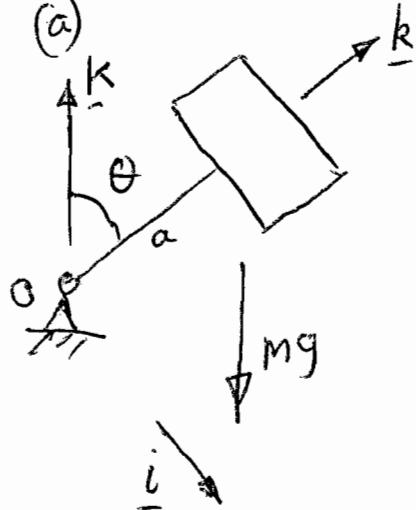


1. (a)



The only couple acting is due to mg

$$\underline{Q} = m g a \sin \theta \underline{j} \quad (1)$$

Gyro equation #3

$$Q_3 = C \dot{\omega}_3 \quad \text{but } (1) \rightarrow Q_3 = 0$$

$$\therefore \dot{\omega}_3 = 0 \quad \therefore \omega_3 = \text{const}$$

$$(b) \quad \underline{Q} = \underline{h} \quad \therefore \underline{\dot{Q}} \cdot \underline{k} = \underline{\dot{h}} \cdot \underline{k}$$

$$\text{and from (1)} \quad \underline{Q} \cdot \underline{k} = 0 \quad \therefore \underline{\dot{h}} \cdot \underline{k} = 0$$

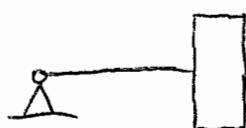
$$\text{but } \underline{k} \text{ is a constant vector} \quad \therefore \frac{d}{dt} (\underline{h} \cdot \underline{k}) = 0$$

$$\therefore \underline{h} \cdot \underline{k} = \text{const}$$

$$\text{Initial value} = C \omega$$

(c)

Gyro equation #2



$$Q_2 = A \dot{\omega}_2 + (A \underline{R}_3 - C \omega_3) \underline{R}_1$$

$$\text{where } \underline{R}_1 = -\dot{\phi} \sin \theta \underline{i}$$

$$\underline{R}_2 = \dot{\theta} \underline{j}$$

$$\underline{R}_3 = \dot{\phi} \cos \theta \underline{k}$$

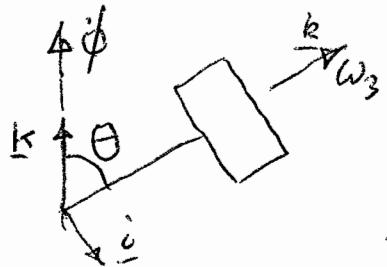
$$\text{Steady state precession} \quad \therefore \dot{\theta} = 0$$

$$\therefore m g a \sin \theta = -(A \dot{\phi} \cos \theta - C \omega_3) \dot{\phi} \sin \theta \quad (2)$$

$$\text{and put } \theta = \frac{\pi}{2} \quad \therefore m g a = C \omega_3 \dot{\phi}$$

$$\therefore \dot{\phi}_{\text{final}} = \frac{m g a}{C \omega_3} \quad \underline{\text{ans}}$$

I(c) cont. Now prove $\Theta = \frac{\pi}{2}$ by conservation of $\underline{h} \cdot \underline{k}$



$$\underline{h} = A \underline{R}_1 \underline{i} + A \underline{R}_2 \underline{j} + C \underline{\omega}_3 \underline{k}$$

$$\underline{h} \cdot \underline{k} = A \underline{R}_1 (\underline{i} \cdot \underline{k}) + C \underline{\omega}_3 (\underline{k} \cdot \underline{k})$$

$$\underline{i} \cdot \underline{k} = -\sin \Theta$$

$$\underline{k} \cdot \underline{k} = \cos \Theta \quad \therefore C \underline{\omega}_3 = +A \dot{\phi} \sin \Theta \sin \Theta + C \underline{\omega}_3 \cos \Theta$$

$$\begin{aligned} \therefore C \underline{\omega}_3 (1 - \cos \Theta) &= A \dot{\phi} \sin^2 \Theta \\ &= A \dot{\phi}^2 (1 - \cos^2 \Theta) \\ &= A \dot{\phi}^2 (1 - \cos \Theta)(1 + \cos \Theta) \end{aligned}$$

$$\therefore C \underline{\omega}_3 = A \dot{\phi}^2 (1 + \cos \Theta) \quad (3)$$

