

4C7 RANDOM AND NON-LINEAR VIBRATIONS - CRIBS 2003

a) $z(t) = x^2(t)$

$R_{zz}(\tau) = E[z(t)z(t+\tau)] = E[x(t)x(t)x(t+\tau)x(t+\tau)]$

Use formula (i) from question to give:-

$R_{zz}(\tau) = E[x(t)x(t+\tau)]E[x(t)x(t+\tau)] + E[x(t)x(t+\tau)]E[x(t)x(t+\tau)] + E[x(t)x(t)]E[x(t+\tau)]$

$R_{zz}(\tau) = 2R_{xx}^2(\tau) + R_{xx}^2(0)$ [20%]

b) $S_{zz}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} R_{zz}(\tau) e^{-i\omega\tau} d\tau$

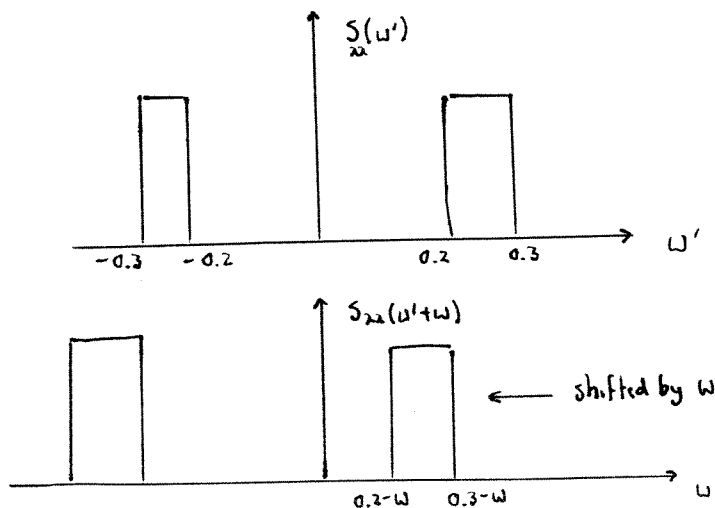
$= \frac{1}{2\pi} \int_{-\infty}^{\infty} 2R_{xx}^2(\tau) e^{-i\omega\tau} d\tau + \frac{1}{2\pi} \int_{-\infty}^{\infty} R_{xx}^2(0) e^{-i\omega\tau} d\tau$

$\frac{1}{\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} S_{xx}(\omega') S_{xx}(\omega'') e^{-i\omega'\tau} e^{i\omega''\tau} e^{-i\omega\tau} d\omega' d\omega'' d\tau$
because $R_{xx}(\tau) = \int_{-\infty}^{\infty} S_{xx}(\omega) e^{i\omega\tau} d\omega$

$= 2 \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} S_{xx}(\omega') S_{xx}(\omega'') \delta(\omega'' - \omega' - \omega) d\omega' d\omega'' + R_{xx}^2(0) \delta(\omega)$

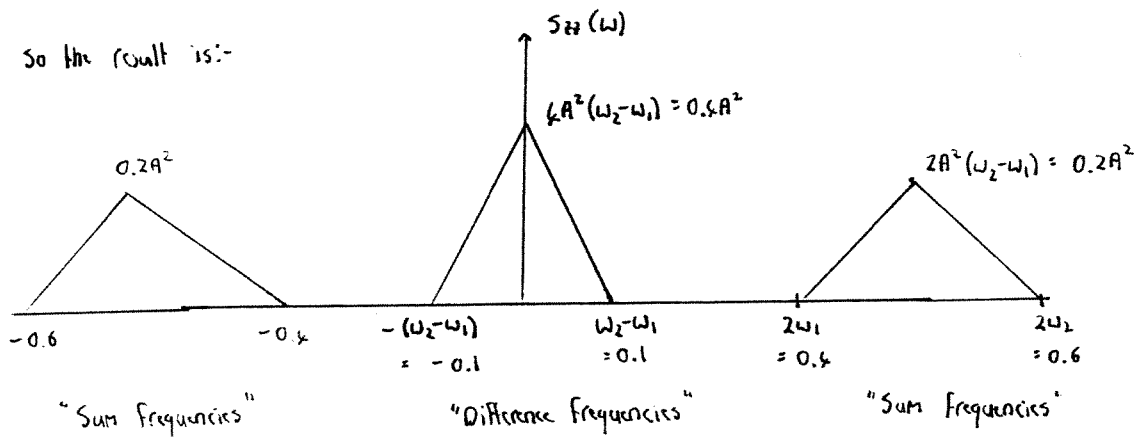
$S_{zz}(\omega) = 2 \int_{-\infty}^{\infty} S_{xx}(\omega') S_{xx}(\omega' + \omega) d\omega' + R_{xx}^2(0) \delta(\omega)$ [35%]

c)



Product of functions
= A^2 where blocks
overlap.
Overlap region = $\omega_2 - \omega_1 - \omega$
etc

So the result is:-



[35%]

d) See above diagram - there is excitation at the resonant frequency.

