

4A12 Turbulence  
Solutions  
2003

4A12

Question 1

$$a) \nabla^2 u_y = \frac{1}{r} \frac{d}{dr} \left[ r \frac{du_y}{dr} \right] = \frac{\partial p}{\partial y} = a$$

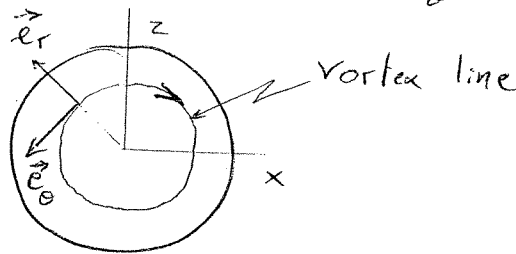
$$u_y = \frac{a}{4} \left( r^2 - \frac{D^2}{4} \right)$$

$$\bar{v} = \frac{1}{\pi D^2/4} \int_0^{D/2} 2\pi r u_y dr = -\frac{a D^2}{32}$$

hence  $u_y = -\frac{8\bar{v}}{D^2} \left( r^2 - \frac{D^2}{4} \right)$  [Given in the question]

Remark:  $Re = \frac{\bar{v} D}{\nu} = \frac{10^{-1} 10^{-2}}{10^{-6}} = 10^3$  laminar

a) Vorticity  $\omega_r = \omega_y = 0$   $\omega_z = \frac{du_y}{dr} = -16\bar{v} \frac{r}{D^2}$



[30%]

b) (i) When  $z > 0$   $\omega_x > 0$  in section A-A'  
When  $z < 0$   $\omega_x < 0$  in section A-A'

Inviscid vortex equation on a short length  $\frac{D\omega_x}{Dt} = (\vec{\omega} \cdot \vec{\nabla}) u_x$

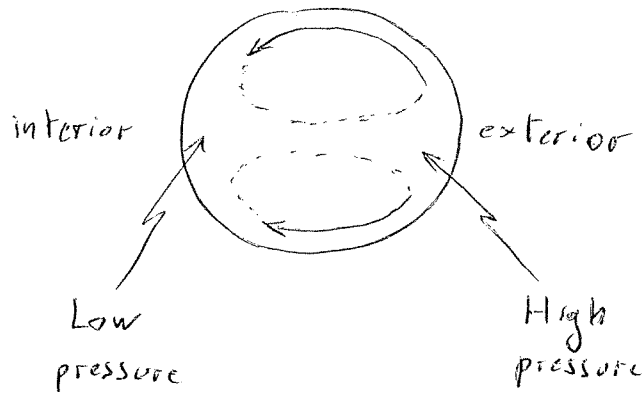
Both for  $z > 0$  and  $z < 0$   $(\vec{\omega} \cdot \vec{\nabla}) u_x$  is strengthening  $\omega_x$ . So in

section B-B' we still have  $\omega_x > 0$  when  $z > 0$   
 $\omega_x < 0$  when  $z < 0$

This explains the recirculations observed.

[25%]

(ii) In the bend, centrifugal forces are balanced by pressure



This pressure difference causes a flow near the boundary where centrifugal forces  $\left(\frac{v^2}{R}\right)$  are weak due to the no-slip condition. [25%]

c) Length of recirculation

Vorticity equation in the x-direction

$$\frac{D \omega_x}{Dt} = (\bar{\omega} \cdot \bar{v}) \omega_x + \nu \nabla^2 \omega_x$$

Vortex stretching is very weak after the bend. Vorticity in the x direction diffuses downstream until it is homogeneous  $\omega_x = 0$

$$\underbrace{U_x \frac{\partial \omega_x}{\partial x}}_{\text{convection downstream}} \approx \underbrace{\nu \nabla^2 \omega_x}_{\text{diffusion within 2 cross-section}}$$

$$\bar{v} \frac{\omega_x}{L} \approx \nu \frac{\omega_x}{(D/2)^2}$$

Diffusion has killed the recirculation when [20%]

$$L \sim \frac{D^2 \bar{v}}{4 \nu} \quad \text{about } 2m \quad \frac{L}{D} \sim \frac{Re}{4}$$

## Question 2

and divergence free

a)  $(-\alpha x, 0, \alpha z)$  is curl free  $\checkmark$  hence  $\vec{v}$  must carry the vorticity and be also divergence free. [20%]

b) Assuming the same  $e^{i(k_0 x + k_0 y)}$  form

$$v_x = B e^{i(k_0 x + k_0 y)}$$

$$v_y = C e^{i(k_0 x + k_0 y)}$$

$$\vec{v} \cdot \vec{v} = (ik_0 B + ik_0 C) e^{i(k_0 x + k_0 y)}$$

$$\hookrightarrow C = -B$$

2D flow, hence  $\omega_z$  only  $\omega_z = \frac{\partial v_y}{\partial x} - \frac{\partial v_x}{\partial y}$

