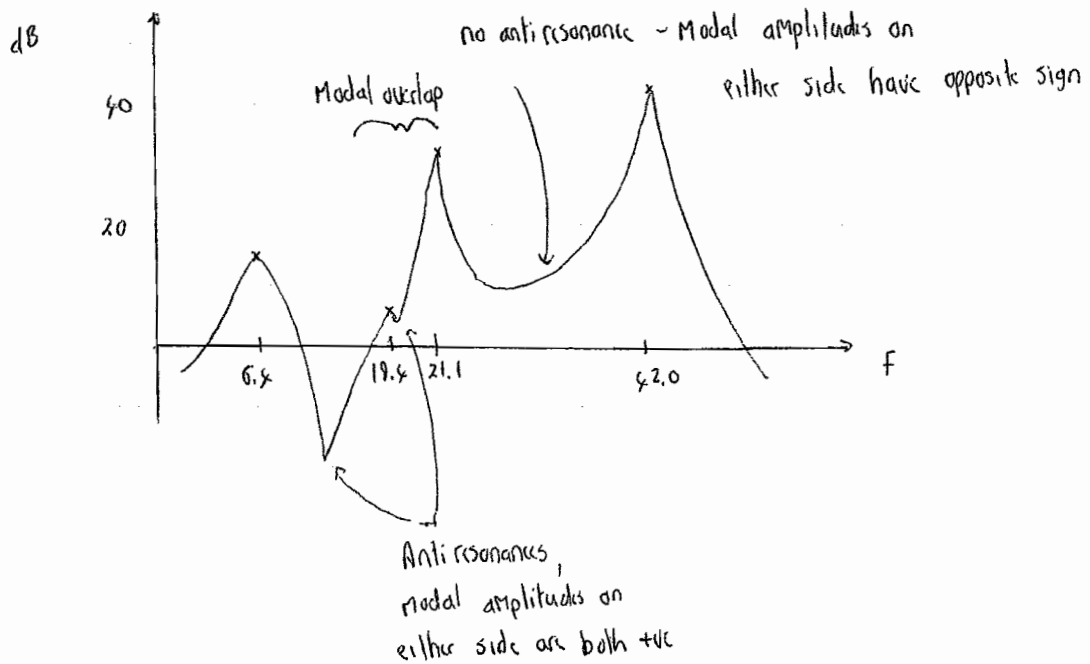


4CG ADVANCED LINEAR VIBRATION - CRTB 2007

1 a)	Mode	$Q = \frac{F_n}{\Delta F}$	$Q a_n$	$20 \log_{10} Q a_n $
	1	21.3	5.3	14 dB
	2	8.1	1.4	3 dB
	3	14.1	34	31 dB
	4	210	-132	42 dB

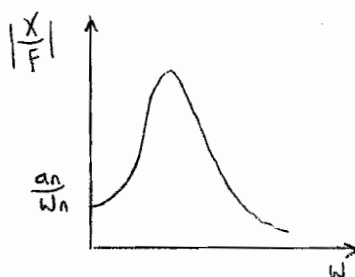
suspect measurement is not a point receptance, since the modal amplitudes should all be +ve for that case

Note Mode 3 lies within the $\frac{1}{2}$ power bandwidth of mode 2 \Rightarrow modal overlap



Note First two modes are relatively heavily damped compared to modes 3 and 4 [30%]

