vapour} / 80\% \text{ liquid}$$

$$\therefore \Delta h_s = 64.6 - 62.2$$

$$= 2.4 \text{ kJ/kg}$$

$$\therefore \Delta h_{\text{TURB}} = 0.85 \times 2.4$$

$$= 2.04 \text{ kJ/kg}$$

$$\begin{aligned} \text{Net work input now} &= w_c - w_T \\ &= 26.44 - 2.04 \\ &= 24.4 \text{ kJ/kg} \end{aligned}$$

Heat output unchanged.

$$\therefore \text{New PER} = \frac{147.24}{24.4} = \underline{\underline{6.034}} \quad (\text{up by } 8\%)$$

[4]

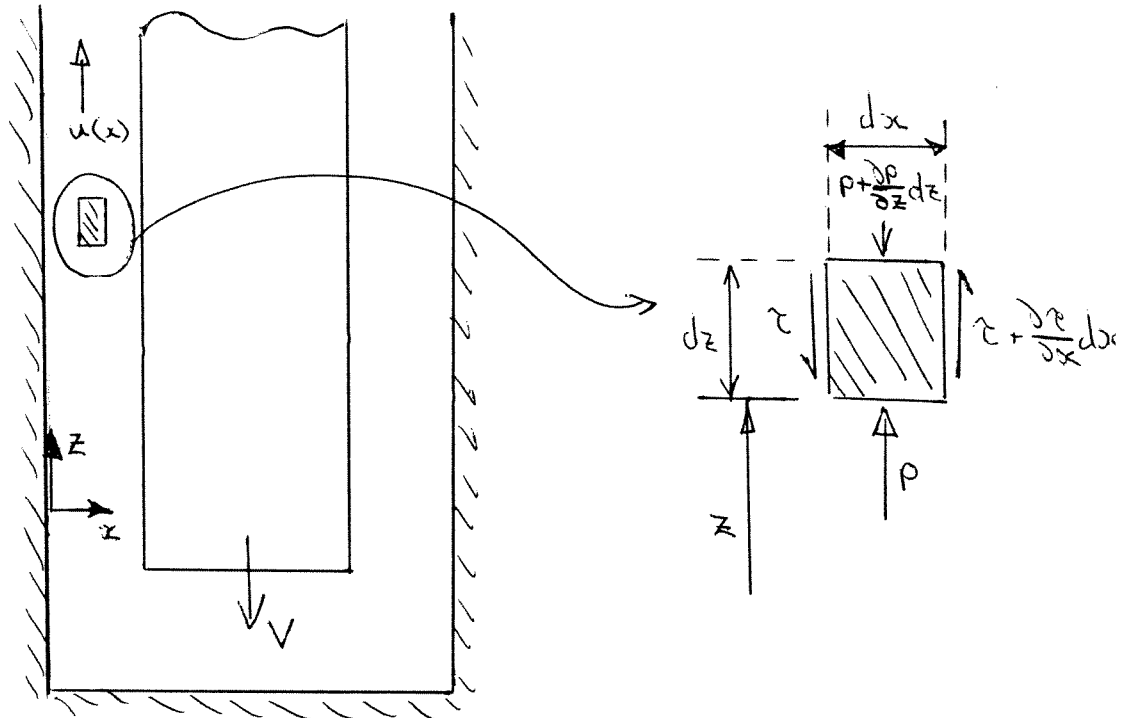
(e) Probably not: (i) Much greater complexity and cost for relatively small benefit.

(ii) High liquid content at turbine exit gives practical difficulties.

[2]

4.

(a)



Force balance on non-accelerating element of fluid. Sum forces in z direction:

$$0 = p dx - (p + \frac{dp}{dz} dz) dx - \tau dz + (\tau + \frac{d\tau}{dx} dx) dz$$

$$\Rightarrow 0 = - \frac{dp}{dz} dz dx + \frac{d\tau}{dx} dx dz$$

$$\Rightarrow \frac{d\tau}{dx} = \frac{dp}{dz}$$

Newton's law of viscosity $\tau = \mu \frac{\partial u_z}{\partial x}$

$$\Rightarrow \underline{\underline{\mu \frac{\partial^2 u_z}{\partial x^2} = \frac{dp}{dz}}}$$

(b) Integrate, noting ρ a function of z only.

$$u_z = \frac{1}{\mu} \frac{dp}{dz} \left(\frac{x^2}{2} \right) + Ax + B$$

Boundary conditions:

$$\begin{cases} u_z = 0 & \text{at } x = 0 \Rightarrow B = 0 \\ u_z = -V & \text{at } x = W \Rightarrow -V = \frac{W^2}{2\mu} \frac{dp}{dz} + AW \end{cases}$$

Substitute for A :

$$\underline{u_z = -\frac{1}{2\mu} \frac{dp}{dz} (Wx - x^2) - V \frac{x}{W}}$$

(c) Integrate to find mass flux across gap.