[10%]

$$2 \text{ a) } M\ddot{x}_1 + C\dot{x}_1 + kx_1 = F + \alpha G \quad \text{--- (1)}$$

$$M\ddot{x}_2 + C\dot{x}_2 + kx_2 = F - \alpha G \quad \text{--- (2)}$$

Since F and G are statistically independent, the spectrum of $F + \alpha G$ is $S_{FF}(\omega) + \alpha^2 S_{GG}(\omega) = (1 + \alpha^2) S_0$

The spectrum of $F - \alpha G$ is also $(1 + \alpha^2) S_0 \Rightarrow$ Each oscillator has the same response spectrum

Using standard result for the response of a linear system to white noise:

$$\sigma_{x_1}^2 = \frac{\pi S_0 (1 + \alpha^2)}{Ck} = \sigma_{x_2}^2 \quad [25\%]$$

b) Subtract equations (1) and (2) to give:-

$$M(\ddot{x}_1 - \ddot{x}_2) + C(\dot{x}_1 - \dot{x}_2) + k(x_1 - x_2) = 2\alpha G$$

Putting $z = x_1 - x_2$ and using standard response to White Noise results:-

$$\sigma_z^2 = \frac{k\alpha^2 \pi S_0}{Ck} \quad \sigma_z^2 = \frac{k\alpha^2 \pi S_0}{CM} \quad [20\%]$$

$$c) \quad \sigma_z^2 = \frac{k(0.5)^2 \pi \times 3 \times 10^{-5}}{0.04 \times k} \Rightarrow \sigma_z^2 = 0.024$$

$$\sigma_{\dot{z}}^2 = \sigma_z^2 \left(\frac{k}{M}\right) \Rightarrow \sigma_{\dot{z}}^2 = 0.048$$

Impact rate = crossing rate ν_b^+ with $b = d = 0.1$

$$\nu_b^+ = \left(\frac{1}{2\pi}\right) \left(\frac{\sigma_{\dot{z}}}{\sigma_z}\right) e^{-\frac{1}{2} \left(\frac{0.1}{0.024}\right)^2} = 5.4 \times 10^{-5}$$

$$\text{Prob of impact} = 1 - e^{-\nu_b^+ T} ; T = 1 \times 60 \times 60 \Rightarrow \text{Prob} = 0.176 \quad [30\%]$$

d) In this case subtracting equations (1) and (2) gives:-

$$M(\bar{x}_1 - \bar{x}_2) + C(\bar{x}_1 - \bar{x}_2) + k(x_1 - x_2) = \underbrace{(\gamma - 1)F + (\gamma + 1)\alpha r}_1$$

$$\text{Input spectrum} = [(\gamma - 1)^2 + (\gamma + 1)^2 \alpha^2] S_0$$

$$\text{Put } \frac{d}{d\gamma} = 0 \Rightarrow 2(\gamma - 1) + 2(\gamma + 1)\alpha^2 = 0$$

$$\Rightarrow \underline{\gamma = \frac{1 - \alpha^2}{1 + \alpha^2}}$$

[25%]

3 (a) $U = x^4 - 2x^2 + kx$, so governing equation is $\ddot{x} + 4x^3 - 4x + k = 0$

If $k=0$, write in standard first-order form as

$$\begin{cases} \dot{x} = y \\ \dot{y} = -4x^3 + 4x \end{cases}$$

Singular points where $y=0$ and $4x^3 = 4x \rightarrow x = 0, \pm 1$.