b) Assume mode 1 represents the deflected shape under a static load

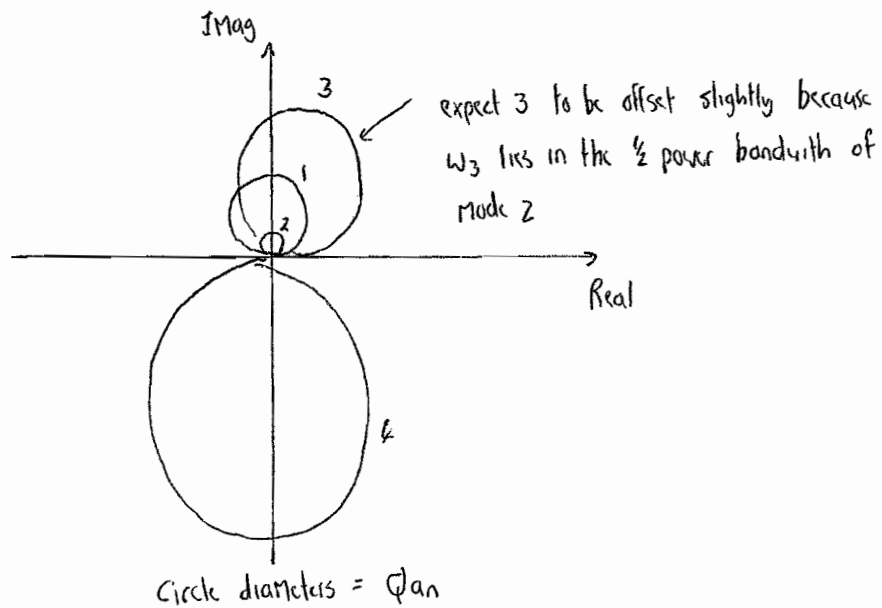


Displacement at $\omega = 0 = \frac{a_1}{\omega_1} = \frac{0.25}{2\pi \times 6.4} = 6.2 \times 10^{-3}$

Stiffness $k = \frac{1}{\text{deflection}} = 161 \text{ N/mm} = \underline{161 \text{ kN/m}}$

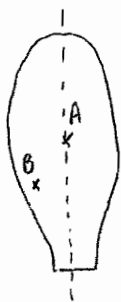
[20%]

c)



[30%]

- d)
- Natural frequencies and Q factors of measured modes unchanged
 - Modal amplitudes expected to change/decrease due to lower response near the floor edge
 - Additional modes may be found, that have a nodal line through A



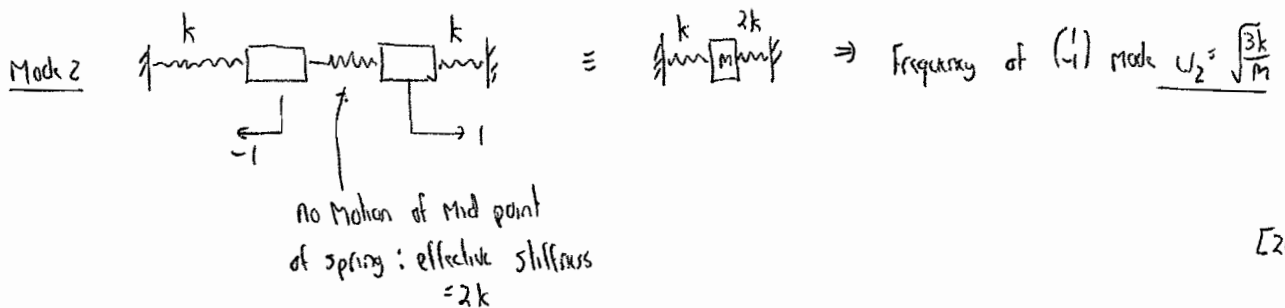
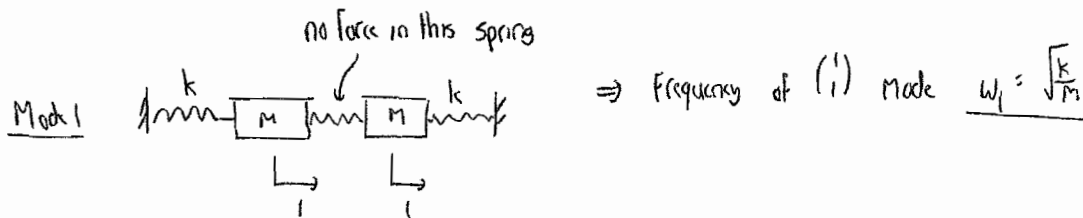
Assuming the floor is near symmetric, it is quite likely that some nodal lines will pass through or near to A - these modes are not "seen" by measurements at A

[20%]

- 2 a) Additional damping comes from
- (i) Friction / microslippage at the structural joints
 - (ii) Air pumping - between the beam and the panel
 - (iii) (Possibly) rattles/impact if there are any loose fittings.
 - (iv) Acoustic radiation - but small below the coincidence frequency
- Together these have a much greater effect than the "material" damping.

