

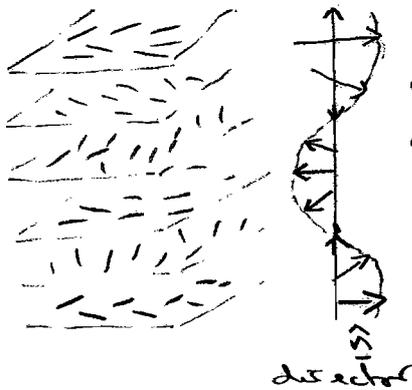
Nematic



Long range orientational order, no correlations between centres of mass Preferred or 'average' pointing direction defined by the director  $\hat{n}$  where  $\hat{n} = -\hat{n}$  (2)  
Molecules anisotropic and phase is fluid thus

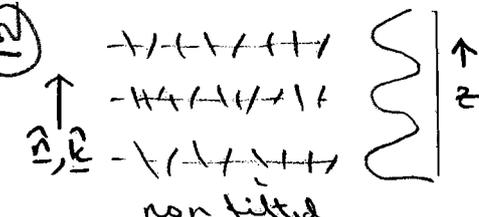
molecules can slide past one another whilst maintaining 'average' parallelism. Least ordered of the thermotropic phases.

Chiral Nematic

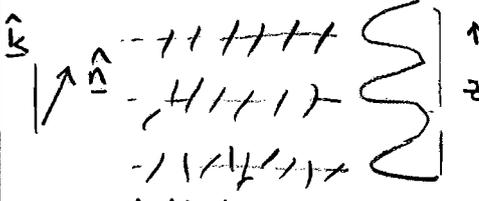


Similar to nematic except a 'screw axis' superimposed on the director. Molecules are all chiral, low twist energy ( $\sim 10^{-5}$  of total associated with  $\hat{n}$ ). Molecules sit in 'pseudo' layers in which each layer has the order of a non chiral nematic. Spiral axis for  $\hat{n}$  gives rise to selective reflector of polarised light. Rotatory power  $\sim 10^3 \times$  sugar solution. (3)

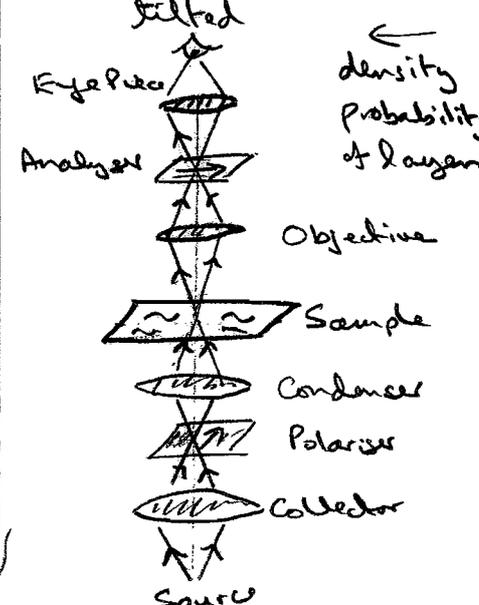
(12)



Stratified layer like structure (probability density wave). In non tilted SA phase  $\hat{n}$  and  $\hat{k}$  (layer normal) parallel. In tilted



SA phase layer normal and  $\hat{n}$  at an angle. layer thickness  $\sim 1$  molecular length (3)  
Layers slide relative to each other  $\therefore$  fluid



