

Q1 (a) From lecture notes: velocity potential  $\phi$  for the flow satisfies  $\nabla^2 \phi = 0$

Boundary conditions:

$$(1) \text{ Top surface } \frac{\partial^2 \phi}{\partial t^2} = -g \frac{\partial \phi}{\partial z} \text{ on } z=0$$

(2) Bottom: deep water, so  $u = \nabla \phi \rightarrow 0$  as  $z \rightarrow -\infty$

(3) Side walls: no flow through walls so  $\frac{\partial \phi}{\partial y} = 0$  on  $y=0, L$

The model neglects viscosity and surface tension. It also assumes that wave amplitudes are small.

(b) Try separable solution  $\phi = f(z)g(y) e^{i(\omega t - kx)}$

$$\text{Then } \nabla^2 \phi = 0 \text{ gives } -k^2 fg + f''g + g''f = 0$$

$$\therefore \frac{f''}{f} + \frac{g''}{g} - k^2 = 0$$

This is in separated form, so deduce that

$$\frac{g''}{g} = -\nu^2 \text{ say, and } \frac{f''}{f} = k^2 + \nu^2 \text{ (r constant)}$$

So  $g = \begin{cases} \sin \nu y \\ \cos \nu y \end{cases}$ . But  $g'(0) = 0$ , so must be  $\cos$ .

Then  $g'(y) = 0$  at  $y=L$ , so  $\sin \nu L = 0$

$$\therefore \nu = \frac{n\pi}{L}, n=0, 1, 2, 3 \dots$$

$$\text{Then } f'' = (k^2 + \nu^2)f \text{ so } f = e^{\pm \sqrt{k^2 + \nu^2} z}$$

But  $f' \rightarrow 0$  as  $z \rightarrow -\infty$  (deep water) so only the + sign is relevant.

So possible solutions are  $\phi = \phi_0 e^{\sqrt{k^2 + \nu^2} z} \cos \nu y e^{i(\omega t - kx)}$

$$\text{with } \nu = \frac{n\pi}{L}, n=0, 1, 2 \dots$$

Now the top surface condition gives

$$-\omega^2 = -g \sqrt{k^2 + \nu^2}$$

$$\therefore k^2 = \frac{\omega^4}{g^2} - \left(\frac{n\pi}{L}\right)^2, \quad n=0, 1, 2, \dots$$

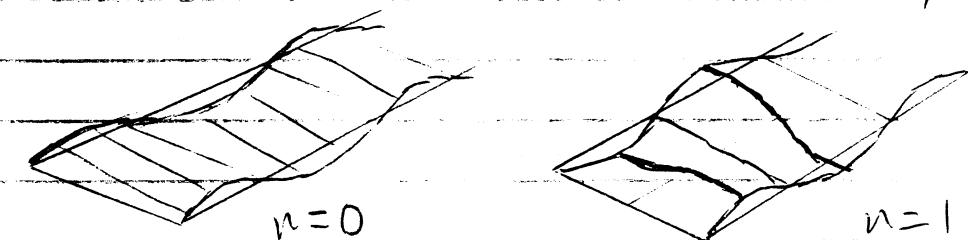
(c) Cut-on frequency where  $k^2$  goes from negative to positive, so when  $\frac{\omega^2}{g} = \frac{n\pi}{L}$ .

Please velocity  $c_p = \frac{\omega}{k} = \sqrt{\frac{\omega^4}{g^2} - \left(\frac{n\pi}{L}\right)^2}$

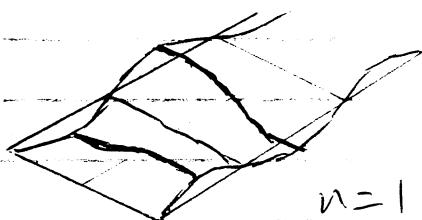
Group velocity  $c_g = \frac{d\omega}{dk} = \frac{d}{dk} \left[ g^2(k^2 + \nu^2) \right]^{1/4}$   
 $= \frac{1}{4} \left[ g^2(k^2 + \nu^2) \right]^{-3/4} \cdot 2g^2 k$   
 $= \frac{g^2}{2} \frac{k}{\omega^3} = \frac{g^2}{2\omega^3} \sqrt{\frac{\omega^4}{g^2} - \left(\frac{n\pi}{L}\right)^2}$

Note that at a cut-on frequency,  $c_p \rightarrow \infty$ ,  $c_g \rightarrow 0$ .

