

①

Geostrophic flow equation

$$2 \underline{\underline{\Omega}} \times \underline{\underline{u}} = - \frac{1}{\rho_0} \nabla P \quad \text{---} \quad \text{①}$$

(a) The pressure gradient force is balanced by the Coriolis force in a rotating frame of reference.

~~No~~ No, this equation needs to be modified by including viscous forces, which are important inside the boundary layers. The modified equation is given in the data card as Ekman layer flow equation.

$$(b) \quad \underline{\underline{\omega}} = \nabla \times \underline{\underline{u}} \Rightarrow \omega_z = \left( \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right)$$

① in component form

$$2 \Omega_z v = \frac{1}{\rho_0} \frac{\partial P}{\partial x} \quad \text{and} \quad 2 \Omega_z u = - \frac{1}{\rho_0} \frac{\partial P}{\partial y}$$

Taking appropriate derivatives of these two components

$$\omega_z = \frac{1}{2 \Omega_z \rho_0} \left( \frac{\partial^2 P}{\partial x^2} + \frac{\partial^2 P}{\partial y^2} \right) = \frac{1}{2 \Omega_z \rho_0} \nabla^2 P.$$

Thus,  $\omega_z$  is related to  $\nabla^2 P$ , the Pressure Laplacian.

$$\nabla \cdot \underline{u} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y}$$

$$= \frac{1}{2\Omega_z \rho_0} \left( -\frac{\partial^2 p}{\partial x \partial y} + \frac{\partial^2 p}{\partial x \partial y} \right) = 0$$

∵  $p$  is continuous

(1)

$$\frac{\partial p}{\partial z} = -\rho g = -\rho_0 (1 - \alpha \theta) g$$

Taking  $z$ -derivative of the components of (1)

$$2\Omega_z \frac{\partial v}{\partial z} = \frac{1}{\rho_0} \frac{\partial}{\partial x} \left( \frac{\partial p}{\partial z} \right) = g\alpha \frac{\partial \theta}{\partial x}$$

$$2\Omega_z \frac{\partial u}{\partial z} = -\frac{1}{\rho_0} \frac{\partial}{\partial y} \left( \frac{\partial p}{\partial z} \right) = -g\alpha \frac{\partial \theta}{\partial y}$$

Thus, the required equations for the thermal wind are

$$\boxed{2\Omega_z \frac{\partial v}{\partial z} = g\alpha \frac{\partial \theta}{\partial x} \quad \& \quad 2\Omega_z \frac{\partial u}{\partial z} = -g\alpha \frac{\partial \theta}{\partial y}}$$

The vertical shear of horizontal velocities ~~are~~ is along the isotherms, and its magnitude is given by the horizontal gradient of the temperature increment  $\theta$  ( $\equiv (T - T_0)$ )

$$\theta = Ax^2 + By$$

$$\Rightarrow \frac{\partial \theta}{\partial x} = 2Ax ; \quad \frac{\partial \theta}{\partial y} = B$$

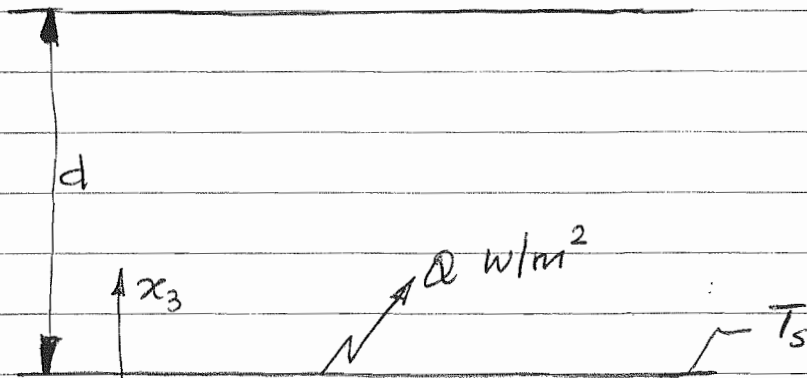
$$\therefore \frac{\partial v}{\partial z} = \frac{g\alpha}{\Omega_z} Ax \Rightarrow \boxed{v = \frac{g\alpha A}{\Omega_z} xz + C_1(x,y)}$$

$$\frac{\partial u}{\partial z} = -\frac{g\alpha}{2\Omega_z} B \Rightarrow \boxed{u = \frac{g\alpha B}{2\Omega_z} z + C_2(x,y)}$$

To determine  $C_1(x,y)$  &  $C_2(x,y)$  one needs to know  $u$  &  $v$  at some height  $z_1$ . Although  $u$  &  $v$  need to satisfy no slip condition at  $z=0$ ,  $C_1$  &  $C_2$  can not be zero since viscous effects are not included in the analysis.