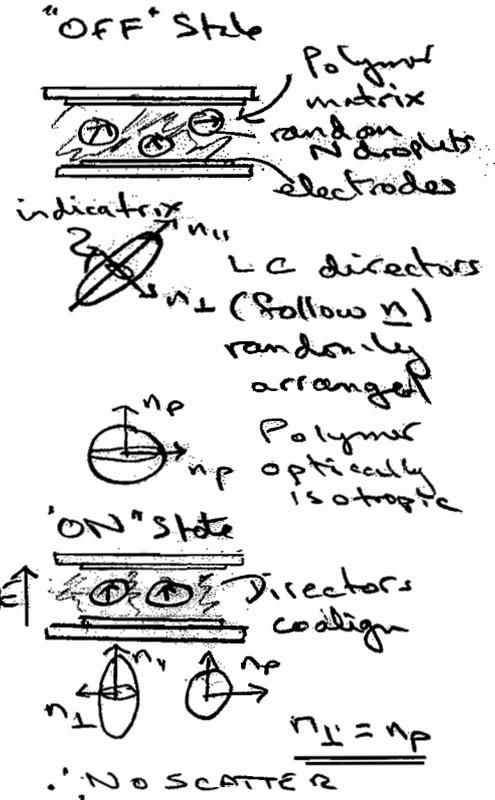
For SA, SC there is no structure in the 2D layer. Higher order phases, SB etc for 2D packing structures, i.e. hexagonal, etc

Phases identified via polarising microscopy (or DSC or X-Ray). Orientation of phase gives different optical textures characteristic of each phase (4)

(2)

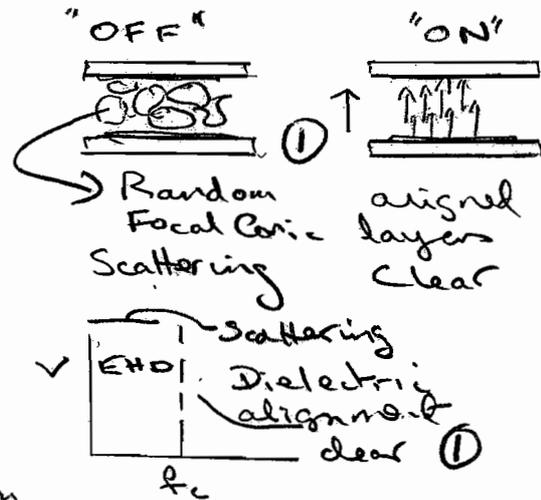
2(b) continued.

(i) Nematic scattering device:  
 Use PDLC (Polymer Dispersed LCs) to create matched ref. indices in the "ON" state, scatter light in the "OFF" state. Polymer matrix  $n_o = n_{\perp}$  of the nematic in the droplet;  $n_{\parallel} \neq n_p \neq n_{\perp}$ . Typical droplet size  $\sim 2-5 \mu\text{m}$  therefore the random "OFF" state R.G. Dalgarno scatters light strongly. Haze at wide angles and higher voltages due to field loss over the polymer



(ii) Smectic A scattering device

In non field effect the random focal conic texture of the SA phase  $\Rightarrow$  scattering domains. Thus the local directors and layer normals are random. Application of a voltage at a frequency



greater than  $f_c$  (for transition from dynamic scattering to homeotropic texture) gives a clear optical state as layers become ordered. Removal of field leads to bistability and ERASURE is achieved by applying a low frequency voltage ( $f < f_c$ )  $\rightarrow$  ionic motion which disorders the layers & scattering reemerges in the EHD (ElectroHydroDynamic) regime. At high frequencies ( $f > f_c$ ) for  $D_e > 0$  materials dielectric alignment  $\Rightarrow$  clear texture.

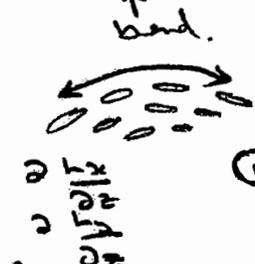
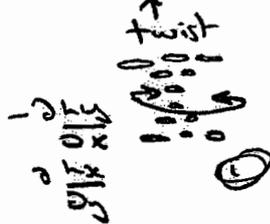
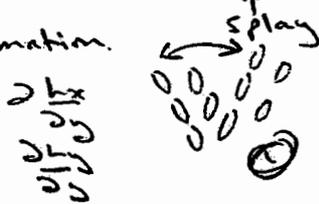
③