If the beams are glued, then (i), (ii) and (iii) will be eliminated and the damping will be reduced, although there could be some additional losses due to hysteresis in the glue [30%]

- b) (i) The structure is symmetric \Rightarrow there will be an antisymmetric mode $\begin{pmatrix} 1 \\ 1 \end{pmatrix}$ and a symmetric mode $\begin{pmatrix} 1 \\ -1 \end{pmatrix}$



(ii) Without damping $k = \begin{pmatrix} 2k & -k \\ -k & 2k \end{pmatrix}$. With damping $k = \begin{pmatrix} 2k & -k \\ -k & k+k(1+i\eta) \end{pmatrix}$

Rayleigh quotient $\frac{u^T k u}{u^T M u} = \omega_n^2 (1+i\eta)_n$

Mode 1 $u^T k u = \begin{pmatrix} 1 & 1 \end{pmatrix} \begin{bmatrix} k \\ k \end{bmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = 2k - k - k + k + k + k i \eta = 2k(1 + \frac{1}{2}i\eta)$
 $u^T M u = 2m$
 $\Rightarrow \omega_n^2 (1+i\eta)_n = \frac{k}{m} (1 + \frac{1}{2}i\eta) \Rightarrow \omega_1 = \frac{1}{2}\omega_2$

Mode 2 $u^T k u = \begin{pmatrix} 1 & -1 \end{pmatrix} \begin{bmatrix} k \\ k \end{bmatrix} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = k + 2k + k + k + k(1+i\eta) = 6k + k i \eta = 6k(1 + \frac{1}{6}i\eta)$
 $u^T M u = 2m \Rightarrow \omega_n^2 (1+i\eta)_n = \frac{3k}{m} (1 + \frac{1}{6}i\eta) \Rightarrow \omega_2 = \frac{1}{2}\omega_1$

[30%]

2cont (iii) In this case we can consider $k \rightarrow k + iC\omega$ rather than $k \rightarrow k + ik\gamma$ for RH spring
 \Rightarrow use previous results with $\gamma \rightarrow C\omega/k$

$$\Rightarrow \gamma_1 = \frac{1}{2}\gamma = \frac{1}{2} C \sqrt{\frac{k}{m}} \frac{1}{k} = \frac{1}{2} \frac{C}{\sqrt{km}}$$

$$\gamma_2 = \frac{1}{2}\gamma = \frac{1}{2} C \sqrt{\frac{3k}{m}} \frac{1}{k} = \frac{1}{2\sqrt{3}} \frac{C}{\sqrt{km}}$$

[20%]

3. a) Method of separation of variables : $w(x,y,t) = X(x)Y(y)T(t)$

$$\frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} - \frac{M}{T} \frac{\partial^2 w}{\partial t^2} = 0$$

$$\Rightarrow X''YT + Y''XT - \frac{M}{T} T''XY = 0$$

$$\Rightarrow \frac{X''}{X} + \frac{Y''}{Y} - \left(\frac{M}{T}\right) \frac{T''}{T} = 0 \quad \text{--- } \textcircled{1}$$

\uparrow \uparrow \uparrow
 $F(x)$ $F(y)$ $F(t)$

$$\Rightarrow \left. \begin{aligned} \frac{X''}{X} &= \text{const} = -\alpha^2 \text{ say} \Rightarrow X = A \cos \alpha x + B \sin \alpha x \\ \frac{Y''}{Y} &= \text{const} = -\beta^2 \text{ say} \Rightarrow Y = C \cos \beta y + D \sin \beta y \\ \frac{T''}{T} &= \text{const} = -\omega^2, \text{ harmonic motion} \end{aligned} \right\} \alpha^2 + \beta^2 = \left(\frac{M}{T}\right) \omega^2$$

to satisfy Eqn $\textcircled{1}$

Boundary conditions $X(0) = 0 \Rightarrow X = B \sin \alpha x$

$$X(L_1) = 0 \Rightarrow \sin \alpha L_1 = 0 \Rightarrow \alpha = \frac{n\pi}{L_1}$$

$$Y(0) = 0 \Rightarrow Y = D \sin \beta y$$

$$Y(L_2) = 0 \Rightarrow \sin \beta L_2 = 0 \Rightarrow \beta = \frac{m\pi}{L_2}$$

$$\Rightarrow \text{Mode shape} = \frac{\sin \frac{n\pi x}{L_1} \sin \frac{m\pi y}{L_2}}$$

$$\text{and } \omega^2 = \left(\frac{T}{M}\right) (\alpha^2 + \beta^2) = \frac{T}{M} \left[\left(\frac{n\pi}{L_1}\right)^2 + \left(\frac{m\pi}{L_2}\right)^2 \right]$$

[60%]

b) ω_1 has $n=1, m=1 \Rightarrow \omega_1^2 = \frac{T}{M} \left[\frac{1}{L_1^2} + \frac{1}{L_2^2} \right] \pi^2$