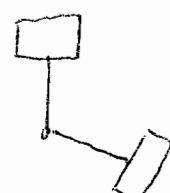
and with $\dot{\phi}_{\text{final}} = \frac{mga}{C \underline{\omega}_3}$, $\Theta_{\text{final}} = \frac{\pi}{2}$

$$\therefore C \underline{\omega}_3 = \frac{A m g a}{C \underline{\omega}_3} \quad *$$

$$\text{but } (C \underline{\omega}_3)^2 = (C \underline{\omega})^2 = \left(C \frac{\underline{\omega}_{\text{init}}}{2}\right)^2 = m g a A \quad (4)$$

So * is satisfied.

$$(d) (KE + PE)_{\text{initial}} = (KE + PE)_{\text{final}}$$



$$\therefore \frac{1}{2} C \underline{\omega}_3^2 + m g a = \frac{1}{2} A \underline{\omega}_1^2 + \frac{1}{2} A \underline{\omega}_2^2 + \frac{1}{2} C \underline{\omega}_3^2 + m g a \cos \Theta$$

$$\text{but at } \Theta_{\text{max}}, \underline{\omega}_2 = 0 \quad \text{and} \quad \underline{\omega}_1^2 = \dot{\phi}^2 \sin^2 \Theta$$

$$\begin{aligned} \therefore m g a (1 - \cos \Theta) &= \frac{1}{2} A \dot{\phi}^2 \sin^2 \Theta \\ &= \frac{1}{2} A \dot{\phi}^2 (1 - \cos^2 \Theta) \end{aligned}$$

$$\therefore m g a = \frac{1}{2} A \dot{\phi}^2 (1 + \cos \Theta)$$

$$\text{but } ③ \rightarrow \dot{\phi}^2 = \left(\frac{c\omega_3}{A(1+\cos\Theta)} \right)^2$$

$$\therefore mga = \frac{1}{2} A \frac{(c\omega_3)^2}{A^2 (1+\cos\Theta)}$$

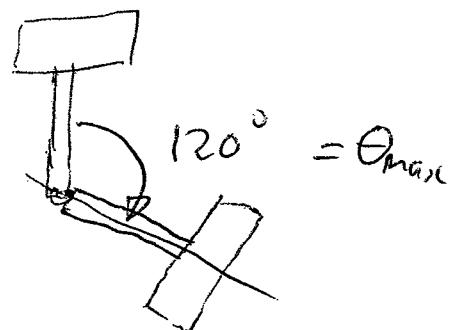
$$\text{and from } ④ \quad = \frac{1}{2A} \frac{mga A}{1+\cos\Theta}$$

$$\therefore 1 + \cos\Theta = \frac{1}{2}$$

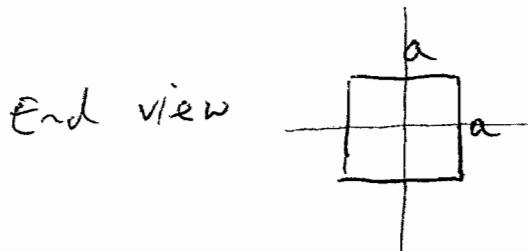
$$\therefore \cos\Theta = -\frac{1}{2}$$

$$\therefore \Theta = \frac{2\pi}{3}$$

=====



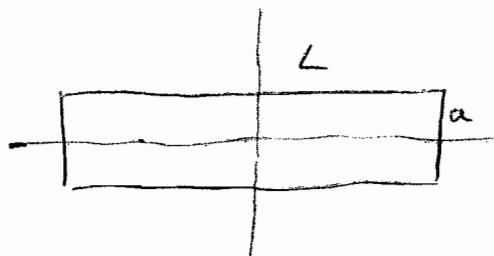
2. (a)



$$\therefore C = \frac{1}{12}Ma^2 + \frac{1}{12}Ma^2 = \underline{\underline{\frac{1}{6}Ma^2}}$$