(i) Near $x=0$ $\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} \approx \begin{bmatrix} 0 & 1 \\ 4 & 0 \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$

Eigenvalues λ satisfy $\begin{vmatrix} -\lambda & 1 \\ 4 & -\lambda \end{vmatrix} = 0$

$\therefore \lambda^2 = 4, \therefore \lambda = \pm 2 = \text{saddle point}$

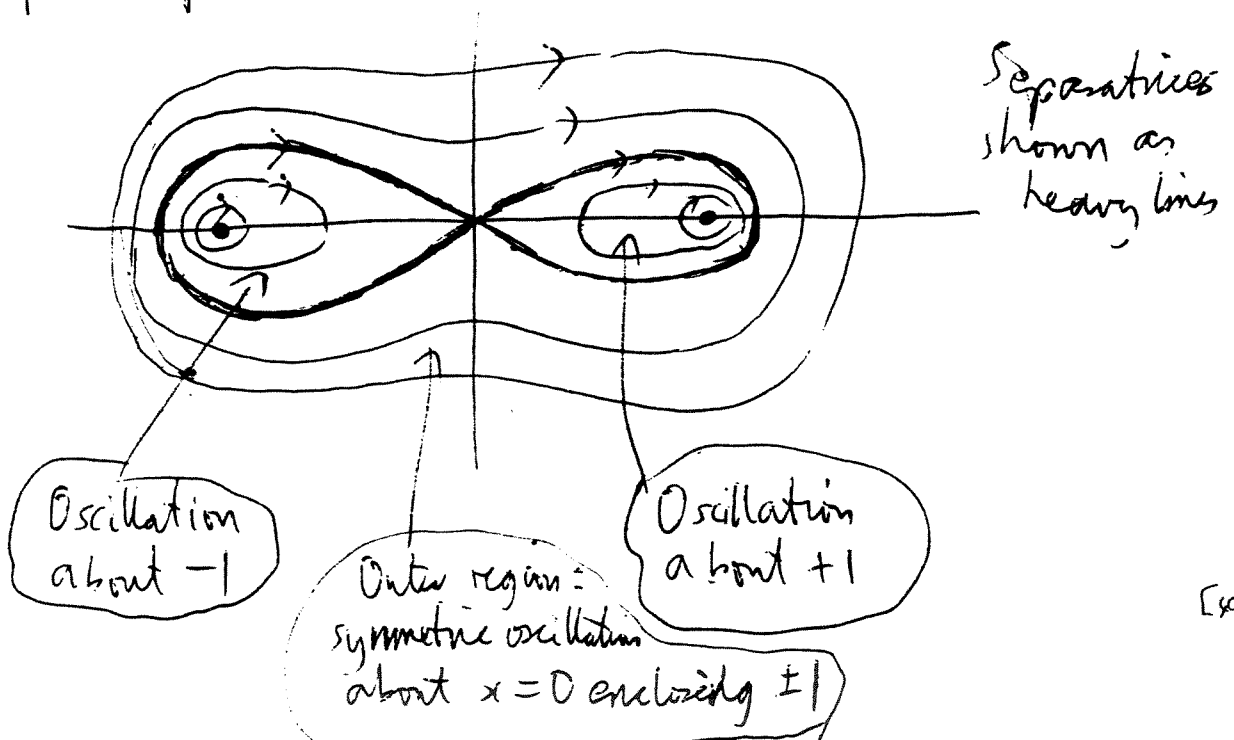
(ii) Near $x=1$: let $x = 1 + \epsilon$

Then $\begin{cases} \dot{\epsilon} = y \\ \dot{y} = -4(1+\epsilon)^3 + 4(1+\epsilon) \approx -8\epsilon \end{cases}$

So eigenvalues λ satisfy $\begin{vmatrix} -\lambda & 1 \\ -8 & -\lambda \end{vmatrix} = 0$

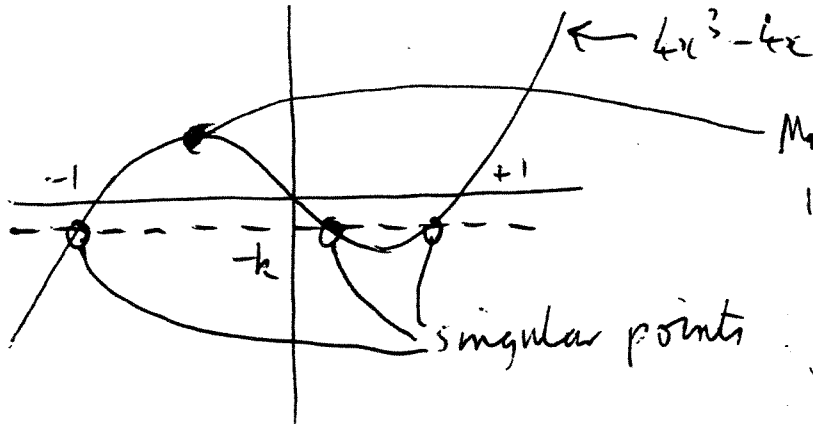
$\therefore \lambda^2 = -8, \text{ so } \lambda = \pm i\sqrt{8} = \text{centre}$

