

ENGINEERING TRIPOS PART IIB

Tuesday 24 April 2007 9 to 10.30

Module 4A8

ENVIRONMENTAL FLUID MECHANICS

*Answer not more than **three** questions.*

All questions carry the same number of marks.

*The **approximate** percentage of marks allocated to each part of a question is indicated in the right margin.*

Attachment: 4A8 Data Card (5 pages).

STATIONERY REQUIREMENTS

Single-sided script paper

SPECIAL REQUIREMENTS

Engineering Data Book

CUED approved calculator allowed

**You may not start to read the questions
printed on the subsequent pages of this
question paper until instructed that you
may do so by the Invigilator**

1 A heavy gas is released on a slope of angle θ to the horizontal and runs down as a thickening layer. You may assume that the flow is homogeneous across the slope and hence the mean flow is essentially two-dimensional. The flow may be considered to be turbulent and the mixing between the gas and the surrounding air is characterised by an entrainment coefficient α . At a distance s down the slope the velocity of the flow is uniform across the thickness of the layer and equal to $U \text{ m s}^{-1}$. The thickness of the layer (normal to the ground) is $h \text{ m}$ and the mean density of the layer is $\rho \text{ kg m}^{-3}$. The air outside the layer is still and its density is $\rho_a \text{ kg m}^{-3}$. The density difference is very small i.e. $(\rho - \rho_a)/\rho_a \ll 1$ so that ρ may be replaced by ρ_a where appropriate.

(a) Derive the equations governing the behaviour of this flow. [20%]

(b) Using the equations for conservation of volume and mass show that

$$\frac{d((\rho - \rho_a)Uh)}{ds} = 0$$

[20%]

(c) Assuming that the parameters describing the layer have a power-law variation with distance down the slope:

(i) Find the variation of the thickness h with distance down the slope. [20%]

(ii) Find the variation of the density difference $(\rho - \rho_a)$ with distance down the slope. [20%]

(iii) Find the variation of the velocity U with distance down the slope. [20%]

2 A steady free convection flow is set up near the ground on a hot day due to the mean heat flux Q from the ground to the air above. The mean velocity is zero. Measurements show that the length-scale of the largest eddying motions is approximately L . The mean temperature of the air is T and it has a heat capacity c_p and a density ρ . The acceleration due to gravity is g . The enthalpy equation for this particular situation is

$$\lambda \frac{\partial T}{\partial x_3} - \rho c_p \overline{u_3 \theta} = -Q$$

where λ is the thermal conductivity of the air, x_3 is the vertical direction, u_3 is the vertical velocity fluctuation, θ is the temperature fluctuation and the overbar denotes a time average. The equation for the turbulent kinetic energy is given on the Data Card.

- (a) Explain the meaning of the two terms on the left-hand side of this equation and identify the regions of the flow where each one dominates and explain why. [20%]
- (b) By considering the dominant terms in the turbulent kinetic energy equation and using the above equation estimate the size of the smallest eddying motion in the turbulent flow. [20%]
- (c) This flow is to be computed using Direct Numerical Simulation. Find an expression for the approximate number of grid points required to accurately compute this flow in terms of Q , g , L , ρ , c_p , T and ν , where ν is the kinematic viscosity of the air. [20%]
- (d) Find an expression for the time-scale of the largest eddying motions in the flow (the large eddy turnover time) in terms of the same variables. [20%]
- (e) Estimate the number of time-steps required in the simulation to capture 10 large eddy turnover times as a function of the same variables. [20%]

(TURN OVER

3 (a) What are the types of pollutant emission sources according to geometrical characterisation? Give an example for each type. [5%]

(b) Explain how to include ground effects in the Gaussian Plume Model for pollutant dispersion. [10%]

(c) Consider a line source at a distance y_0 metres above the ground level emitting (Q/L) kg of a pollutant per unit time and per unit length. The pollutant spreads along y . There is a mean wind at a speed of U m s⁻¹ along the x -direction. Using the solutions given on the data card,

(i) Show that, when the ground is non-absorbing, the pollutant concentration on the ground varies as

$$\bar{\phi}(x, 0) = \frac{A}{\sigma_y} \exp\left(-\frac{y_0^2}{2\sigma_y^2}\right),$$

where $A = 2Q/(UL\sqrt{2\pi})$ and σ_y is the dispersion coefficient. [20%]

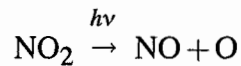
(ii) By taking $\sigma_y = Bx$ where B is a constant, show that the maximum pollutant concentration occurs at $x = y_0/B$ and the maximum concentration is equal to $0.607A/y_0$. [50%]