by b axis theorem &
data book

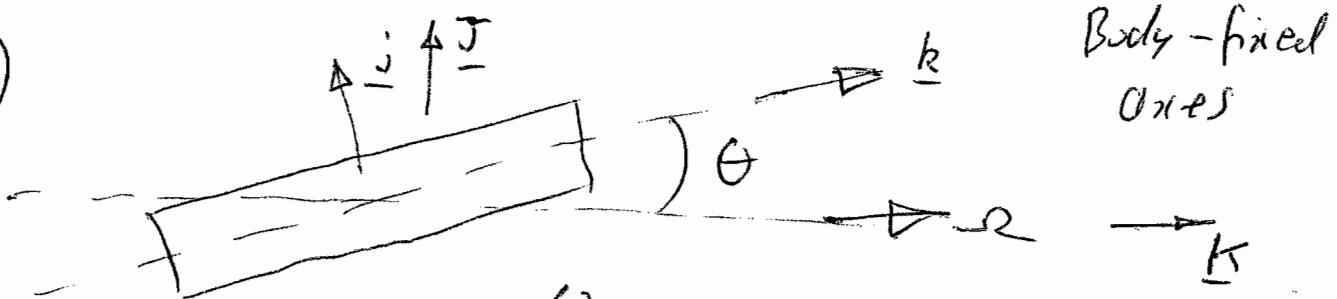
Side view



$$A = \frac{1}{12}ma^2 + \frac{1}{12}mc^2$$
$$= \underline{\underline{\frac{1}{12}m(a^2 + c^2)}}$$

The body is an AAC body so
behaves like a cylinder.

(b)



$$\omega_1 = 0$$

$$\omega_2 = -\underline{R} \sin \theta$$

$$\omega_3 = \underline{R} \cos \theta$$

$$\therefore \underline{h} = A\omega_1 \underline{i} + A\omega_2 \underline{j} + C\omega_3 \underline{k}$$
$$= -A\underline{R} \sin \theta \underline{i} + C\underline{R} \cos \theta \underline{k}$$

and for \underline{h} resolve \underline{h} into $\underline{i} \underline{j} \underline{k}$

$$\begin{aligned}
 2(b) \text{ ii ant. } \underline{h} \cdot \underline{k} &= -A\omega \sin\theta (\underline{i} \cdot \underline{k}) + c\omega \cos\theta (\underline{k} \cdot \underline{k}) \\
 &= -A\omega \sin\theta (-\sin\theta) + c\omega \cos\theta (\cos\theta) \\
 &= A\omega \sin^2\theta + c\omega \cos^2\theta
 \end{aligned}$$

$$\begin{aligned}
 \underline{h} \cdot \underline{j} &= -A\omega \sin\theta (\underline{i} \cdot \underline{j}) + c\omega \cos\theta (\underline{k} \cdot \underline{j}) \\
 &= -A\omega \sin\theta (\cos\theta) + c\omega \cos\theta (\sin\theta) \\
 &= (c-A)\omega \cos\theta \sin\theta
 \end{aligned}$$

$$\underline{h} \cdot \underline{i} = 0$$

$$\therefore \underline{h} = \underbrace{(c-A)\omega \cos\theta \sin\theta \underline{j}}_{\text{ans}} + \underbrace{(A\omega \sin^2\theta + c\omega \cos^2\theta) \underline{k}}_{\text{ans}} \quad (1)$$

$$\text{Now, } \underline{h} = \begin{bmatrix} \text{Inertia} \\ \text{Matrix} \end{bmatrix} \underline{\omega} \underline{k}$$

$$= \begin{bmatrix} \ddots & -e \\ \ddots & -f \\ \ddots & g \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ -\omega \end{bmatrix}$$

$$= e\omega \underline{i} \cancel{+} f\omega \underline{j} \cancel{+} g\omega \underline{k}$$

Compare with (1)