Q2

Start with continuum theory for director field  
 i.e. nit molecular. Frank continuum theory

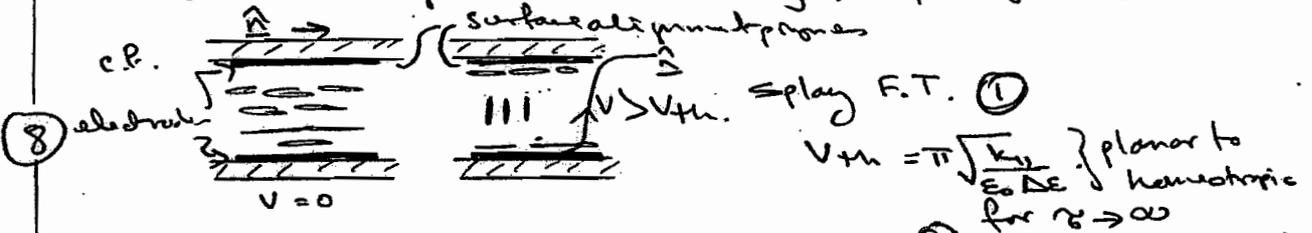
$$F = \frac{1}{2} [ k_{11} (\text{div} \hat{n})^2 + k_{22} (\hat{n} \cdot \text{curl} \hat{n})^2 + k_{33} (\nabla_{\perp} \text{curl} \hat{n})^2 ]$$

Deformation



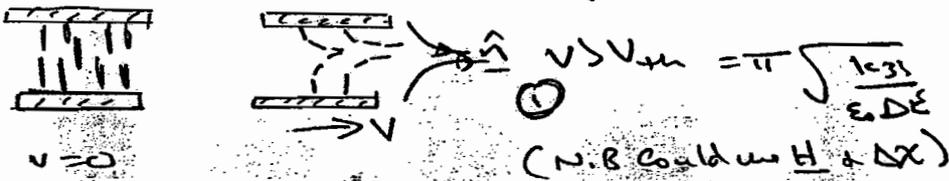
where  $\hat{n}$  is unit vector describing local director.

A Fredericksz-transition occurs when an external voltage (i.e. energy) is applied to induce a (splay / twist / bend) deformation to overcome alignment induced by surface forces.



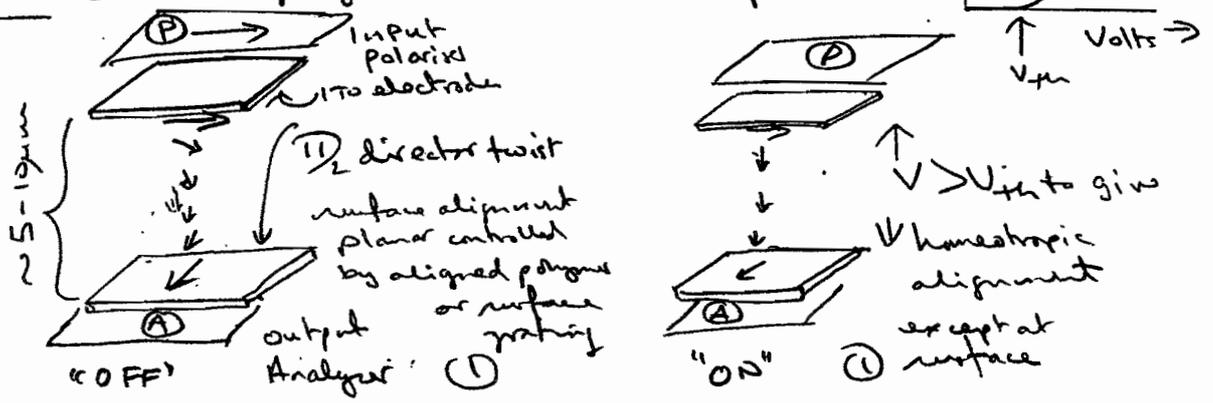
For  $k_{22}$  use twisted alignment layer (c.f. TN cell below)

for  $k_{33}$  use homeotropic alignment and transverse field.