The velocities  $u$  &  $v$  varies linearly with vertical distance  $z$ .

②



a) Mean temperature equation

$$\rho c_p \frac{D\bar{T}}{Dt} = \nabla \cdot (\lambda \nabla \bar{T} - \overline{u\theta} \rho c_p)$$

$$\left. \begin{array}{l} \text{Stationary flow} \\ \text{no mean motion (Velocities)} \end{array} \right\} \Rightarrow \frac{D\bar{T}}{Dt} = 0$$

$$\Rightarrow \lambda \frac{d\bar{T}}{dx_3} - \overline{u_3\theta} \rho c_p = -Q$$

Neglect molecular heat flux

$$\Rightarrow \boxed{\overline{u_3\theta} = \frac{Q}{\rho c_p}}$$

b) Turbulent Kinetic energy equation

$$\rho \frac{DK}{Dt} = -\rho \overline{u_i u_k} \frac{\partial \bar{u}_i}{\partial x_k} - \rho \epsilon + \overline{f_i u_i}$$

No mean velocities & stationary

$$\Rightarrow \boxed{\epsilon = \frac{\overline{f_3 u_3}}{\rho}} = \frac{g}{T} \overline{u_3 \theta}$$

dissipation balances  
buoyancy

Variance equation:

$$\frac{D\theta}{Dt} = -2 \overline{u_j \theta} \frac{\partial \bar{T}}{\partial x_j} - \epsilon_\theta + \text{molecular diffusion}$$

$$\Rightarrow \boxed{\epsilon_\theta = -2 \overline{u_3 \theta} \frac{d\bar{T}}{dx_3}}$$

dissipation balances  
production

$$\epsilon \approx \frac{g}{T} \overline{u_3 \theta} \Rightarrow \frac{u^3}{L} \sim \frac{g}{T} \overline{u_3 \theta}$$

$$\Rightarrow u^2 \sim \left( \frac{g}{T} L \overline{u_3 \theta} \right)^{2/3}$$

$$\boxed{K \sim \left( \frac{g}{T} d \overline{u_3 \theta} \right)^{2/3} \text{ or } \left( \frac{g}{T} \frac{d}{\rho c_p} Q \right)^{2/3}}$$

$$\boxed{\epsilon \sim \frac{g}{T} \frac{Q}{\rho c_p}}$$

c)

$$\epsilon_\sigma \approx -2U_3 \theta \frac{dT}{dx_3}$$

$$C_d \sigma \frac{\epsilon}{k} \approx -2 \frac{Q}{\rho C_p} \frac{dT}{dx_3}$$

from (b)  $\epsilon/k \approx \frac{1}{d} \left( \frac{g d Q}{T \rho C_p} \right)^{1/3}$

$$\Rightarrow C_d \sigma \frac{1}{d} \left( \frac{g d Q}{\rho C_p} \right)^{1/3} T^{-4/3} = -2 \left( \frac{Q}{\rho C_p} \right) \frac{dT}{dx_3}$$

$$-A dx_3 = T^{1/3} dT$$

where  $A \equiv \frac{C_d \sigma}{2} \left( \frac{g}{d^2} \right)^{1/3} \left( \frac{\rho C_p}{Q} \right)^{2/3}$

integrate

$$\frac{3}{4} T^{4/3} \Big|_{T_s}^{\bar{T}} = -A x_3$$

$$\bar{T}^{4/3} - T_s^{4/3} = -\frac{4}{3} A x_3$$

$$\frac{\bar{T}}{T_s} = (1 - A^* x_3)^{3/4} \quad A^* \equiv \frac{4}{3} \frac{A}{T_s^{4/3}}$$

@  $x_3 = d/2$

$$\frac{\bar{T}}{T_s} = (1 - A d)^{3/4}$$

$$A = A^* / 2 = \frac{C_d \sigma}{3} \frac{1}{T_s^{4/3}} \left( \frac{g}{d^2} \right)^{1/3} \left( \frac{\rho C_p}{Q} \right)^{2/3}$$