$$\therefore f = -(c-A)\omega \cos\theta \sin\theta$$

which is one of the products
of inertia

(c) Use Euler's equations in body-fixed axes i j k

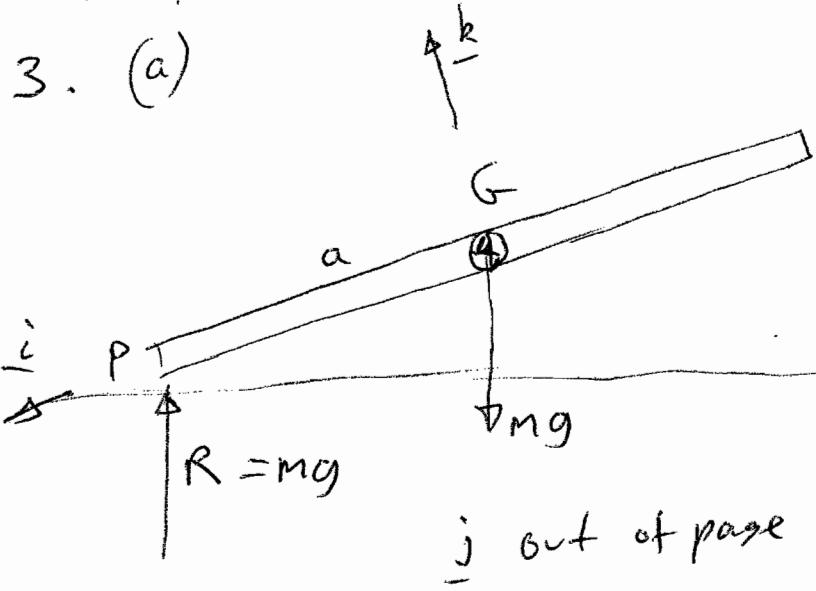
$$\begin{aligned} Q_1 &= A\dot{\omega}_1 - (A-C)\omega_2\omega_3 \\ &= 0 - (A-C)(-\omega \sin \theta)(\omega \cos \theta) \end{aligned}$$

$$Q_1 = \underline{(A-C) \omega^2 \sin \theta \cos \theta} \quad \text{ans}$$

(d) $Q_1 = 0$ if $A=C$

$$\therefore a=L$$

This is when the rotor is a cube
which is an AAA body, identical
to a sphere which can spin
about any axis with no
dynamic moment.



G is at rest

$$\therefore R = mg$$

(no net force on
the coin)

$$Q = -mga \cos\theta \quad \text{factoring moments about G}$$

(b) Gyro equation "2"

$$Q_2 = A\dot{\alpha}_2 + (A\dot{\alpha}_3 - \omega_3)\alpha_1$$

$$\text{with } \alpha_1 = -\dot{\phi} \sin\theta$$

$$\alpha_2 = \dot{\phi} = 0 \quad \text{instead state}$$

$$\alpha_3 = \dot{\phi} \cos\theta$$

$$\therefore -mga \cos\theta = (\dot{\phi} \cos\theta - \omega_3)(-\dot{\phi} \sin\theta) \quad (1)$$

No slip at P:

$$\omega \times a_i = 0 \quad \therefore (\alpha_{1i} + \alpha_{2i} + \omega_3 k) \times a_i = 0$$

$$\therefore \alpha_2 = 0$$

$$\text{and } \omega_3 = 0$$

$$3(b) \text{ cont. so } (1) \rightarrow m g a \cos \theta = A \phi \sin \theta$$

$$\therefore \dot{\phi}^2 = \frac{m g a}{A} \cot \theta$$

$$\text{but for a thin coin } C = \frac{1}{2} m a^2$$

$$A = \frac{1}{4} m a^2$$

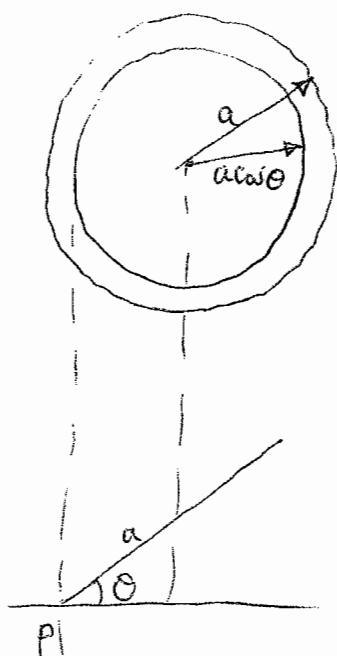
$$\therefore \dot{\phi}^2 = \frac{4 m g a \cot \theta}{m a^2} = 4 \frac{g}{a} \cot \theta$$

$$\therefore \dot{\phi} = \underline{2 \sqrt{\frac{g}{a} \cot \theta}} \quad \text{ans}$$

This is the rate of turning of the reference frame
ie the wobbling rate

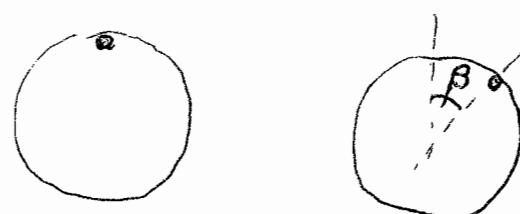
3(c) For the turning rate of the head of the coin
consider the circle on which point P rolls.