$$\begin{aligned} \dot{m}/\rho &= \int_0^W u_z dx = -\frac{1}{2\mu} \frac{dp}{dz} \left(\frac{W^3}{2} - \frac{W^3}{3} \right) - V \frac{W}{2} \\ &= -\frac{1}{12\mu} \frac{dp}{dz} W^3 - \frac{VW}{2} \end{aligned}$$

But continuity of mass gives: $2 \frac{\dot{m}}{\rho} = 2WV \Rightarrow \frac{\dot{m}}{\rho} = WV$

Thus,

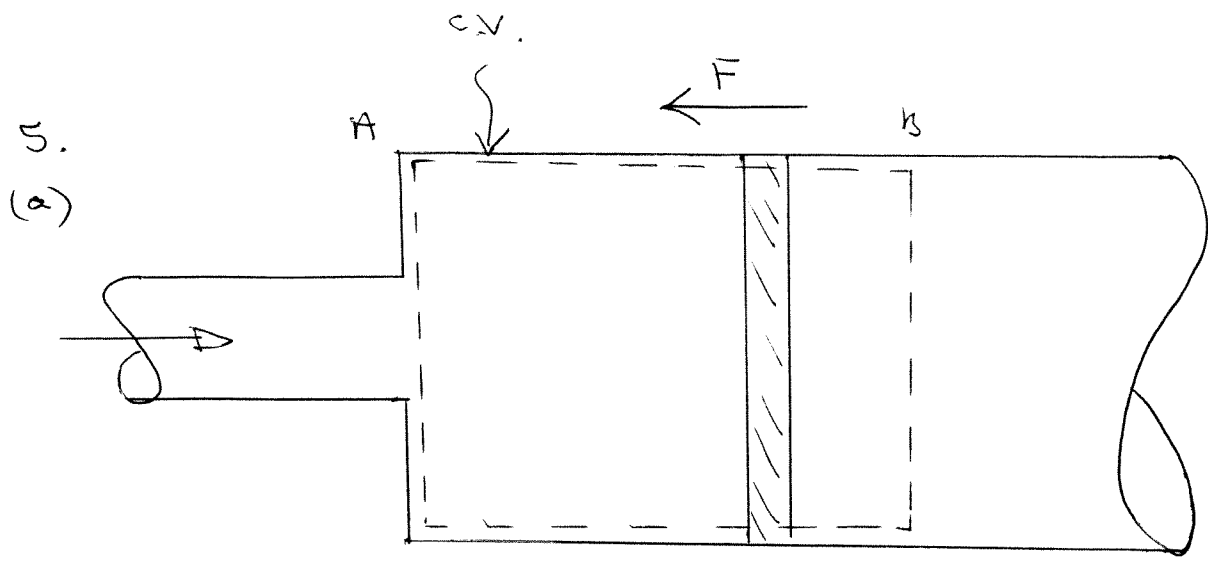
$$-\frac{1}{12\mu} \frac{dp}{dz} W^3 - \frac{VW}{2} = VW$$

$$\Rightarrow \frac{1}{\mu} \frac{dp}{dz} = -18 \frac{V}{W^2} \quad (\text{negative, as required})$$

Substitute for dp/dz :

$$u_z = 9 \frac{V}{W^2} (Wx - x^2) - V \frac{x}{W}$$

$$\Rightarrow \underline{u_z = \frac{V}{W^2} (8Wx - 9x^2)}$$



Apply F.M.E. to C.V. in streamwise direction

$$P_A A_B - P_B A_B - F = \dot{m} V_B - \dot{m} V_A$$

↑
damping force needed to keep filter in place

$$F = c_0 \left(\frac{1}{2} \rho V_B^2 \right) A_B, \quad \dot{m} = \rho A_B V_B$$

$$\Rightarrow P_A A_B - P_B A_B - c_0 \left(\frac{1}{2} \rho V_B^2 \right) A_B = \rho A_B V_B (V_B - V_A)$$

$$\Rightarrow \underline{\underline{P_A - P_B = c_0 \left(\frac{1}{2} \rho V_B^2 \right) + \rho V_B (V_B - V_A)}}$$

(b) Apply extended Bernoulli equation:

$$\left[P_A + \frac{1}{2} \rho V_A^2 \right] - \left[P_B + \frac{1}{2} \rho V_B^2 \right] = \text{energy loss (A} \rightarrow \text{B)}$$

$$\begin{aligned} \Rightarrow \text{energy loss} &= c_0 \left(\frac{1}{2} \rho V_B^2 \right) + \rho V_B (V_B - V_A) + \frac{1}{2} \rho (V_A^2 - V_B^2) \\ &= \frac{1}{2} c_0 \rho V_B^2 + \frac{1}{2} \rho (V_A^2 - V_B^2 - 2V_A V_B + 2V_B^2) \\ &= \frac{1}{2} c_0 \rho V_B^2 + \frac{1}{2} \rho (V_A^2 - 2V_A V_B + V_B^2) \\ &= \underline{\underline{\frac{1}{2} c_0 \rho V_B^2 + \frac{1}{2} \rho (V_A - V_B)^2}} \end{aligned}$$