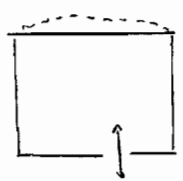
Put $L_2 = (L_1/A)^{-1} \Rightarrow \omega_1^2 = \left(\frac{T}{M}\right) \pi^2 \left[\frac{1}{L_1^2} + \frac{L_1^2}{A^2} \right]$

For minimum $\frac{d}{dL_1} = 0 \Rightarrow -2L_1^{-3} + 2L_1/A^2 = 0 \Rightarrow L_1^4 = A^2$
 $\Rightarrow \underline{L_1 = L_2}$

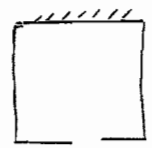
Square membrane has lowest ω_1 , with $\omega_1 = \sqrt{\frac{T}{M}} \pi \frac{\sqrt{2}}{\sqrt{A}}$

Circular membrane has $\omega_1 = \frac{2.4}{a} \sqrt{\frac{T}{M}} = \sqrt{\frac{T}{M}} \frac{2.4}{\sqrt{A}} \sqrt{\pi} \Rightarrow \frac{\omega_{\text{square}}}{\omega_{\text{circle}}} = \frac{\pi \sqrt{2}}{2.4 \sqrt{\pi}} = \underline{1.046}$

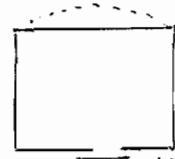
3: c) Consider the coupled system and two constrained systems



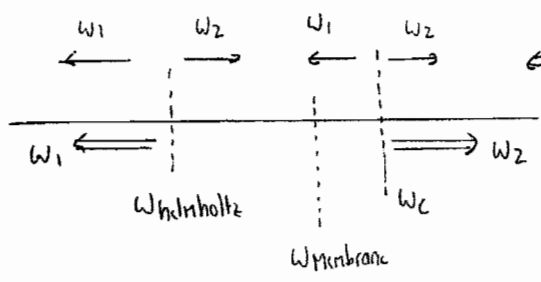
original coupled system: ω_1 and ω_2



constrained system has $\omega_n = \omega_{helmholtz}$



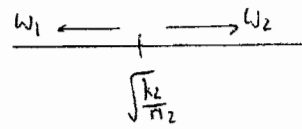
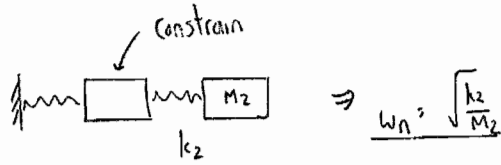
constrained system has $\omega_n > \omega_{membrane}$ due to added acoustic stiffness
Say $\omega_n = \omega_c$



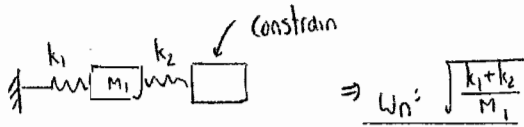
← application of the interlacing theorem

← location of coupled frequencies

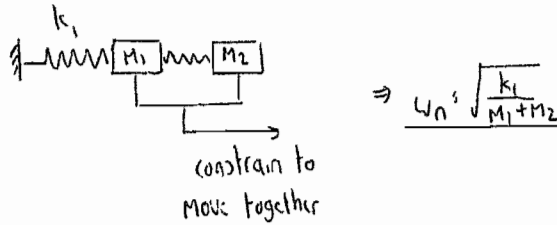
4 a) (i)



Frequency after constraint must interlace with frequencies before constraint



Again, $\omega_1 \leq \omega_n \leq \omega_2$



Again, $\omega_1 \leq \omega_n \leq \omega_2$

Thus $\omega_1 \leq \sqrt{\frac{k_2}{M_2}} \leq \sqrt{\frac{k_1+k_2}{M_1}} \leq \sqrt{\frac{k_1}{M_1+M_2}} \leq \omega_2$

[20%

(ii) For $\omega_1 = \omega_2$ we must have

$$\begin{matrix} \text{①} & \frac{k_2}{M_2} & = & \frac{k_1+k_2}{M_1} & = & \frac{k_1}{M_1+M_2} & \text{③} & \text{①} = \text{③} \Rightarrow & k_1 = k_2 \left(1 + \frac{M_1}{M_2}\right) \\ & \text{②} & & & & & & \text{①} = \text{②} \Rightarrow & k_1+k_2 = k_2 \left(\frac{M_1}{M_2}\right) \end{matrix}$$