In one wobble the circumference
of length $2\pi a$ rolls on a circle
of circumference $2\pi a \cos \theta$



So a spot on the edge of the coin
has advanced a distance

$$2\pi a - 2\pi a \cos \theta$$



after one wobble

So the apparent angle β through which the coin has turned is $\frac{2\pi a(1 - \cos\theta)}{a}$

The denominator could equally be $a \cos\theta$, but at this stage assume small θ

$$\therefore \cos\theta \approx 1 - \frac{1}{2}\theta^2$$

Ignore the θ^2 bit in the denominator, but not in the numerator

$$\therefore \beta \approx 2\pi \frac{1}{2}\theta^2 = \pi\theta^2$$

The time taken for one wobble of the coin

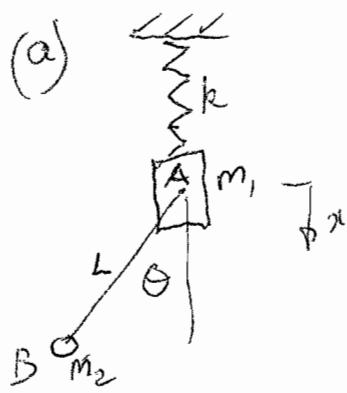
$$T = \frac{2\pi}{\dot{\phi}} = \frac{2\pi}{2\sqrt{\frac{g}{a\theta}}} \quad \text{again for small } \theta$$

$$= \pi \sqrt{\frac{a\theta}{g}} \quad \cot\theta \approx \frac{1}{\theta}$$

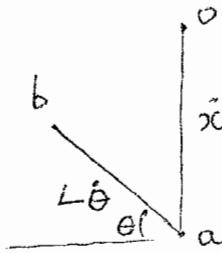
$$\text{So } \dot{\beta} = \frac{\beta}{T} = \frac{\pi\theta^2}{\pi\sqrt{\frac{a\theta}{g}}} = \sqrt{\frac{g\theta^3}{a}} \quad \text{ans}$$

This is the rate of turning of the head on the coin. $\dot{\beta} \rightarrow 0$ as $\theta \rightarrow 0$
as observed.

4. (a)



Velocity diagram



$$T = \frac{1}{2} m_1 \dot{x}^2 + \frac{1}{2} m_2 \left[(\dot{x} - L\dot{\theta} \sin \theta)^2 + (L\dot{\theta} \cos \theta)^2 \right]$$

$$= \frac{1}{2} (m_1 + m_2) \dot{x}^2 + \frac{1}{2} m_2 (2\dot{x}L\dot{\theta} \sin \theta) + \frac{1}{2} m_2 L^2 \dot{\theta}^2 (\sin^2 \theta + \cos^2 \theta)$$

$$\underline{T = \frac{1}{2} [(m_1 + m_2) \dot{x}^2 + m_2 (L^2 \dot{\theta}^2 - 2\dot{x}L\dot{\theta} \sin \theta)]} \quad \underline{\text{ans}}$$

$$V = \underline{\underline{\frac{1}{2} k x^2}} - m_1 g x - m_2 g (x - \frac{1}{2} L(1 - \cos \theta)) \quad \underline{\text{ans}}$$

(b) x -equation

$$\boxed{\text{Lagrange: } \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{x}_j} \right) - \frac{\partial T}{\partial x_j} + \frac{\partial V}{\partial x_j} = Q_j}$$

$$\frac{\partial T}{\partial x} = (m_1 + m_2) \dot{x} - m_2 L \dot{\theta} \sin \theta$$

$$\frac{\partial T}{\partial \dot{x}} = 0$$

$$\frac{\partial V}{\partial x} = kx - (m_1 + m_2)g$$

$$\therefore \underline{(m_1 + m_2) \ddot{x} - m_2 L \ddot{\theta} \sin \theta - m_2 L \dot{\theta}^2 \sin \theta \cos \theta + kx} \quad \underline{\text{exp}} \\ \underline{-(m_1 + m_2)g = 0} \quad \underline{\text{ans.}}$$

check: put $\theta = 0$

$$\therefore (m_1 + m_2) \ddot{x} + kx = (m_1 + m_2)g \quad \text{usual equation for mass on spring}$$

