

ENGINEERING TRIPOS ENGINEERING TRIPOS

PART IIB PART IIA

Friday 30 April 2004

2.30 to 4

Module 4C14

NATURAL AND MICRO-ARCHITECTURED MATERIALS

Answer not more than three questions.

All questions carry the same number of marks.

The approximate percentage of marks allocated to each part of a question is indicated in the right margin.

Attachments:

Special datasheet(s) (2 pages).

You may not start to read the questions printed on the subsequent pages of this question paper until instructed that you may do so by the Invigilator



#### 1 Explain the following

- The cell membrane of red blood cells has a very low elastic modulus, but a large lock-up strain and a high ultimate strength. [25%]
- Wood is strongly anisotropic, with a compressive strength along the grain an order of magnitude greater than that across the grain.

[25%]

Proteins are transported within cells at much faster rates than diffusion can provide.

[25%]

Plants and animal cells have different strategies for harvesting energy. (d)

[25%]

- Describe the physical basis for the Young's modulus of biological tissues by 2 explaining how the following concepts dictate the modulus:
  - (i) persistence length

[25%]

nodal connectivity (ii)

[25%]

Qualitatively describe the myosin crossbridge cycle with reference to [25%] conversions of ATP to ADP.

Describe the tension versus length curve of a single muscle fibre. Suppose that the tension decreased nonlinearly with increasing length for striation spacings greater than 2.5 µm. Would this invalidate the theory that the crossbridges working independently generate the tension?

[25%]

- In the Huxley crossbridge model for a muscle, n(x) is the fraction of attached crossbridges, where x is the position of an actin binding site from the equilibrium position of a myosin head. Assume that the attachment and detachment of the crossbridges is governed by a first order kinetic scheme with attachment and detachment constants f(x) and g(x), respectively.
- (a) Determine the steady-state n(x) in terms of f(x) and g(x) for a muscle in isometric tension. [10%]

(b) Given that:

$$f(x) = 0$$
;  $g(x) = g_1$ ;  $x < 0$ ;  
 $f(x) = f_0$ ;  $g(x) = g_0$ ;  $0 \le x \le h$ ;  
 $f(x) = 0$ ;  $g(x) = g_0$   $x > h$ ;

determine n(x) for shortening at a constant velocity V = -dx/dt. Here  $g_0$ ,  $f_0$  and h are constants. [50%]

(c) Writing any appropriate equations, explain how one might use the Huxley crossbridge dynamics model to calculate the response of a muscle in Hill's quick-release experiments (step change in tension). [40%]

4 (a) Describe the functions of arteries and veins in blood flow. Comment on why there is no reverse flow in the veins.

[15%]

(b) Describe the role of capillaries in blood flow.

[15%]

(c) Consider a thin-walled cylindrical tube of length L, radius r and wall thickness t subjected to an internal pressure P. The radius of the tube is  $r_0$  at zero pressure, and the tube is made of an elastic material having Young's modulus E. If the cross-sectional area of the tube varies linearly with pressure, show that the compliance per unit length of the tube can be approximated as  $c = 2\pi r^3/(Et)$ .

[25%]

(d) Comment on why pulmonary and systemic capillaries can be modelled as resistance vessels whereas large arteries cannot, and explain why veins are more compliant than arteries (for a fixed vessel radius).

[15%]

(e) There is blood flow across the cylindrical vessel described in (c), with inlet pressure  $P_o$  and outlet pressure.  $P_1$ . It can be shown that the flux of blood is given by:

$$Q = \frac{\pi r^4}{24\mu L} \frac{(1 + \gamma P_0)^3 - (1 + \gamma P_1)^3}{\gamma}$$

where  $\gamma = c/(\pi r^2)$  and  $\mu$  is the viscosity of blood.

- (i) Explain why the pressure drop in the veins can be much less than in the arteries.
- (ii) Show that, in the limit  $c \to 0$ , the blood vessel reduces to a resistance vessel, with its resistance per unit length given by  $\rho = 8\mu/(\pi r^4)$ .