(d) Calculate the maximum ground concentration if the line source is located at 10 m above the ground level, the pollutant emission rate is 1 kg m⁻¹ s⁻¹ and the mean wind speed is 10 km hr⁻¹ [10%]

(e) What happens to the maximum ground level concentration if the height of the pollutant source is doubled? [5%]

4 (a) Explain the physical significance of the Damköhler number, Da , and its various limits. [20%]

(b) The mean concentration $\bar{\phi}$ of NO_2 in the air above a city at 5:00 AM is measured to be $\bar{\phi}_0 \text{ kg m}^{-3}$ and the concentration is considered to be homogeneous in the mean but finite small-scale fluctuations exist. The turbulence in the city atmosphere has a characteristic timescale T_{turb} . The weather is clear and still. Thus, one would expect the following photolytic reaction with a rate constant $k \text{ s}^{-1}$



to occur during the day.

(i) Calculate the time, in terms of k , for the mean NO_2 concentration to become 5% of its initial value. Take k to be independent of the time of the day. [35%]

(ii) From the modelled variance equation given in the data card, obtain an expression for the variation of NO_2 variance, g , with time. Sketch the variation of g/g_0 with time for $k = 0$, $k \gg 1$, and $k < 1$. g_0 is the initial variance. [45%]

END OF PAPER

4A8: Environmental Fluid Mechanics

DATA CARD

(5 pages)

Part I: Turbulence and Fluid Mechanics

Rotating Flows

Geostrophic Flow

$$-\frac{1}{\rho} \nabla p = 2\Omega \times \underline{u}$$

Ekman Layer Flow

$$-2\Omega_z v = -\frac{1}{\rho} \frac{\partial p}{\partial x} + v \frac{\partial^2 u}{\partial z^2}$$

$$2\Omega_z u = -\frac{1}{\rho} \frac{\partial p}{\partial y} + v \frac{\partial^2 v}{\partial z^2}$$

$$0 = -\frac{1}{\rho} \frac{\partial p}{\partial z}$$

OR

$$-2\Omega_z v = v \frac{\partial^2 u}{\partial z^2}$$

$$-2\Omega_z v(u_0 - u) = v \frac{\partial^2 v}{\partial z^2}$$

GEOSTROPHIC VELOCITY

Solution is

$$u = u_0 \left[1 - e^{-z/\Delta} \cos \frac{z}{\Delta} \right]$$

$$v = u_0 e^{-z/\Delta} \sin \frac{z}{\Delta}$$

$$\Delta = \left(\frac{v}{\Omega_z} \right)^{1/2}$$

Turbulent Flows – Incompressible

Continuity Equation $\nabla \cdot \underline{U} = \frac{\partial U_i}{\partial x_i} = 0$

Momentum Equation $\rho \frac{DU}{Dt} = -\nabla P + \mu \nabla^2 \underline{U} + \underline{F}$

$$\rho \frac{DU_i}{Dt} = -\frac{\partial P}{\partial x_i} + \mu \frac{\partial^2 U_i}{\partial x_j^2} + F_i$$

Enthalpy Equation $\rho c_p \frac{DT}{Dt} = -k \frac{\partial^2 T}{\partial x_i^2}$

Reynolds Transformation $U_i = \overline{U}_i + u_i$ etc

Reynolds Stress $= -\overline{\rho u_i u_j}$

Reynolds Heat Flux $= -\overline{\rho c_p u_j \theta}$

Turbulent Kinetic Energy Equation $\frac{D}{Dt} \frac{q^2}{2} = -\overline{u_i u_k} \frac{\partial \overline{U}_i}{\partial x_k} - \varepsilon + \frac{\overline{f_i u_i}}{\rho} + \text{transport terms}$

$$\varepsilon = \overline{v \left(\frac{\partial u_i}{\partial x_k} + \frac{\partial u_k}{\partial x_i} \right) \left(\frac{\partial u_i}{\partial x_k} \right)}, \quad \frac{q^2}{2} = \frac{1}{2} (u_1^2 + u_2^2 + u_3^2)$$

In flows with thermally driven motion

$$\frac{f_i u_i}{\rho} = \frac{g}{T} \cdot \overline{\theta u_i}$$

$i = \text{Vertical direction}$

Dissipation of turbulent kinetic energy $\varepsilon \approx \frac{u'^3}{\ell}$

Kolmogorov microscale $\eta = \left(\frac{v^3}{\varepsilon} \right)^{1/4}$

Taylor microscale (λ) $\varepsilon = 15v \frac{u'^2}{\lambda^2}$ (v is the kinematic viscosity)