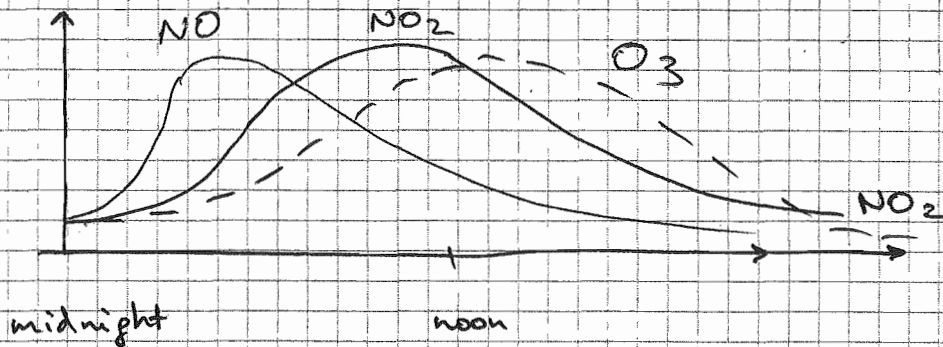
Q3

(a)  $\text{NO}$ : emitted from transport & power generation

$\text{O}_3$ : not emitted, exists in background & created by photochemistry

$\text{NO}_2$ : emitted little, created by  $\text{NO} + \text{O}_3 \rightarrow \text{NO}_2 + \text{O}$   
destroyed by photolysis  $\text{NO}_2 \rightarrow \text{NO} + \text{O}$

$\text{O} + \text{O}_2 \rightarrow \text{O}_3$ , which further reacts with  $\text{NO}$ .

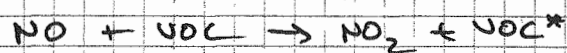


$\text{O}_3$  is irritating for respiratory system and causes impaired vision.

$\text{NO}_2$  (brown gas) causes lung irritation at high concentrations and contributes to acid rain

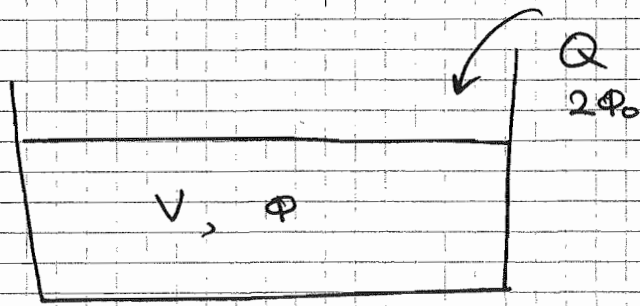
$\text{NO}$  could cause liver damage in high doses

VOC: emitted by transport, industry, domestic sector.  
participates in smog formation



Can be toxic, carcinogenic, irritant.

(b)



- Water conservation:  $\frac{dV}{dt} = Q \Rightarrow V = V_0 + Qt \quad (1)$

- Pollutant conservation:  $\frac{d(\phi V)}{dt} = Q \cdot 2\phi_0$

$\Rightarrow \phi \frac{dV}{dt} + V \frac{d\phi}{dt} = Q \phi_1 \quad (\phi_1 = 2\phi_0)$

$\Rightarrow (V_0 + Qt) \frac{d\phi}{dt} = Q \phi_1 - \phi \cdot Q$

$\Rightarrow \frac{d\phi}{\phi_1 - \phi} = \frac{dt}{\frac{V_0}{Q} + t}$

$\Rightarrow -\ln(\phi_1 - \phi) = \ln\left(\frac{V_0}{Q} + t\right) + C$

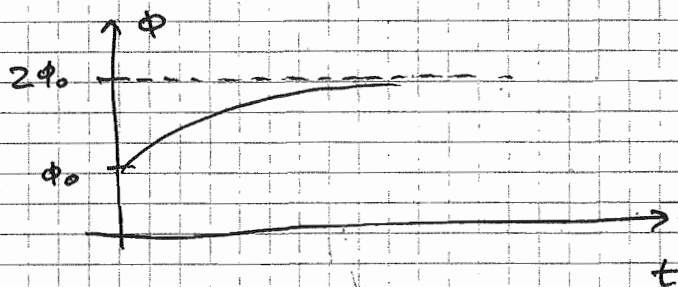
$\Rightarrow \phi_1 - \phi = A \left(\frac{V_0}{Q} + t\right)^{-1}$

At  $t=0, \phi = \phi_0$

$\& \text{ using } \phi_1 = 2\phi_0, \phi_0 = A \frac{Q}{V_0} \Rightarrow A = \phi_0 \frac{V_0}{Q}$