[Hint: 
$$x^3 - y^3 = (x - y)(x^2 + xy + y^2)$$
.] [30%]

### **END OF PAPER**

Engineering Tripos Part IIA

## Paper G4: Mechanics of Solids

# ELASTICITY and PLASTICITY FORMULAE

# 1. Axi-symmetric deformation: discs, tubes and spheres

Equilibrium  $\sigma_{\theta\theta} = \frac{\mathrm{d}(r\sigma_{\Pi})}{\mathrm{d}r} + \rho\omega^2 r^2 \qquad \qquad \sigma_{\theta\theta} = \frac{1}{2r} \frac{\mathrm{d}(r^2\sigma_{\Pi})}{\mathrm{d}r}$   $\mathrm{Lam's\ equations\ (in\ elasticity)} \qquad \sigma_{\Pi} = \mathrm{A} - \frac{\mathrm{B}}{r^2} - \frac{3+\nu}{8} \ \rho\omega^2 r^2 - \frac{E\alpha}{r^2} \int\limits_{\mathrm{c}}^{\mathrm{r}} r \mathrm{T} \mathrm{d}r \qquad \qquad \sigma_{\Pi} = \mathrm{A} - \frac{\mathrm{B}}{r^3}$   $\sigma_{\theta\theta} = \mathrm{A} + \frac{\mathrm{B}}{r^2} - \frac{1+3\nu}{8} \ \rho\omega^2 r^2 + \frac{E\alpha}{r^2} \int\limits_{\mathrm{c}}^{\mathrm{r}} r \mathrm{T} \mathrm{d}r - E\alpha T \qquad \sigma_{\theta\theta} = \mathrm{A} + \frac{\mathrm{B}}{2r^3}$ 

## 2. Plane stress and plane strain

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	Cartesian coordinates	Polar coordinates
Strains	$\varepsilon_{\rm XX} = \frac{\partial u}{\partial x}$	$ \varepsilon_{\rm rr} = \frac{\partial u}{\partial r} $
	$ \varepsilon_{yy} = \frac{\partial v}{\partial y} $	$\varepsilon_{\theta\theta} = \frac{u}{r} + \frac{1}{r} \frac{\partial v}{\partial \theta}$
	$\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$	$\gamma_{r\theta} = \frac{\partial v}{\partial r} + \frac{1}{r} \frac{\partial u}{\partial \theta} - \frac{v}{r}$
Compatibility	$\frac{\partial^2 \gamma_{xy}}{\partial x \partial y} = \frac{\partial^2 \varepsilon_{xx}}{\partial y^2} + \frac{\partial^2 \varepsilon_{yy}}{\partial x^2}$	$\frac{\partial}{\partial r} \left\{ r \frac{\partial \gamma_{r\theta}}{\partial \theta} \right\} = \frac{\partial}{\partial r} \left\{ r^2 \frac{\partial \varepsilon_{\theta\theta}}{\partial r} \right\} - r \frac{\partial \varepsilon_{\pi}}{\partial r} + \frac{\partial^2 \varepsilon_{\pi}}{\partial \theta^2}$
or (in elasticity)	$\left\{ \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right\} (\sigma_{xx} + \sigma_{yy}) = 0$	$\left\{ \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right\} (\sigma_{\text{tr}} + \sigma_{\theta\theta}) = 0$
Equilibrium	$\frac{\partial \sigma_{xx}}{\partial r} + \frac{\partial \sigma_{xy}}{\partial v} = 0$	$\frac{\partial}{\partial r}(r\sigma_{\text{T}}) + \frac{\partial\sigma_{\text{r}\theta}}{\partial\theta} - \sigma_{\theta\theta} = 0$
<u> 244merum</u>	$\frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{xy}}{\partial x} = 0$	$\frac{\partial \sigma_{\theta\theta}}{\partial \theta} + \frac{\partial}{\partial r}(r\sigma_{r\theta}) + \sigma_{r\theta} = 0$
$ abla^4 \phi = 0$ (in elasticity)	$\left\{ \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right\} \left\{ \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} \right\} = 0$	$\left\{ \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right\}$
	, ,	$\times \left\{ \frac{\partial^2 \phi}{\partial r^2} + \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} \right\} = 0$
Airy Stress Function	$\sigma_{XX} = \frac{\partial^2 \phi}{\partial y^2}$	$\sigma_{\rm ff} = \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2}$
	$\sigma_{yy} = \frac{\partial^2 \phi}{\partial x^2}$	$\sigma_{\Theta\Theta} = \frac{\partial^2 \phi}{\partial r^2}$
	$\sigma_{xy} = -\frac{\partial^2 \phi}{\partial x \partial y}$	$\sigma_{r\theta} = -\frac{\partial}{\partial r} \left\{ \frac{1}{r} \frac{\partial \phi}{\partial \theta} \right\}$



## 3. Torsion of prismatic bars

Prandtl stress function: 
$$\sigma_{zx} \ (= \tau_x) = \frac{dF}{dy}$$
,  $\sigma_{zy} \ (= \tau_y) = -\frac{dF}{dx}$ 

Equilibrium: 
$$T = 2 \int F dA$$

Governing equation for elastic torsion:  $\nabla^2 F = -2G\beta$  where  $\beta$  is the angle of twist per unit length.