Density Influenced Flows

Atmospheric Boundary Layer

$$\left. \frac{dT}{dz} \right|_{\text{NEUTRAL STABILITY}} = -\frac{g}{C_p} = \left. \frac{dT}{dz} \right|_{\text{DALR}}$$

$$R_i = \frac{g}{T} \frac{\left. \frac{dT}{dz} - \left. \frac{dT}{dz} \right|_{\text{DALR}} \right|}{\left(\left. \frac{dU}{dz} \right)^2} = \text{RICHARDSON NUMBER}$$

Neutral Stability

$$U = \frac{u_*}{\kappa} \ln \frac{z}{z_0}; \quad \frac{dU}{dz} = \frac{u_*}{\kappa z}$$

$$u_* = \sqrt{\frac{\tau_w}{\rho}}; \quad \kappa = \text{von Kármán Constant} = 0.40$$

Non-Neutral Stability

$$L = \text{Monin-Obukhov length} = -\frac{u_*^3}{\kappa \frac{g}{T} \frac{Q}{\rho c_p}}$$

Q = surface heat flux

$$\frac{dU}{dz} = \frac{u_*}{\kappa z} \left(1 - 15 \frac{z}{L} \right)^{-1/4} \quad \text{Unstable}$$

$$= \frac{u_*}{\kappa z} \left(1 + 4.7 \frac{z}{L} \right) \quad \text{Stable}$$

Buoyant plume for a point source

$$\frac{d}{dz} (\pi R^2 w) = 2\pi R u_e \quad (i)$$

$$\frac{d}{dz} (\rho \pi R^2 w) = \rho_a 2\pi R u_e \quad (ii)$$

$$\frac{d}{dz} (\rho \pi R^2 w^2) = g(\rho_a - \rho) \pi R^2 \quad (iii)$$

(i) and (iii) give

$$\pi R^2 w \left(\frac{\rho_a - \rho}{\rho_a} \right) g = \text{constant} = F_0 \text{ (buoyancy flux)}$$

$$u_e = \alpha w$$

(α = Entrainment coefficient)

Buoyancy Frequency

$$N^2 = -\frac{g}{\rho} \frac{d\rho}{dz} = \frac{g}{T} \frac{dT}{dz}$$

Actually $\frac{g}{T} \left(\frac{dT}{dz} - \frac{dT}{dz} \Big|_{\text{DALR}} \right)$

N = Brunt – Väisälä Frequency or Buoyancy Frequency

Part II: Dispersion of Pollution in the Atmospheric Environment

Transport equation for the mean of the reactive scalar ϕ :

$$\frac{\partial \bar{\phi}}{\partial t} + \bar{u}_j \frac{\partial \bar{\phi}}{\partial x_j} = \frac{\partial}{\partial x_j} \left(K \frac{\partial \bar{\phi}}{\partial x_j} \right) + \bar{w}$$

Transport equation for the variance of the reactive scalar ϕ :

$$\frac{\partial g}{\partial t} + \bar{u}_j \frac{\partial g}{\partial x_j} = \frac{\partial}{\partial x_j} \left(K \frac{\partial g}{\partial x_j} \right) + 2K \left(\frac{\partial \bar{\phi}}{\partial x_j} \right)^2 - \frac{2}{T_{turb}} g + 2\overline{\phi'w'}$$

Mean concentration of pollutant after instantaneous release of Q kg at $t=0$:

$$\bar{\phi}(x, y, z, t) = \frac{Q}{8(\pi t)^{3/2} (K_x K_y K_z)^{1/2}} \exp \left[-\frac{1}{4t} \left(\frac{(x-x_0)^2}{K_x} + \frac{(y-y_0)^2}{K_y} + \frac{(z-z_0)^2}{K_z} \right) \right]$$

Gaussian plume spreading in two dimensions from a source at $(0,0,z_0)$ emitting Q kg/s:

$$\bar{\phi}(x, y, z) = \frac{Q}{2\pi U \sigma_y \sigma_z} \exp \left[-\left(\frac{y^2}{2\sigma_y^2} + \frac{(z-z_0)^2}{2\sigma_z^2} \right) \right]$$

One-dimensional spreading from line source at $y=0$ emitting Q/L kg/s/m :

$$\bar{\phi}(x, y) = \frac{Q}{UL} \frac{1}{\sqrt{2\pi}\sigma_y} \exp \left(-\frac{y^2}{2\sigma_y^2} \right)$$

Relationship between eddy diffusivity and dispersion coefficient:

$$\sigma^2 = 2 \frac{x}{U} K$$