Thus $k_2 \left(1 + \frac{M_1}{M_2}\right) + k_2 = k_2 \left(\frac{M_1}{M_2}\right) \Rightarrow 2k_2 = 0 \Rightarrow \underline{k_1 = k_2 = 0}$

$\omega_1 = \omega_2$ only possible when $\underline{\omega_1 = \omega_2 = 0}$

[20%

(iii) Equation of motion

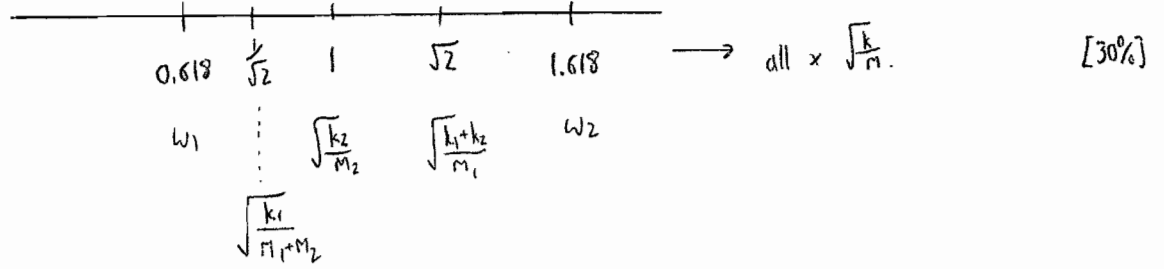
$$\begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \begin{pmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{pmatrix} + \begin{pmatrix} 2k & -k \\ -k & k \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$| -\omega^2 m + k | = 0 \Rightarrow \begin{vmatrix} -\omega^2 m + 2k & -k \\ -k & -\omega^2 m + k \end{vmatrix} = 0 \Rightarrow (-\omega^2 m + 2k)(-\omega^2 m + k) - k^2 = 0$$

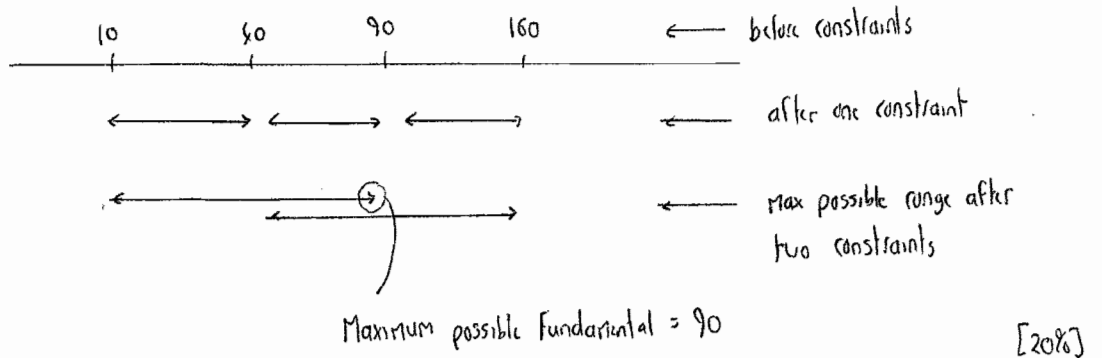
$$\Rightarrow \omega^4 m^2 - 3\omega^2 m k + k^2 = 0$$

$$\Rightarrow \omega_{1,2} = \frac{3mk \pm \sqrt{9m^2 k^2 - 4m^2 k^2}}{2m^2} = \frac{k}{m} \left(\frac{3 \pm \sqrt{5}}{2} \right) \Rightarrow \omega_1 = \frac{k}{m} \left(\frac{3 + \sqrt{5}}{2} \right) \quad \omega_2 = \frac{k}{m} \left(\frac{3 - \sqrt{5}}{2} \right)$$

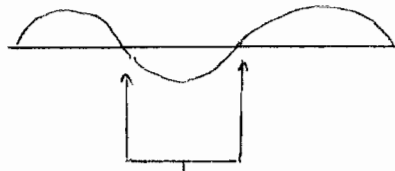
4 cont



b) (i) Treat the stiff springs as constraints



(iii) $90 = \omega_3$: mode shape:



put constraints here i.e. at $L/3$ and $2L/3$ [10%]