In all three cases apply increasing field and measure optical change due to each F.T.

Since birefringent materials use X polariser



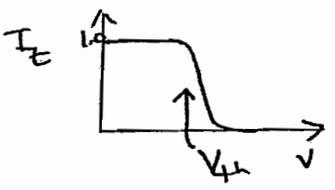
(4)

Q2 Continued

TN cell base on polarisation guiding down twisted structure induced using surface alignment layer (+ chiral additive to ensure  $1/4$  twist & minimise optical losses)  $\Rightarrow$  Transition

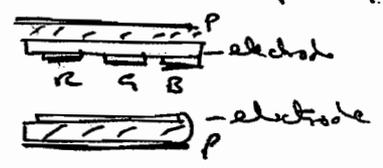
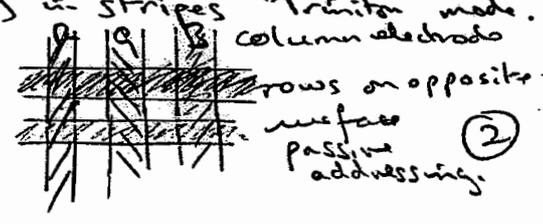
For  $V > V_{th}$  helix unwinds to homeotropic.  $\therefore$  Twist and guiding lost.  $\therefore$  optical polarisation of input light (P) is maintained.  $\therefore$  P & A, as drawn, crossed  $\rightarrow$  extinction. (2)

(6)



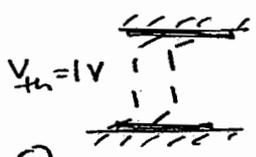
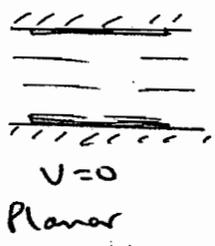
On field removal the surface pinned molecules & chiral (twist) revert back to original "off" state.

To construct colour display use R.S.B. filter at each pixel. Normally in stripes "Trinitron" mode.



(6)

Problem: Planar sample, DE +ve  $\therefore$  splay <sup>(1)</sup> Friedericksz Transition



$$V_{th} = \pi \left[ \frac{k_{11}}{\epsilon_0 \Delta \epsilon} \right]^{1/2} \quad (1)$$

$$\therefore k_{11} = \frac{V_{th}^2 \cdot \epsilon_0 \cdot \Delta \epsilon}{\pi^2} \quad (1)$$

$$= \frac{1 \times 8.854 \times 10^{-12} \cdot 16}{\pi^2} \text{ [SI]}$$

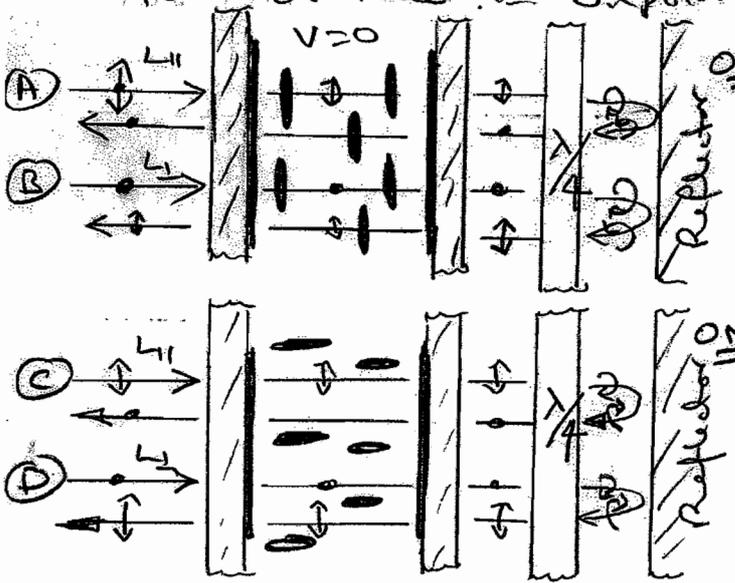
$$k_{11} = 14.4 \times 10^{-12} \text{ N} \quad (2)$$