By symmetry, $x = -1$ is an identical centre.
So phase portrait is:



(b) With $k \neq 0$, singular points satisfy $\begin{cases} y = 0 \\ -4x^3 + 4x - k = 0 \end{cases}$

ie where $4x^3 - 4x = -k$



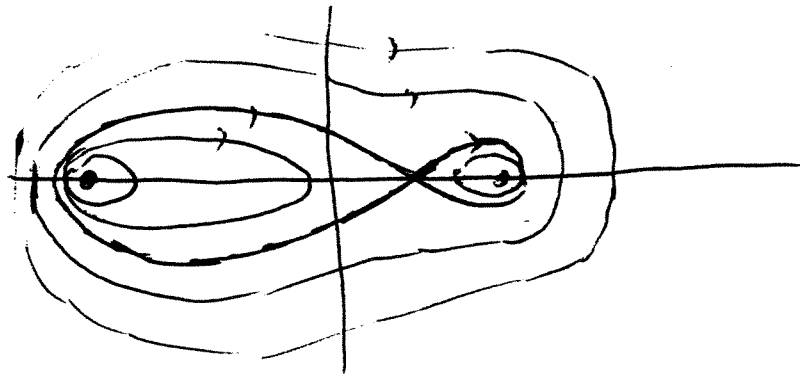
Max value where
 $12x^2 - 4 = 0$
 ie $x = \pm \frac{1}{\sqrt{3}}$

Value then is
 $\frac{+4}{\sqrt{3}} \left(\frac{1}{3} - 1 \right) = \frac{+8}{3\sqrt{3}} = \pm k_{crit}$

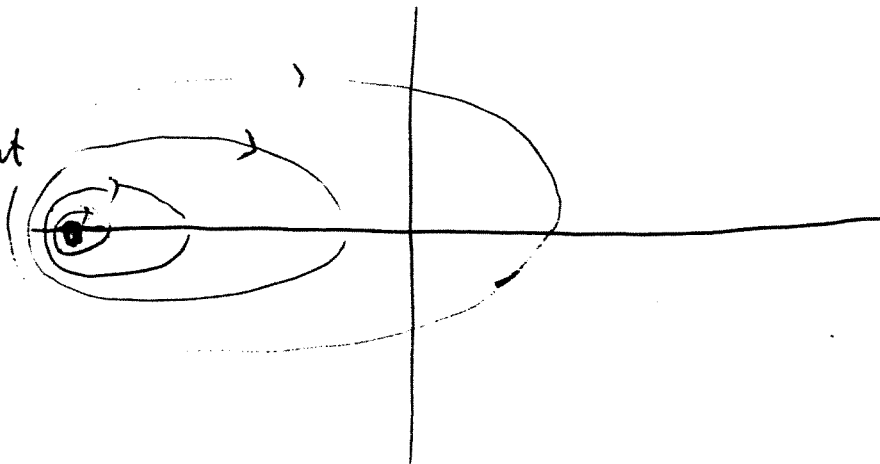
For $|k| < k_{crit}$, get 3 singular points with same character as part (a)

For $|k| > k_{crit}$, only one singular point which is a centre [20%]

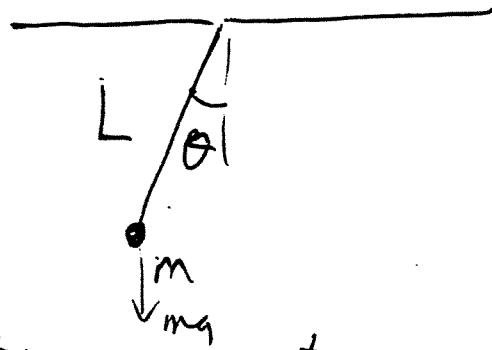
(c)
 $|k| < k_{crit}$
 $k > 0$



$k > k_{crit}$



4 (a)



$$L = L_0 + a \sin \omega t$$

Acceleration components are $L\ddot{\theta} + 2\dot{L}\dot{\theta}$



To avoid needing to know tension in string, resolve perpendicular to string:

$$m(L\ddot{\theta} + 2\dot{L}\dot{\theta}) = -mg \sin \theta$$

$$\therefore L\ddot{\theta} + 2\dot{L}\dot{\theta} + g \sin \theta = 0$$

$$\therefore (L_0 + a \sin \omega t) \ddot{\theta} + 2a\omega \cos \omega t \dot{\theta} + g \sin \theta = 0 \quad [25\%]$$