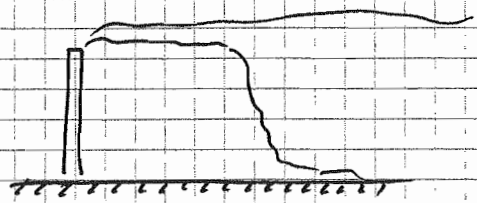
$\therefore 2\phi_0 - \phi = \frac{\phi_0 \cdot V_0}{Q} \left(\frac{V_0}{Q} + t\right)^{-1}$

$\Rightarrow \frac{\phi}{\phi_0} = 2 - \left(1 + \frac{tQ}{V_0}\right)^{-1} \quad \text{Q.E.D.}$

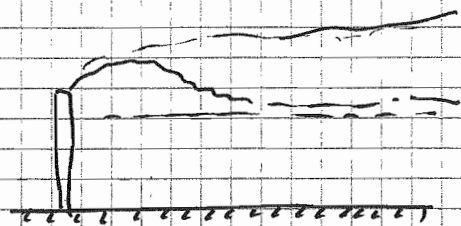




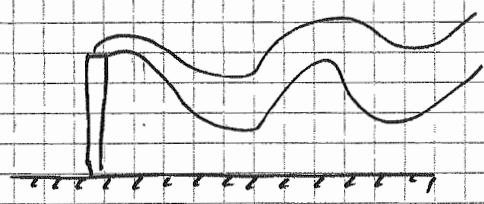
Q4  
(a)



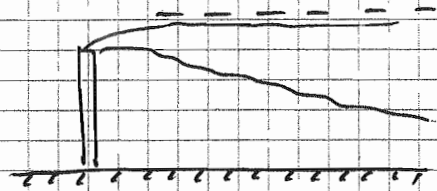
Fumigation: when mixing height increases at dawn and reaches undiluted plume



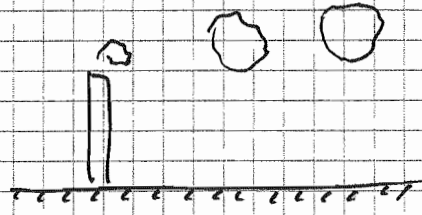
Lofting: source above inversion, plume avoids ground



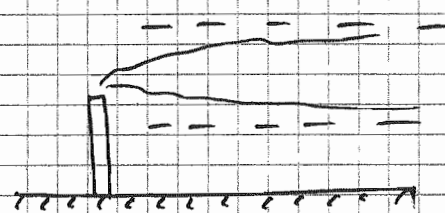
Looping: large-scale motion, vigorous ~~stables~~ turbulence neutral or unstable



Trapped plume: source below inversion, spreads to the ground



Thermalling: very buoyant plume, breaks apart



Fanning: plume between two inversions

$$(b) \quad U \frac{\partial \phi}{\partial x} = K \frac{\partial^2 \phi}{\partial y^2} \quad (1)$$

If  $\phi = \frac{Q}{U \sigma \sqrt{2\pi}} \exp\left(-\frac{y^2}{2\sigma^2}\right)$   $Q$ : kg/s/m of source

Then:  $\frac{\partial \phi}{\partial y} = -2y \phi \frac{1}{2\sigma^2} = -\frac{y}{\sigma^2} \phi \Rightarrow \frac{\partial^2 \phi}{\partial y^2} = -\frac{\phi}{\sigma^2} \left(1 - \frac{y^2}{\sigma^2}\right)$

$$\frac{\partial \phi}{\partial x} = \frac{\partial \phi}{\partial \sigma} \frac{d\sigma}{dx}$$

$$= \left[ -\frac{\phi}{\sigma} + \phi \left(\frac{y^2}{\sigma^3}\right) \right] \frac{d\sigma}{dx} = -\frac{\phi}{\sigma} \left(1 - \frac{y^2}{\sigma^2}\right) \frac{d\sigma}{dx}$$

Q4 cont'd

$$C \frac{\partial \phi}{\partial x} = K \frac{\partial^2 \phi}{\partial y^2} \Rightarrow C \left( -\frac{\phi}{\sigma} \right) \left( 1 - \frac{y^2}{\sigma^2} \right) \frac{d\sigma}{dx} = K \left( -\frac{\phi}{\sigma^2} \right) \left( 1 - \frac{y^2}{\sigma^2} \right)$$

$$\text{①} \quad \sigma \frac{d\sigma}{dx} = \frac{C}{K}$$

$$\text{②} \quad \frac{d\sigma^2}{dx} = 2 \frac{C}{K}$$

This is the condition for the Gaussian plane equation to be a solution to Eq (1).

If  $K = \text{constant}$  and  $\sigma = 0$  at  $x = 0$ ,

$$\sigma^2 = 2 \frac{C}{K} x$$