(6)

### ③ Dye Guest-Host Nematic

{ NB The candidates could describe a DGH in incorporating a  $\lambda/4$  plate or a White Taylor Display. Either is acceptable }

Polarizer Free. i.e. Unpolarised light



OFF In Dye Guest Host absorbing dye with a +ve absorption anisotropy i.e. uniaxial

$\epsilon_{||} - \epsilon_{\perp} = \Delta\epsilon > 0$   
 $\epsilon_{||} \neq 0$   
 is incorporated in a nematic matrix. Thus the major absorption axis follows the director  $\underline{n}$ .

②  $V > V_{th}$ . Quarter wave p.A. at  $45^\circ$  to dir axis

Therefore in OFF state no planar alignment. ①

Construction in case pointer  $\lambda/4$  plate matched to peak absorption  $\lambda$ .

OFF State For unpolarised light consider two components  $\parallel$  and  $\perp$  to director  
 $L_{||}$  to  $\underline{n}$  Route ①  $L_{||}$  absorbed partially traversing device to  $\lambda/4$  plate.  $L_{||}$  changed to circular polarisation of one handedness which is reversed on reflection ( $\pi$  phase change). After reflection & passage thro'  $\lambda/4$  plate component left is rotated into  $L_{\perp}$  direction which then passes back out. ②

$L_{\perp}$  to  $\underline{n}$  Route ②.  $L_{\perp}$  not absorbed on initial pass - follows same process as ① but reaches nematic layer with polarisation rotated by  $90^\circ$  (i.e. in  $L_{||}$  direction) after reflection & passage via  $\lambda/4$  plate. This component is then absorbed by planar dye. Thus in 'OFF' state all polarisations absorbed. ②

In ON case  $L_{||}$  and  $L_{\perp}$  is always in direction of  $\underline{\epsilon}_{\perp}$  due to uniaxial symmetry.  $\therefore$  No absorption. So for white incident light the 'OFF' state is coloured (by subtraction) & ON state is white.  $\therefore$  Use either Red, Green, or Blue

transmitting dyes. Wavelength specific because of  $\lambda/4$  plate. ①

6

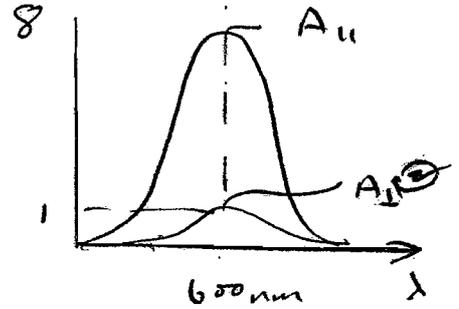
3 cont....

$\lambda_{peak} = 600\text{nm}$ , Dichroic Ratio =  $\frac{Abs_{||}}{Abs_{\perp}} = \frac{1+2S}{1-S} = DR$   
 (Definition) ②

where  $S = \text{Order Parameter}$ .

6  $A_{||}$  for planar 'off' alignment ①

$A_{\perp}$  for perpendicular 'on' alignment in electric field.



Calc  $\Rightarrow S = \frac{1+2S}{1-S}$

$\therefore S = 0.7$  ①

DGH versus TN Cell.