4(b) cont. θ equation

$$\frac{\partial T}{\partial \dot{\theta}} = m_2 L^2 \ddot{\theta} - m_2 L \dot{x} \sin \theta$$

$$\frac{\partial T}{\partial \theta} = -m_2 L \dot{x} \dot{\theta} \cos \theta$$

$$\frac{\partial V}{\partial \theta} = +m_2 g L \sin \theta$$

$$\therefore m_2 L^2 \ddot{\theta} - m_2 L \dot{x} \sin \theta - m_2 L \dot{x} \dot{\theta} \cos \theta + m_2 L \dot{x} \dot{\theta} \cos \theta + m_2 g L \sin \theta = 0$$

$$\therefore \underline{m_2 L^2 \ddot{\theta} - \dot{x} \sin \theta + g \sin \theta = 0} \quad \underline{\text{ans.}}$$

Check : put $x=0 \quad \therefore \ddot{\theta} + \frac{g}{L} \sin \theta = 0$, usual pendulum equation

$$4(c) \quad \begin{aligned} & \text{put } \sin \theta \sim \theta \\ & \cos \theta \sim 1 \quad \text{and ignore small terms.} \\ \therefore (m_1 + m_2) \ddot{x} + kx &= (m_1 + m_2) g \quad \left. \right\} \text{ans} \\ \text{and } L \ddot{\theta} + g \theta &= 0 \end{aligned}$$

4(d) Put $\sin \theta \approx \theta$ and $\cos \theta \approx 1 - \frac{\theta^2}{2}$
 into T & V , keeping up to order θ^2

$$T \approx \frac{1}{2} [(m_1 + m_2) \dot{x}^2 + m_2 L^2 \dot{\theta}^2]$$

$$V \approx \frac{1}{2} k x^2 - (m_1 + m_2) g x + \frac{1}{2} m_2 g L \theta^2$$

Make a substitution $x = y + x_0$

$$x^2 = y^2 + 2x_0 y + x_0^2$$

$$\therefore V = \frac{1}{2} k (y^2 + 2x_0 y + x_0^2) - (m_1 + m_2) g (y + x_0) + \frac{1}{2} m_2 g L \theta^2$$

$$\text{choose } x_0 = \frac{(m_1 + m_2) g}{k}$$

and arbitrarily shift datum for V to $\frac{1}{2} k x_0^2 - (m_1 + m_2) x_0$

$$\therefore V = \frac{1}{2} k y^2 + \frac{1}{2} m_2 g L \theta^2$$

and since $\dot{x} = \dot{y}$

$$T = \frac{1}{2} [(m_1 + m_2) \dot{y}^2 + m_2 L^2 \dot{\theta}^2]$$

T & V give diagonal mass and stiffness matrices

$$\{M\} = \begin{bmatrix} m_1 + m_2 & 0 \\ 0 & m_2 L^2 \end{bmatrix} \quad \text{and} \quad \{F\} = \begin{bmatrix} k & 0 \\ 0 & m_2 g L \end{bmatrix}$$

4(d) Put $\sin\theta \approx \theta$ and $\cos\theta \approx 1 - \frac{\theta^2}{2}$
 into $T \delta V$, keeping terms up to order θ^2

$$T \approx \frac{1}{2} ((m_1 + m_2) \ddot{x}^2 + m_2 L^2 \dot{\theta}^2)$$

$$V \approx \frac{1}{2} k x^2 - (m_1 + m_2) g x + \frac{1}{2} m_2 g L \theta^2$$

only quadratic terms count in the mass & stiffness
 matrices so

$$[M] = \begin{bmatrix} m_1 + m_2 & 0 \\ 0 & m_2 L^2 \end{bmatrix} \quad [k] = \begin{bmatrix} k & 0 \\ 0 & m_2 g L \end{bmatrix}$$

$$\text{and } [M] \begin{bmatrix} \ddot{x} \\ \ddot{\theta} \end{bmatrix} + [k] \begin{bmatrix} x \\ \theta \end{bmatrix} = \begin{bmatrix} (m_1 + m_2) g \\ 0 \end{bmatrix}$$



this comes from $\frac{\partial V}{\partial x} = -(m_1 + m_2)g$
 which is a constant throughout
 the motion.