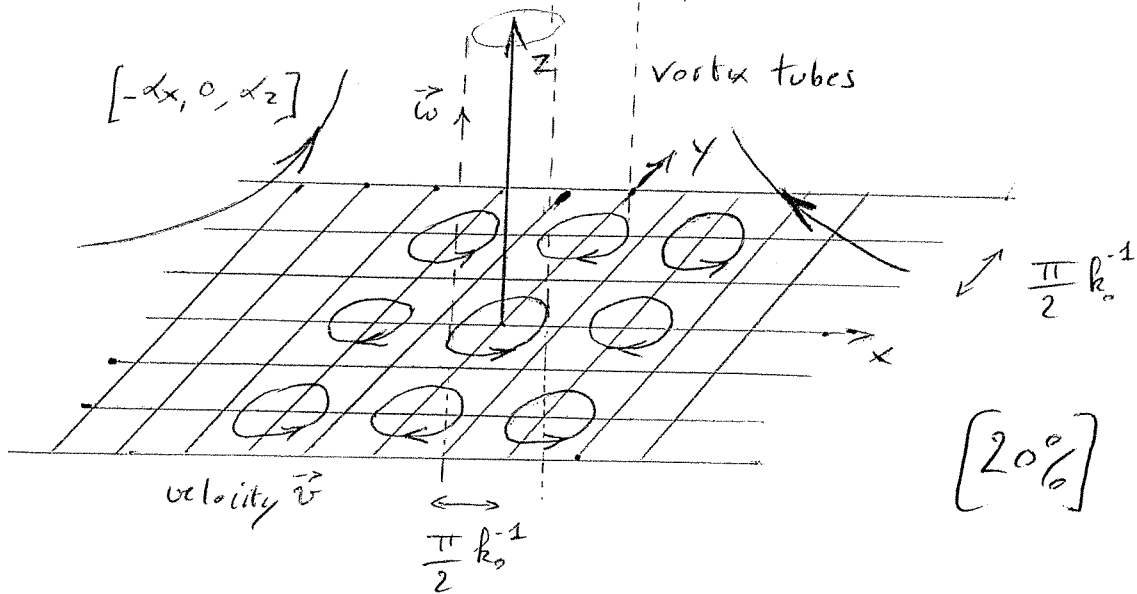
$$\omega_z = (ik_0 C - ik_0 B) e^{i(k_0 x + k_0 y)}$$

$$= -2ik_0 B e^{i(k_0 x + k_0 y + \frac{\pi}{2})}$$

Hence  $v_x = -\frac{A_0}{2k_0} \operatorname{Re} \left[ e^{i(k_0 x + k_0 y + \frac{\pi}{2})} \right]$ ,  $v_y = \frac{A_0}{2k_0} \operatorname{Re} \left[ e^{i(k_0 x + k_0 y + \frac{\pi}{2})} \right]$

c)  $\operatorname{Re} \left[ e^{i(k_0 x + k_0 y)} \right] = \cos(k_0 x + k_0 y)$

[30%]



d)  $\frac{\partial \omega_z}{\partial t} + \vec{v} \cdot \nabla \omega_z - \alpha x \frac{\partial \omega_z}{\partial x} = \omega_z \frac{\partial v_z}{\partial z}$

$\downarrow$   
 $k_x v_x + k_y v_y = 0$  (continuity)

$$\hookrightarrow \frac{dA}{dt} + A i \frac{dk_x}{dt} x = \alpha A + \alpha x i k_x A$$

$$\rightarrow \frac{dA}{dt} = \alpha A \quad \text{and} \quad \frac{dk_x}{dt} = k_x \alpha$$