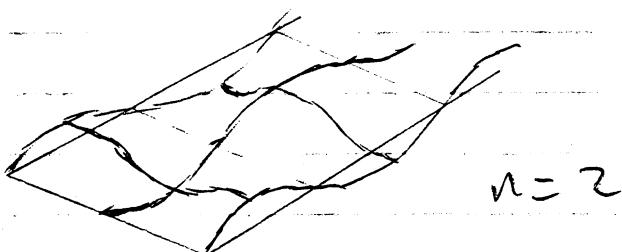
(d)  $L = 5\text{m}$ , so cut-on frequencies are  $n=0, 0\text{Hz}$ ,  
 $n=1, 0.395\text{Hz}$ ,  
 $n=2, 0.559\text{Hz}$



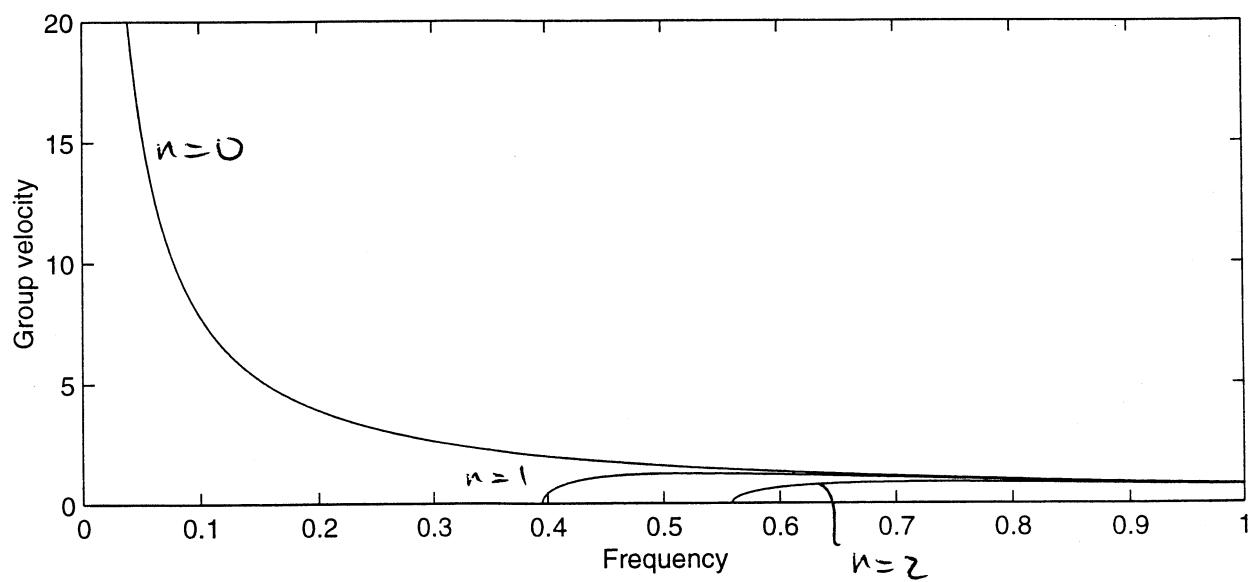
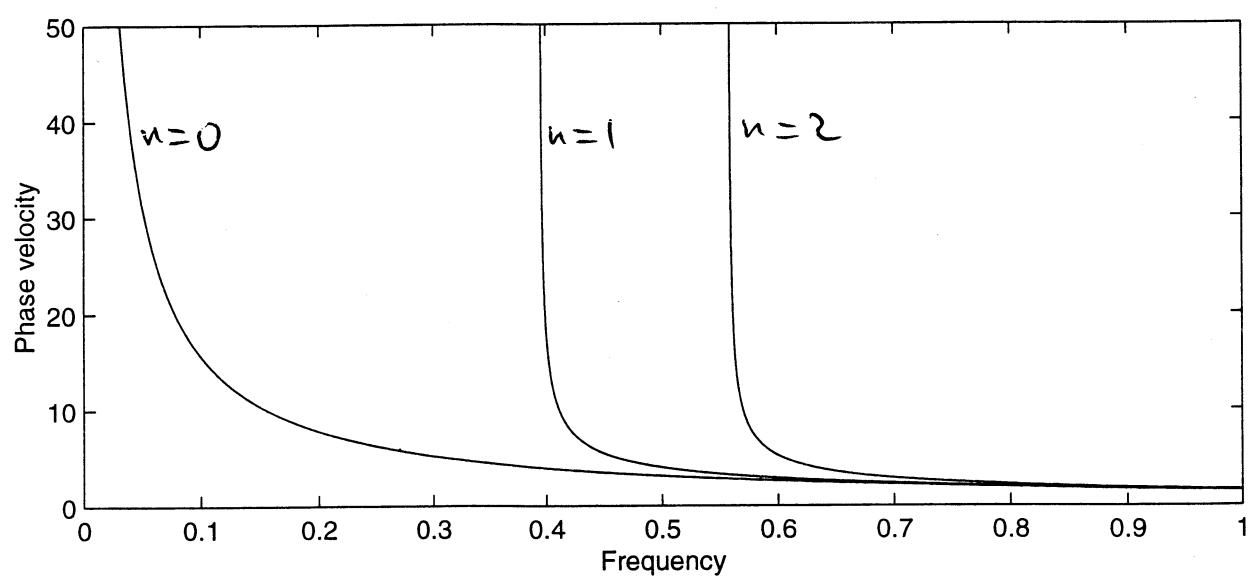
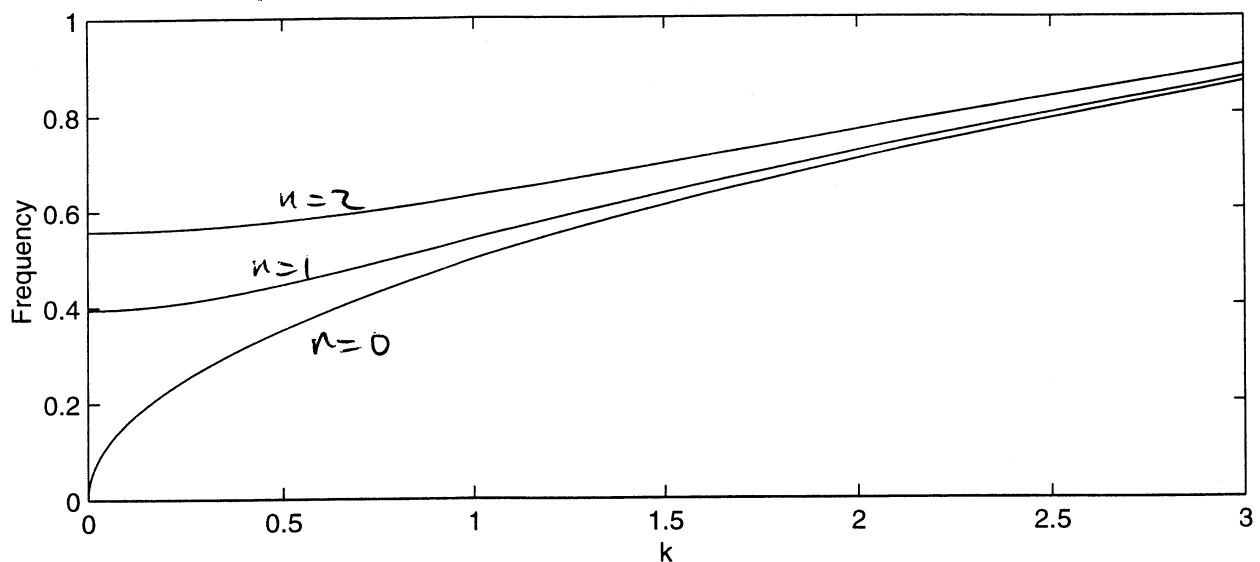
$n=0$