# 4. Total potential energy of a body

$$\Pi = U - W$$

where 
$$U = \frac{1}{2} \int_{V} \mathcal{E}^{T}[D] \mathcal{E} dV$$
,  $W = P^{T} \mathcal{U}$  and  $[D]$  is the elastic stiffness matrix.

# 5. Principal stresses and stress invariants

Values of the principal stresses, op, can be obtained from the equation

$$\begin{vmatrix} \sigma_{XX} - \sigma_P & \sigma_{Xy} & \sigma_{Xz} \\ \sigma_{Xy} & \sigma_{yy} - \sigma_P & \sigma_{yz} \\ \sigma_{xz} & \sigma_{yz} & \sigma_{zz} - \sigma_P \end{vmatrix} = 0$$

This is equivalent to a cubic equation whose roots are the values of the 3 principal stresses, i.e. the possible values of op.

Expanding: 
$$\sigma_P^3 - I_1 \sigma_P^2 + I_2 \sigma_P - I_3 = 0$$
 where  $I_1 = \sigma_{xx} + \sigma_{yy} + \sigma_{zz}$ ,

$$I_{2} = \begin{vmatrix} \sigma_{yy} & \sigma_{yz} \\ \sigma_{yz} & \sigma_{zz} \end{vmatrix} + \begin{vmatrix} \sigma_{xx} & \sigma_{xz} \\ \sigma_{xz} & \sigma_{zz} \end{vmatrix} + \begin{vmatrix} \sigma_{xx} & \sigma_{xy} \\ \sigma_{xy} & \sigma_{yy} \end{vmatrix} \quad \text{and} \quad I_{3} = \begin{vmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{xy} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{xz} & \sigma_{yz} & \sigma_{zz} \end{vmatrix}.$$

#### 6. Equivalent stress and strain

Equivalent stress 
$$\bar{\sigma} = \sqrt{\frac{1}{2}} \{ (\sigma_1 - \sigma_2)^2 + (\sigma_2 - \sigma_3)^2 + (\sigma_3 - \sigma_1)^2 \}^{-1/2}$$

Equivalent strain increment 
$$d\bar{\varepsilon} = \sqrt{\frac{2}{3}} \{ d\varepsilon_1^2 + d\varepsilon_2^2 + d\varepsilon_3^2 \}^{-1/2}$$

### 7. Yield criteria and flow rules

Tresca

Material yields when maximum value of  $|\sigma_1 - \sigma_2|$ ,  $|\sigma_2 - \sigma_3|$  or  $|\sigma_3 - \sigma_1| = Y = 2k$ , and then,

if 
$$\sigma_3$$
 is the intermediate stress,  $d\varepsilon_1: d\varepsilon_2: d\varepsilon_3 = \lambda(1:-1:0)$  where  $\lambda \neq 0$ .

von Mises

Material yields when, 
$$(\sigma_1 - \sigma_2)^2 + (\sigma_2 - \sigma_3)^2 + (\sigma_3 - \sigma_1)^2 = 2Y^2 = 6k^2$$
, and then

$$\frac{\mathrm{d}\varepsilon_1}{\sigma_1'} = \frac{\mathrm{d}\varepsilon_2}{\sigma_2'} = \frac{\mathrm{d}\varepsilon_3}{\sigma_3'} = \frac{\mathrm{d}\varepsilon_1 - \mathrm{d}\varepsilon_2}{\sigma_1 - \sigma_2} = \frac{\mathrm{d}\varepsilon_2 - \mathrm{d}\varepsilon_3}{\sigma_2 - \sigma_3} = \frac{\mathrm{d}\varepsilon_3 - \mathrm{d}\varepsilon_1}{\sigma_3 - \sigma_1} = \lambda = \frac{3}{2} \frac{\mathrm{d}\bar{\varepsilon}}{\bar{\sigma}} .$$

October 2000