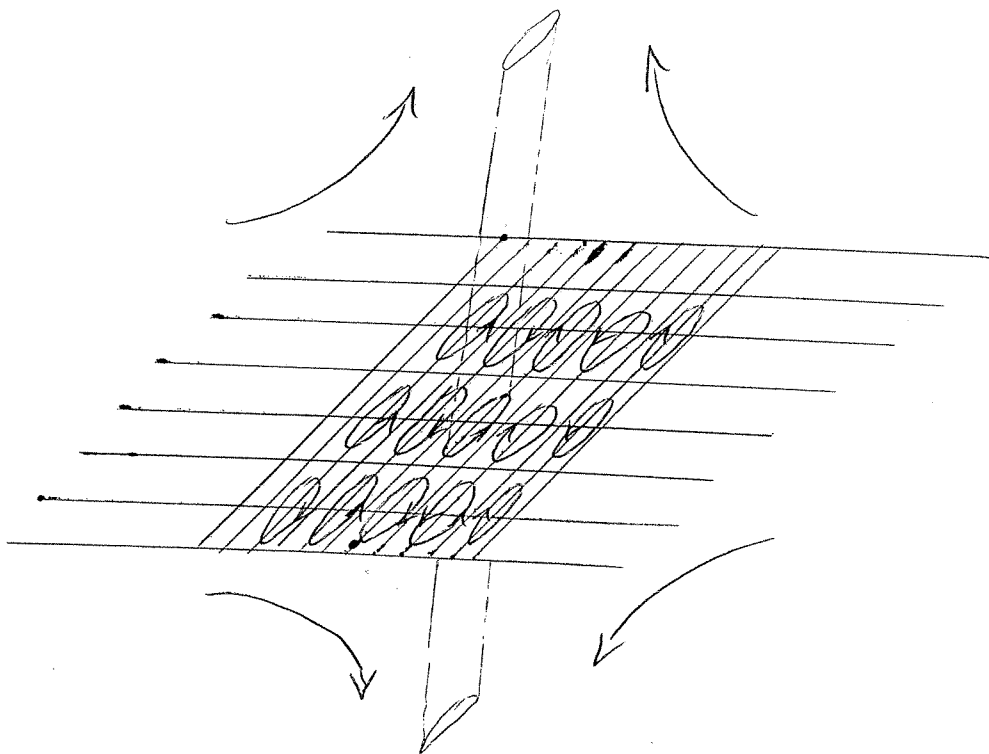
Hence  $A = A_0 e^{\alpha t}$

$$k_x = k_0 e^{\alpha t}$$

Physical interpretation: the straining part of the flow  $[-\alpha x, 0, \alpha z]$  squeezes the vortex tubes towards the plane  $x=0$  and stretches these tubes along the  $z$ -direction:

hence  $A$  grows (exponentially)

and the width of the vortex tubes in the  $x$ -direction,  $\propto k_x^{-1}$  decreases (exponentially)



[30°]

### Question 3

a) No shear stress on  $\mathcal{C}$ , hence conservation of the total angular momentum, hence  $\int_S \omega dS$  is conserved [20%]

b)  $\int_S \vec{U} \cdot [\text{Navier-Stokes}] dS$

$$= \int_S \frac{\partial}{\partial t} \left( \frac{U^2}{2} \right) + \underbrace{\vec{U} \cdot \vec{\nabla}}_0 \frac{U^2}{2} + \underbrace{\vec{U} \cdot \vec{\nabla}}_0 p - \nu \vec{U} \cdot \vec{\nabla}^2 \vec{U} dS = 0$$

$$\rightarrow \frac{\partial}{\partial t} \int_S \frac{U^2}{2} dS = -\nu \int_S \vec{\nabla} \vec{U} : \vec{\nabla} \vec{U} dS \Rightarrow \int_S \omega^2 dS \quad [20\%]$$

c)  $\int_S \omega$  [vorticity equation]  $dS$

$$= \int_S \frac{\partial}{\partial t} \left( \frac{\omega^2}{2} \right) + \underbrace{(\vec{U} \cdot \vec{\nabla})}_{\downarrow 0} \frac{\omega^2}{2} - \nu \omega \nabla^2 \omega dS = 0$$

$$D = \int_S \omega^2 dS, \text{ hence } \frac{\partial D}{\partial t} = -2\nu \int_S \vec{\nabla} \omega \cdot \vec{\nabla} \omega dS$$

From lectures  $\vec{\nabla} \omega$  is subjected to "stretching"  $\rightarrow \vec{\nabla} \omega$  can grow on a fast turnover timescale [20%]

d) Solid rotation:  $U_\theta = \Omega r \rightarrow \omega = 2\Omega$

$$\left. \begin{aligned} \Phi_0 &= \int_S \omega dS = \int_0^R \omega 2\pi r dr = 2\pi R^2 \Omega \\ E_0 &= \int_S \frac{U_\theta^2}{2} dS = \int_0^R \frac{U_\theta^2}{2} 2\pi r dr = \frac{\pi}{4} R^4 \Omega^2 \end{aligned} \right\} \Phi_0^2 = 16\pi E_0$$

Uniform vorticity distribution corresponds to the minimum  $\int_S \omega^2 dS$  [20%]

e) When  $\Phi_0^2 > 2\pi E_0$ , we want less velocity with the same amount of vorticity  $\rightarrow$  concentrate vorticity on the boundary



When  $\Phi_0^2 < 2\pi E_0$ , vorticity is concentrated at the center



[20%]

### Question 4

a)  $k = \frac{U^2}{z}$  . In the log region  $U \sim U_*$   
hence  $k \sim U_*^2$   $k = \alpha U_*^2$  [20%]

b)  $\frac{U}{U_*} = A \ln y + B$   
 $= A \ln \left( \frac{y U_*}{\nu} \right) + B$

At the first grid point  $\frac{U}{\sqrt{k/2}} = A \ln \left( \frac{y_0 \sqrt{k/2}}{\nu} \right) + B$  [20%]

c) Energy dissipation  $\varepsilon = \frac{U^3}{l}$

In the log-law region,  $l \sim y$ , hence the energy dissipation is  $\varepsilon = \frac{U_*^3}{y_0}$ . [30%]

or Energy production: Reynolds stress  $\times$  mean velocity deformation

$$U_*^2 \frac{\partial U}{\partial y} \approx U_*^2 \frac{U_*}{y_0}$$

$$\varepsilon = A k^2 \frac{\sqrt{k/2}}{y_0}$$

d)  $\frac{\partial U}{\partial y}(y_0) = A \frac{U_*}{y_0} = A \frac{\sqrt{k/2}}{y_0}$

$$\frac{\partial U}{\partial y}(y_0) = A \frac{\sqrt{k/2}}{y_0}$$

[30%]

3 boundary conditions for the 3 unknowns  $U, \varepsilon, k$