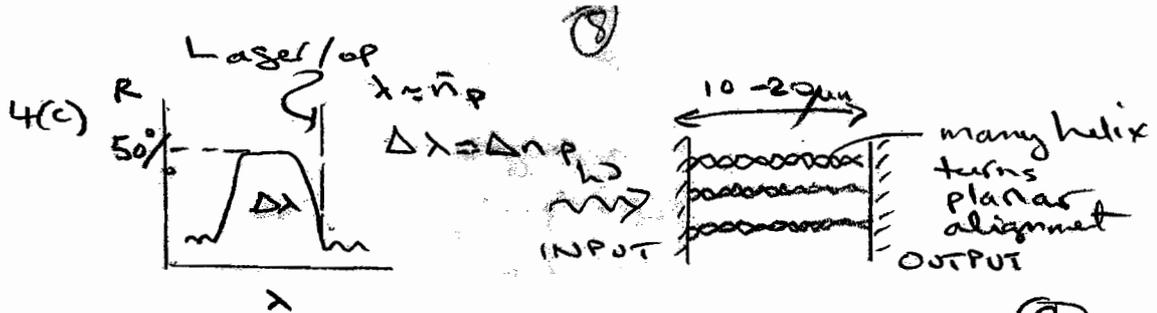
4

Construction  $\rightarrow$

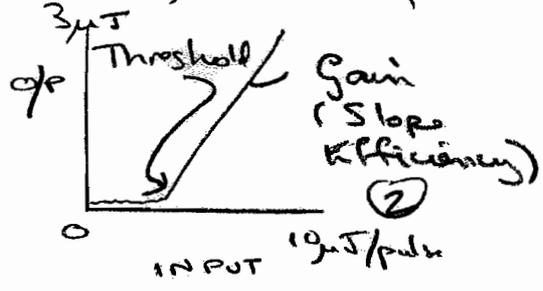
	DH	TN
Brightness	V. good	< good
Contrast	Medium	Good
View Angle	V. good	Restricted
Complexity	1 or 0 Polaris	2 Polaris
Multiplexing	Poor	good.

3

Make device emissive by using a fluorescent dye  
 $\therefore$  absorbed energy re-emitted at lower or higher  $\lambda$  ①

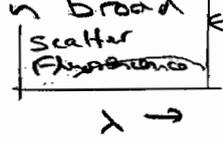


The  $N^*$  chiral structure sets up a PB for CP light with planar surface alignment.  
 Gain medium: Add fluorescent laser dye to  $N^*$  phase (DCM, Pyromethane etc). Pump with frequency doubled YAG laser ( $\lambda_{in} = 532\text{nm}$ ). Choose dye that absorbs at 532nm. Up conversion to longer  $\lambda$  for output fluorescence. Light internally reflected until DOS  $\rightarrow \alpha$  at the band edge and emission occurs ( $\sim 600\text{nm}$ ). Position output fluorescence near BE for max gain

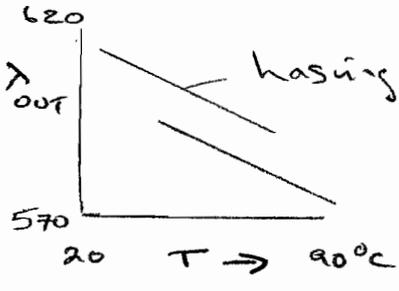
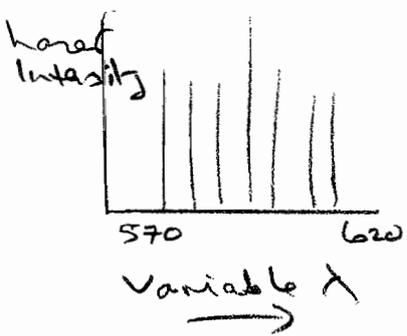


Well defined threshold and high gain & slope efficiencies  
 Beam profile ( $\sim 1\text{mm}$ ) gaussian  
 Directional, coherent  
 Narrow linewidths ( $\sim 0.01\text{nm}$ )

Differentiated from other emission which is broad band and multidirectional.  
 O/P pulses  $\sim \text{ns}$  Efficiency depends on  $(S)$   
 V. high  $\eta \sim 60-70\%$



Tuning can be achieved by choosing an  $N^*$  material in which  $p$  the helix pitch is a fct. Slowly ramp  $T \rightarrow \lambda$  changes as helix starts to unwind. Thus the resonant conditions change and  $\lambda \propto p(T) = f(T)$



Pitch jumps at surface give different modes

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