S(a)

Think in terms of θ

$$T = \frac{1}{2} I \dot{\theta}^2$$

Now do a coordinate transformation

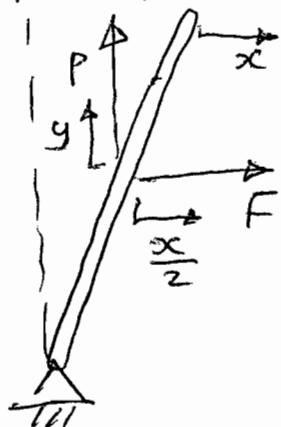
$$x = L \sin \theta$$

$$\dot{x} = L \dot{\theta} \cos \theta$$

$$\dot{x}^2 = L^2 \dot{\theta}^2 \cos^2 \theta = \dot{\theta}^2 (L^2 - L^2 \sin^2 \theta) \\ = \dot{\theta}^2 (L^2 - x^2)$$

$$\therefore T = \frac{1}{2} I \frac{\dot{x}^2}{L^2 - x^2} \quad \underline{\underline{\text{ans}}}$$

For generalized forces use a work argument



$$dW = F \frac{dx}{2} + P dy$$

↗
 easy
 by
 inspection

↗
 needs thought

$$y = \frac{L}{2} (\cos \theta - 1)$$

$$dy = \frac{L}{2} (-\sin \theta d\theta) = -\frac{x}{2} d\theta$$

$$\text{and } dx = L \cos \theta d\theta = L \sqrt{1 - \sin^2 \theta} d\theta$$

$$\therefore d\theta = \frac{dx}{\sqrt{L^2 - x^2}}$$

$$\therefore dW = Q dx = \frac{F}{2} dx - \frac{P x}{2 \sqrt{L^2 - x^2}} dx$$

$$\therefore Q = \frac{F}{2} - \frac{P x}{2 \sqrt{L^2 - x^2}} \quad \underline{\underline{\text{ans}}}$$

S(b)

$$\frac{\partial T}{\partial \dot{x}} = \frac{I \dot{x}}{L^2 - x^2}$$

$$\frac{\partial T}{\partial x} = \frac{I \dot{x}^2 x}{(L^2 - x^2)^2}$$

V=0 because all forces are included in Q

$$\text{so } \frac{\partial V}{\partial x} \approx 0$$

$$5(b) \text{ cont'd} \quad \text{Lagrange} : \frac{d}{dt} \left(\frac{I\ddot{x}}{L^2 - x^2} \right) - \frac{I\dot{x}^2 x}{(L^2 - x^2)^2} = Q$$

$$\therefore \frac{I\ddot{x}(L^2 - x^2) - I\dot{x}(-2x)}{(L^2 - x^2)^2} - \frac{I\dot{x}^2 x}{(L^2 - x^2)^2} = Q$$

$$\therefore \frac{I\ddot{x}}{L^2 - x^2} + \frac{I\dot{x}^2 x}{(L^2 - x^2)^2} = \frac{F}{2} \cancel{-} - \frac{Px}{2\sqrt{L^2 - x^2}} \quad \underline{\underline{\text{ans}}}$$

Small vibration $\therefore x \ll L$ and ignore $\dot{x}^2 x$ terms

$$\therefore \frac{I\ddot{x}}{L^2} + \frac{Px}{2L} = \frac{F}{2} \quad \therefore \ddot{x} + \frac{PL}{2I} x = \frac{FL^2}{2I}$$

$$\therefore \omega_n = \sqrt{\frac{PL}{2I}} \quad \underline{\underline{\text{ans}}} \quad \omega_n \stackrel{\uparrow}{\sim} \frac{P}{L^2}$$

$$5(c) \quad \text{use } \theta \quad \therefore T = \frac{1}{2} I \dot{\theta}^2$$