(c) energy lost = $\frac{1}{2} \rho v_B^2 + \frac{1}{2} \rho (v_A - v_B)^2 = \frac{1}{2} \rho (\frac{1}{2} v_A)^2 + \frac{1}{2} \rho (\frac{1}{2} v_A)^2$
 ($c_0 = 1$) ($v_B = \frac{1}{2} v_A$)

\Rightarrow energy lost = $\frac{1}{4} \rho v_A^2$

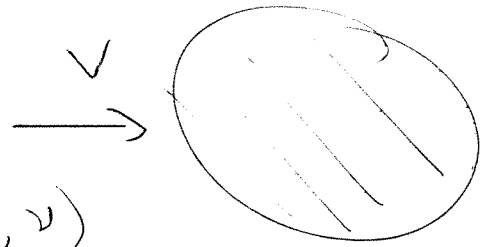
This loss of mechanical energy (per unit volume) leads to a rise of internal energy so,

~~$\rho c_v \Delta T = \frac{1}{4} \rho v_A^2$~~

$\Rightarrow \Delta T = \frac{v_A^2}{4 c_v} = \underline{\underline{0.1^\circ C}}$

6.

(a)



$$\dot{Q} = f(\Delta T, V, d, \lambda, \alpha, \nu)$$

Π theorem: Parameters = 7

Dimensions = 4 (kg, s, m, °C)

Number of groups = 7 - 4 = 3.

Spot check

- { ν / α is dimensionless (Prandtl number)
- { Vd / ν is dimensionless (Reynolds number)

Need one more group. It must involve \dot{Q} , ΔT and λ .

$$\dot{Q} \equiv \text{Watts/m}$$

$$\Delta T \equiv ^\circ\text{C}$$

$$\lambda \equiv \text{Watts/m/}^\circ\text{C}$$

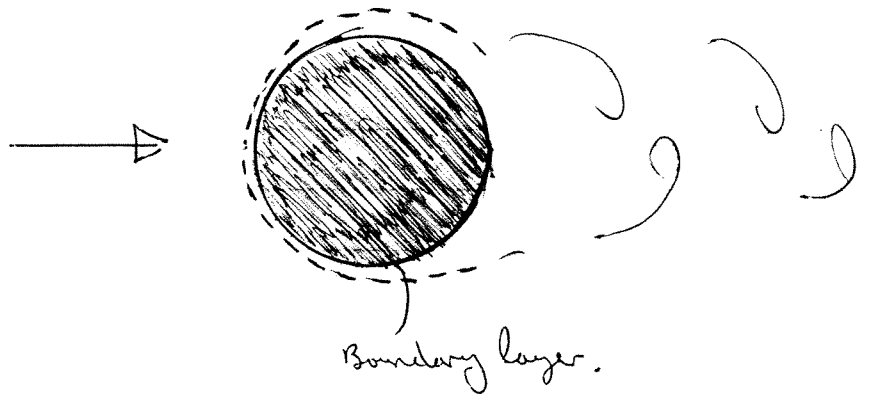
This third group is $\frac{\dot{Q}}{\lambda \Delta T}$ or $\frac{\lambda \Delta T}{\dot{Q}}$.

In summary,

$$\underline{\underline{\frac{\dot{Q}}{\lambda \Delta T} = f\left(\frac{Vd}{\nu}, \frac{\nu}{\alpha}\right)}}$$

This ΔT is proportional to \dot{Q} .

(b)



The boundary layer represents a significant part of the thermal resistance. The larger v , the thicker the boundary layer, and hence the higher the thermal resistance.

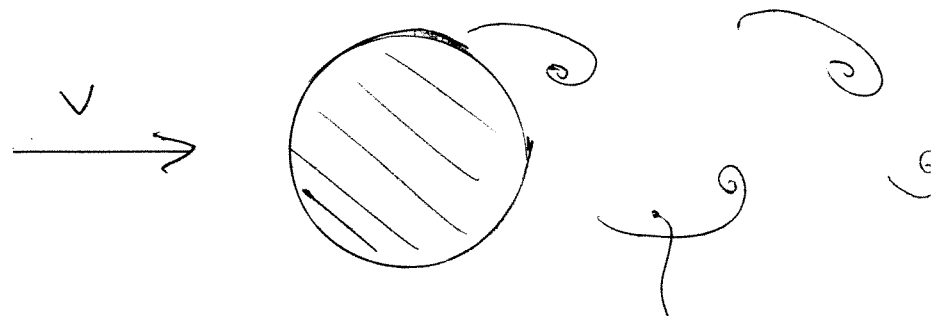
(c) If v is not a relevant parameter,

$$\frac{\dot{Q}}{\lambda \Delta T} = f\left(\frac{vd}{\alpha}\right)$$

↑ Peclet number.

If $\dot{Q}/\lambda \Delta T$ is proportional to $v^{0.4}$ then it is also proportional to $d^{0.4}$ and $\alpha^{-0.4}$.

(d)



Kármán vortex street

The oscillations are due to the periodic shedding of vortices from the rear of the cylinder which carry heat downstream.