$n=1$



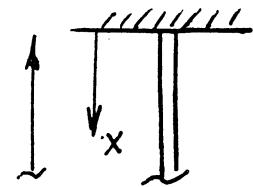
$n=2$



Q2 (a) conservation of translational momentum

$$M \tilde{V}_0 = 2M V_0 \rightarrow V_0 = \tilde{V}_0/2$$

$$\sigma_0 = \frac{Mg}{A}$$



(b) eq of motion for mass,  $t > 0$

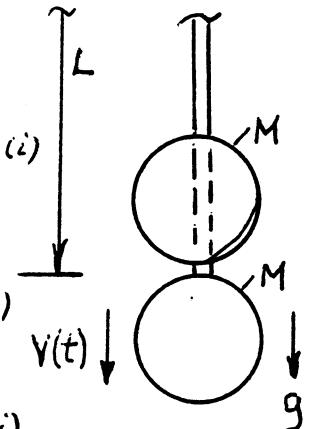
$$2M \frac{dV}{dt} = -A [\sigma_i(L,t) + \sigma_r(L,t)] - A\sigma_0 + 2Mg \quad (i)$$

continuity of velocity

$$V(t) = \dot{u}_i(L,t) + \dot{u}_r(L,t) \quad (ii)$$

cons. momentum across wavefront,  $t > 0$

$$\sigma_i = -\rho c_0 \dot{u}_i, \quad \sigma_r = \rho c_0 \dot{u}_r \quad (iii)$$



$$\therefore \frac{dV}{dt} = \frac{\rho c_0 A}{2M} \left\{ \dot{u}_i(L,t) - [V(t) - \dot{u}_i(L,t)] \right\} + \frac{g}{2}$$

$$\text{or } \frac{dV}{dt} + \frac{c_0}{mL} V = 2\dot{u}_i(L,t) + \frac{g}{2} \quad \text{where } m = \frac{2M}{\rho AL}$$

Sol'n

$$V = \frac{g mL}{2c_0} + B e^{-\frac{c_0 t}{mL}}$$

$$\text{I.C. } V_0 = V(0) = \frac{g mL}{2c_0} + B \Rightarrow B = V_0 - \frac{g mL}{2c_0}$$

$$\therefore V(t) = \alpha + (V_0 - \alpha) e^{-\beta t},$$

$$\alpha = \frac{g mL}{2c_0}, \quad \beta = \frac{c_0}{mL}$$

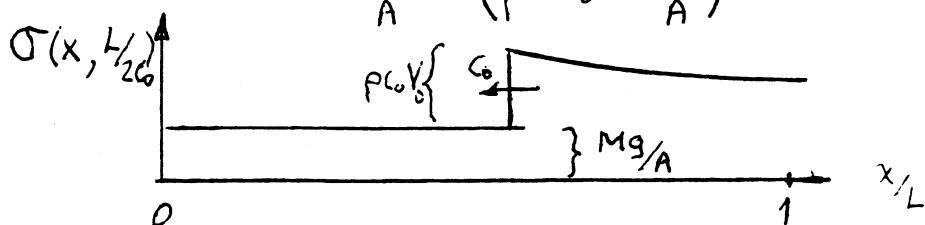
$$t < 2L/c_0$$

(c)

$$\sigma_r(L,t) = \rho c_0 V(t)$$

$$\sigma_r(x,t) = \rho c_0 \left\{ \alpha + (V_0 - \alpha) e^{-\frac{(c_0 t + x - L)}{mL}} \right\} \quad x > L - c_0 t$$

$$= \frac{Mg}{A} + \left( \rho c_0 V_0 - \frac{Mg}{A} \right) e^{-\frac{(c_0 t + x - L)}{mL}}$$