$V = 0$ as before

$$dN = Q d\theta = F \frac{dx}{2} + p dy$$

$$\begin{aligned} dx &= L \cos \theta d\theta \\ dy &= -\frac{L}{2} \sin \theta d\theta \end{aligned} \quad \left. \begin{array}{l} \\ \end{array} \right\} \text{as before}$$

$$\therefore Q = F \frac{L}{2} \cos \theta - P \frac{L}{2} \sin \theta$$

$$\frac{\delta T}{\delta \dot{\theta}} = I \dot{\theta} \quad \frac{\delta T}{\delta \theta} = 0 \quad \frac{\delta V}{\delta \theta} = 0$$

$$\therefore I \ddot{\theta} = F \frac{L}{2} \cos \theta - P \frac{L}{2} \sin \theta \quad \underline{\underline{\text{ans}}}$$

$$\left(\text{Check, small } \theta \therefore I \ddot{\theta} + \frac{PL}{2} \theta = \frac{FL}{2} \right)$$

as before with $x = L \theta$

S(c) contd

$$\text{Substitute } \ddot{x} = L \sin \theta$$

$$\dot{x} = L \dot{\theta} \cos \theta$$

$$\ddot{x} = L \ddot{\theta} \cos \theta - L \dot{\theta}^2 \sin \theta$$

into the \ddot{x} -equation from S(b)

$$\frac{I}{L^2 - \dot{x}^2} (L \ddot{\theta} \cos \theta - L \dot{\theta}^2 \sin \theta) + \frac{I(L \dot{\theta} \cos \theta)^2 L \sin \theta}{(L^2 - \dot{x}^2)^2} = \\ = \frac{F}{2} - \frac{PL \sin \theta}{2\sqrt{L^2 - (L \sin \theta)^2}}$$

$$\therefore \frac{I}{L \cos^2 \theta} (\ddot{\theta} \cos \theta - \dot{\theta}^2 \sin \theta) + \frac{I \dot{\theta}^2 \sin \theta}{L \cos^2 \theta} = \frac{F}{2} - \frac{PL \sin \theta}{2 \cos \theta}$$

$$\therefore I \ddot{\theta} = \frac{F L \cos \theta}{2} - \frac{P L \sin \theta}{2} \quad \underline{\text{as required}}$$



Notes • Several candidates obtained full marks, simply, by not making mistakes.

• Fudging the "show that" was obvious

• " $V=0$ because buoyancy forces are equal and opposite to gravity earned" no marks

Engineering Tripos Part IIA 2007

Answers

1. (c) $\dot{\phi} = \frac{mga}{C\omega} = \sqrt{\frac{mga}{A}} = \frac{C\omega}{A}$ (d) 120°

2. (a) $m(a^2 + L^2)/12, m(a^2 + L^2)/12, ma^2/6$
 (b) $\omega = 0, i - \Omega \sin \theta j + \Omega \cos \theta k$
 (i) $\mathbf{h} = -A\Omega \sin \theta j + C\Omega \cos \theta k$
 (ii) $\mathbf{h} = (C - A)\Omega \cos \theta \sin \theta \mathbf{j} + (A\Omega \sin^2 \theta + C\Omega \cos^2 \theta) \mathbf{k}$
 product of inertia = $-(C - A)\cos \theta \sin \theta$
 (c) $Q_1 = (A - C)\Omega^2 \cos \theta \sin \theta$
 (d) $L = a$

3. (a) $\mathbf{Q} = -mga \cos \theta \mathbf{j}$ (b) $2\sqrt{\frac{g}{a \tan \theta}}$ (c) $\sqrt{\frac{g\theta^3}{a}}$

4. (a) $V = \frac{1}{2}k\theta^2 - m_1gx - m_2g(x - L(1 - \cos \theta))$
 (b) $(m_1 + m_2)\ddot{x} - m_2L\ddot{\theta} \sin \theta - m_2L\dot{\theta}^2 \sin \theta \cos \theta + kx - (m_1 + m_2)g = 0$
 $L\ddot{\theta} - \ddot{x} \sin \theta + g \sin \theta = 0$
 (c) $(m_1 + m_2)\ddot{x} + kx = (m_1 + m_2)g$
 $L\ddot{\theta} + g\theta = 0$
 (d) $M = \begin{bmatrix} m_1 + m_2 & 0 \\ 0 & m_2L^2 \end{bmatrix} \quad K = \begin{bmatrix} k & 0 \\ 0 & m_2gL \end{bmatrix}$

5. (b) $\frac{I\ddot{x}}{L^2 - x^2} + \frac{I\dot{x}^2 x}{(L^2 - x^2)^2} = \frac{F}{2} - \frac{Px}{2\sqrt{L^2 - x^2}}$
 (c) $T = \frac{1}{2}I\dot{\theta}^2 \quad Q = \frac{FL \cos \theta}{2} - \frac{PL \sin \theta}{2}$