(b) For $|\theta| \ll 1$, replace $\sin \theta \approx \theta$

Then equation can be written

$$L_0 \ddot{\theta} + g \theta \approx -a \sin \omega t \ddot{\theta} - 2a\omega \cos \omega t \dot{\theta}$$

$$\therefore \ddot{\theta} + \Omega^2 \theta \approx -\frac{a}{L} [\ddot{\theta} \sin \omega t + 2\omega \dot{\theta} \cos \omega t]$$

$$(\Omega^2 = g/L_0)$$

With $\frac{a}{L} \ll 1$, solve by iteration.

$$\text{Zero order: } \ddot{\theta}_0 + \Omega^2 \theta_0 = 0$$

\rightarrow general solution $\theta_0 = A \cos(\Omega t + \phi)$, A, ϕ constants

$$\text{First order: solve } \ddot{\theta}_1 + \Omega^2 \theta_1 \approx -\frac{a}{L} [\ddot{\theta}_0 \sin \omega t + 2\omega \dot{\theta}_0 \cos \omega t]$$

$$\begin{aligned} \therefore \ddot{\theta}_1 + \Omega^2 \theta_1 &= \frac{Aa}{L} \left[\Omega^2 \cos(\Omega t + \phi) \sin \omega t + 2\omega\Omega \sin(\Omega t + \phi) \cos \omega t \right] \\ &= \frac{Aa}{L} \left\{ \frac{\Omega^2}{2} \left[\sin(\omega t + \Omega t + \phi) + \sin(\omega t - \Omega t - \phi) \right] \right. \\ &\quad \left. + \omega\Omega \left[\sin(\omega t + \Omega t + \phi) - \sin(\omega t - \Omega t - \phi) \right] \right\} \end{aligned}$$

Provided $\omega + \Omega \neq \Omega$ and $\omega - \Omega \neq \Omega$, we can find a P.I. by trying $\theta_1 = \alpha \sin(\omega t + \Omega t + \phi) + \beta \sin(\omega t - \Omega t - \phi)$
 Substitute:

$$\begin{aligned} -\alpha[(\omega + \Omega)^2 - \Omega^2] \sin(\omega t + \Omega t + \phi) - \beta[(\omega - \Omega)^2 - \Omega^2] \sin(\omega t - \Omega t - \phi) \\ = \text{RHS} \\ \therefore \alpha = -\frac{Aa\Omega}{2L} \frac{(\Omega + 2\omega)}{(\omega^2 + 2\omega\Omega)}, \quad \beta = -\frac{Aa\Omega}{2L} \frac{(\Omega - 2\omega)}{(\omega^2 - 2\omega\Omega)} \end{aligned}$$

This solution describes steady, non-growing response.

$\omega + \Omega = \Omega$ is not possible, but $\omega - \Omega = \Omega \Rightarrow \omega = 2\Omega$ is possible.
 Then for a P.I. have to try $\theta_1 = \alpha \sin(3\Omega t + \phi) + \gamma t \cos(\Omega t + \phi)$
 α is the same as before.

For γ term:
$$\begin{cases} \dot{\theta}_1 = -\gamma t \Omega \sin(\Omega t + \phi) + \gamma \dots \cos(\Omega t + \phi) \\ \ddot{\theta}_1 = -\gamma t \Omega^2 \cos(\Omega t + \phi) - 2\gamma \Omega \sin(\Omega t + \phi) \end{cases}$$

\therefore need
$$-2\gamma \Omega \sin(\Omega t + \phi) = \frac{Aa\Omega(\Omega - 2\omega)}{2L} \sin(3\Omega t + \phi)$$

$$\therefore \gamma = \frac{3Aa\Omega}{4L}$$

[50%]

(c) Growing oscillation only in the case $\omega = 2\Omega$.
 Phase ϕ plays no role - the value of γ does not depend on ϕ , rather surprisingly.

[25%]