(d) at clamped end when wavefront arrives,  $t = L/c_0$

$$\sigma(0, \frac{L}{c_0}^+) = \sigma_F = \frac{Mg}{A} + 2\rho c_0 V_0 = \frac{Mg}{A} + 2\rho c_0 \frac{\tilde{V}_0 F}{2}$$

$$\therefore \tilde{V}_0 F = \frac{\sigma_F}{\rho c_0} - \frac{Mg}{\rho c_0 A} = \frac{\sigma_F}{\rho c_0} - \frac{mgh}{2c_0}$$

break located at  $x = 0$  (the clamp) at  $t = L/c_0^+$

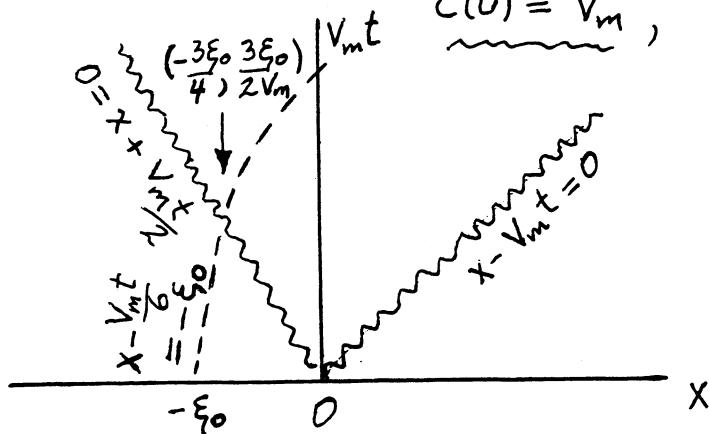
Q3 (a) flux:  $q = \rho V = \rho_m V_m (y - y^{3/2})$ ,  $y = \rho/\rho_m$

$$0 = \frac{d(q)}{dy} = \rho_m V_m \left(1 - \frac{3}{2} y^{1/2}\right)$$

$$\therefore \rho = \frac{4}{9} \rho_m \rightarrow (\text{flux})_{\max}$$

(b) characteristic speeds  $c(\rho) = dq/d\rho = V_m \left(1 - \frac{3}{2} \sqrt{\rho/\rho_m}\right)$

$$c(0) = V_m, c(\rho_m) = -V_m/2$$



constrained speed

$$V = V_m \left(1 - \sqrt{\rho/\rho_m}\right) = V_m \left(1 - \frac{5}{6}\right) = V_m/6$$

(c) on characteristic lines in fan emanating from (0,0)

$$x = \frac{dq}{d\rho} t = V_m t \left(1 - \frac{3}{2} \sqrt{\rho/\rho_m}\right) \rightarrow \sqrt{\rho/\rho_m} = \frac{2}{3} \left(1 - \frac{x}{V_m t}\right)$$

speed of beam located at  $x = \xi(t)$

$$\frac{d\xi}{dt} = V_m \left[1 - \frac{2}{3} \left(1 - \frac{\xi}{V_m t}\right)\right] = \frac{V_m}{3} \left[1 + \frac{2\xi}{V_m t}\right]$$

$$\therefore 3t \frac{d\xi}{dt} - 2\xi = V_m t$$

sol'n:  $\xi = B_1 V_m t + B_2 t^{\eta}$  }  $3B_1 V_m t + 2B_2 \eta t^{\eta-1} - 2B_2 = V_m t$   
 $\frac{d\xi}{dt} = B_1 V_m + B_2 \eta t^{\eta-1}$  }  $-2B_2 \eta = V_m t$

$$\therefore B_1 = 1, \eta = 2/3$$

$$\xi(t) = V_m t + B_2 t^{2/3}$$

$$(d) \quad \text{in fan:} \quad \xi(t) = V_m t + B_2 t^{2/3}$$

initial cond:

$$\text{bean enters fan at } \left( -\frac{3\xi_0}{4}, \frac{3\xi_0}{2V_m} \right)$$

$$\therefore -\frac{3\xi_0}{4} = \frac{3\xi_0}{2V_m} + B_2 \left( \frac{3\xi_0}{2V_m} \right)^{2/3}$$

$$B_2 = -\frac{9\xi_0}{4} \left( \frac{2V_m}{3\xi_0} \right)^{2/3}$$

$$\xi(t) = V_m t - \underbrace{\frac{9\xi_0}{4} \left( \frac{2V_m t}{3\xi_0} \right)^{2/3}}_{t > \frac{3\xi_0}{2V_m}},$$

